

# Innovative NDT technique based on ferrofluids for detection of surface cracks

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## Abstract

An innovative NDT technique is proposed for surface inspection of materials not necessarily magnetic or conductive, based on local magnetic field variations due to ferrofluid deposited in the cracks. The feasibility of the technique is assessed preliminarily, based on signal detectability without applied external magnetic field, and under applied DC fields. The signals (local magnetic flux density variations) are quantified analytically, experimentally and numerically. The model agrees well with the tests, showing that the signal increases with the applied field strength, up to the saturation magnetization of the ferrofluid, and decreases with the distance to the crack longitudinal axis, and thus it can provide useful estimations of the signal. The proposed technique, requiring application of external fields to magnetize the ferrofluid to enhance the signal, seems promising: the model suggests that signals associated to cracks significantly smaller than surface cracks in a target application like aircraft skin panel inspection NASA STD-5009 are easily detectable with commercial magnetometers.

**Keywords:** surface flaw; magnetic particle; ferrofluid; aluminium alloys; composite materials

## 1. INTRODUCTION

Early crack detection and monitoring is critical for, *inter alia*, aviation safety and, for this purpose, Non-Destructive Testing (NDT) is an indispensable tool in both production and maintenance. That is why the aerospace sector is one of the largest customers for the NDT industry: in 2008 and 2010, the global expenditure on NDT equipment was slightly over \$1 billion [1,2], and was forecasted to grow up to \$1.3 billion by 2013 [2], and \$1.4 billion by 2017 [1]. Thus, research on NDT solutions for aerospace components that enhance safety and reduce costs is of paramount importance for the NDT and aerospace industries.

In this work, an innovative NDT technique is proposed for surface inspection, based on detection of local magnetic field variations due to accumulation of a ferrofluid in surface cracks. Ferrofluids are colloidal suspensions of small surfacted magnetic particles in a liquid carrier [3]; typically, iron oxide nanoparticles in a Newtonian fluid that can be polar (water) or non-polar (organic solvents). Ferrofluids can be magnetized by applying an external magnetic field, as it forces the magnetic dipole moments of the particles in suspension to align with the direction of the applied field [3]. The objective of this work is to make a preliminary assessment of the feasibility of the proposed NDT technique, based on signal detectability without applied magnetic fields, and under applied direct current (DC) magnetic fields. For this purpose, investigations are conducted to quantify analytically, experimentally and numerically the local magnetic field variations due to presence of a ferrofluid in surface cracks machined in plates of aluminium alloy (AA) 2024-T3. This alloy was chosen as it is widely used in applications for which fatigue resistance is critical, like in skin panels of military aircraft [4] and commercial civil aviation aircraft [5]. The ultimate goal of this research is to implement the proposed technique, meeting the accuracy, reliability and safety requirements of NDT applications in the aerospace industry, while trying to reduce the inspection costs. The latter can be achieved through a combination of reductions in equipment cost, inspection time, operator training, etc., while guaranteeing suitability to a wide range of aerospace NDT applications. Namely, the proposed technique would be applicable for surface inspection of materials not necessarily magnetic or conductive, e.g., aluminium alloys and Carbon or Glass Fibre Reinforced Polymers.

### **1.1 Dimensions of the studied cracks**

The minimum detectable crack depends on the inspection method [6]. In aircraft structural design, the initial crack depends on the component and the type of flaw evaluated [7]. Namely, for fail-safe involving surface flaws, an initial damage of 1.27 mm (3.18 mm for slow-flaw growth) is assumed for pre-service inspections with high standard NDT, while 6.35 mm is assumed for in-service inspections. In structural applications of aluminium alloy panels in aircraft, the most common NDT method for crack monitoring is General Visual Inspection (GVI) [8]. For GVI, the length of the detectable crack ranges from 5.08 to 12.70 mm. For other NDT techniques comparable to the subject of this work, the minimum detectable surface crack is: 1) for dye penetrant testing (PT), 6.36 mm long, 0.64 mm deep, or 3.82 mm long, 1.91 mm deep; 2) for eddy current testing (ET), 5.08 mm long, 0.51 mm deep, or 2.54 mm long, 1.27 mm deep; and 3) for magnetic particle testing (MT), 9.56 mm long, 0.97 mm deep,

or 6.36 mm long, 1.91 mm deep [6]. Our purpose is to determine if the proposed NDT concept would allow detection of surface cracks with these dimensions or smaller, given that one of the target applications is aircraft skin panel inspection. However, for the preliminary study for the proof of concept, cracks of larger dimensions were used. Once the experimental measurements validated the model results for the local magnetic field variations, further computations were made for cracks with the aforementioned dimensions.

## **2. MATERIALS & METHODS**

### **2.1 Tested specimens**

The tested specimens are rectangular plates 100 mm long, 20 mm wide and 2 mm thick, machine cut from sheet of as-received commercial AA 2024-T3. The thickness matches typical values for aircraft skin panels, e.g., 1–1.6 mm for plain panels without holes and 2–3 mm for heavy loaded panels like those in the wing [9]. Using a metal saw, simulated cracks were machined in the surface of the samples along the longitudinal symmetry axis. Simulated cracks can be used in research instead of real cracks grown by fatigue during operation or dynamic testing [10]. The reference crack was 60 mm in length  $l$ , 1.50 mm in width  $w$ , and 0.70 mm in depth  $d$ . In subsequent series of tests, cracks with  $l$  in the range 34.86–66.40 mm,  $w$  in the range 1.97–2.85 mm, and  $d$  in the range 0.52–0.65 mm were used. Finally, for the transversal tests, a crack 12.75 mm long, 0.95 mm wide and 0.60 mm deep, machined in a plate 20 mm long, 20 mm wide, was used.

### **2.2 Ferrofluids**

The magnetic particles in ferrofluids are generally spherical and with diameters ranging from 5 to 15 nm [3,11,12]. Also, each particle is generally a single magnetic domain, i.e., it is an entity with a single magnetic moment, and thus behaves as a single magnetic dipole. The volume fraction of ferrofluids, i.e., the volume percentage of magnetic solid material with respect to the total volume, is usually 5 to 15%. Four ferrofluids have been considered: three generic suspensions of ferromagnetic nanoparticles and the commercial ferrofluid N-503 from Sigma Hi-Chemical Inc. (the properties of the latter are shown in Table 1, as provided by the manufacturer). The generic suspensions are made of iron ( $\alpha$ -Fe), magnetite ( $\text{Fe}_3\text{O}_4$ ) and maghemite ( $\gamma$ - $\text{Fe}_2\text{O}_3$ ) nanoparticles, respectively (their properties are shown in Table 2). For simplicity, the idea of synthesizing a custom-made ferrofluid was disregarded and it was decided to use only the commercial ferrofluid for the experiments, and consequently also for

the numerical simulations. Nevertheless, for comparison purposes, the theoretical development and calculations in Section 3 were applied to all four ferrofluids.

**Table 1** Properties of commercial ferrofluid N-503 supplied by Sigma Hi-Chemical Inc.

Concept	Value
Liquid carrier	Iso-paraffin
Type of magnetic particles	Magnetite ( $\text{Fe}_3\text{O}_4$ ) particles
Average diameter	10 nm
Volume fraction	8.9%
Saturation magnetization $M_{sat}$	43.8 kA/m
Dynamic viscosity $\eta$	20.6 mPa·s, at 293 K

**Table 2** Properties of iron ( $\alpha$ -Fe), magnetite ( $\text{Fe}_3\text{O}_4$ ) and maghemite ( $\gamma$ - $\text{Fe}_2\text{O}_3$ ) particles at room temperature (RT).

Type	Concept	Value	Reference	Observations
$\alpha$ -Fe	Critical diameter <sup>a</sup>	7 nm	[13]	
	Mass/specific saturation magnetization	91.3 A·m <sup>2</sup> /kg	[14]	8.4 nm diameter particles
	Density	7870 kg/m <sup>3</sup>		
$\text{Fe}_3\text{O}_4$	Critical diameter <sup>a</sup>	128 nm		
	Mass/specific saturation magnetization	46 A·m <sup>2</sup> /kg	[15]	19 nm diameter particles
	Density	5000 kg/m <sup>3</sup>		
$\gamma$ - $\text{Fe}_2\text{O}_3$	Critical diameter <sup>a</sup>	166 nm		
	Mass/specific saturation magnetization	36.6 A·m <sup>2</sup> /kg	[16]	9 nm diameter particles
	Density	4600 kg/m <sup>3</sup>		

<sup>a</sup> The critical diameter is a threshold below which a magnetic particle is a single magnetic domain. Above this critical diameter, the particle consists of multiple magnetic domains.

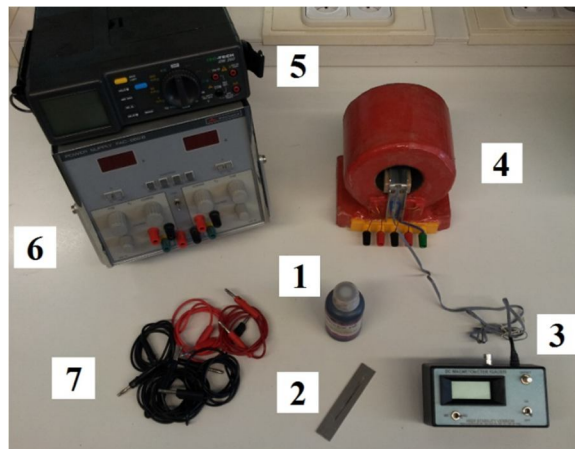
### 2.3 Experimental setup & methodology

The magnetic flux density  $\bar{B}$  was measured before and during application of a DC magnetic field, with and without ferrofluid in the reference crack. Fig. 1 shows the experimental set.

The custom-made bracket (see Appendix A and Online Resource 1) consists of two concentric Cu wire coils covered by a protective shell, and the specimen support, located such that the crack in the sample is aligned with the revolution axes of the coils. A dispenser was used for depositing the ferrofluid in the cracks. A power source supplied DC power to the coils for generating DC magnetic fields. The intensity and voltage of the current were measured with a multimeter, while  $\bar{B}$  was measured with the AlphaLab magnetometer (with resolution 0.001 mT and accuracy  $\pm 2\%$ ). The reference position of the Hall probe was 3.5 mm below the crack centre, oriented to measure the component of  $\bar{B}$  in the direction of the revolution axes of the coils. Measurements were taken sequentially in a series of cases summarized in Table 3.

For the cases where power is supplied to the coils, the tested voltages ranged from 1 to 17 V, with the coils generating applied external fields with strengths  $H$  ranging from around 1 to 16 kA/m. Higher applied fields were not used because the gain in ferrofluid magnetization is counterbalanced by a much poorer magnetometer resolution for the signals above 20 mT.

For another sample, tests were realized with the Hall probe located at increasing distance from the crack, with  $H$  of 8 kA/m. Also, tests with surface cracks of different dimensions were conducted, with the Hall probe located back in the reference position, with  $H$  of 8 and 16 kA/m. Finally, tests were made with a crack oriented perpendicular to the revolution axes of the coils, and thus to the applied field, to explore the response if the defect does not lay parallel to this field. In this case, the crack and plate were smaller to allow introducing the plate inside the support transversally. The tests were performed in a laboratory with low electromagnetic noise. Prior to testing, the background field was zeroed by applying an appropriate offset in the magnetometer. Each of the test results shown in Section 4 is the average of three individual measurements (the error bars in the figures represent one standard deviation).



**Fig. 1** Experimental set: ferrofluid (1), test specimen (2), magnetometer and Hall probe (3), custom-made bracket (4), multimeter (5), power source (6) and connection cables (7)

**Table 3** Summary of experimental conditions applied sequentially during the tests.

DC test cases	Sample on support	Ferrofluid in the crack	DC power supplied to coils
DC.1	No	No	No
DC.2	Yes	No	No
DC.3	Yes	No	Yes
DC.4	Yes	Yes <sup>a</sup>	No (before magnetizing)
DC.5	Yes	Yes <sup>a</sup>	Yes (ferrofluid magnetized)

<sup>a</sup> Namely, the volume of commercial ferrofluid in the reference crack is approximately 50 mm<sup>3</sup>.

## 2.4 Simulations with COMSOL Multi-physics

The local magnetic field variations due to the presence of ferrofluid in the reference crack were computed numerically with the AC/DC Module of COMSOL Multi-physics, a commercial finite element analysis software package for coupled physics phenomena and engineering applications, developed by COMSOL, Inc., Palo Alto, CA, USA. The mesh was a user-controlled unstructured 3D mesh of tetrahedral elements. The properties of the ferrofluid, the air and the test plates used in the simulations are summarized in Table 4. The electrical conductivity, permittivity and permeability of the commercial ferrofluid were not supplied by the manufacturer. The conductivity, measured with a Hach HQ440d Multi-Parameter Meter, was  $5 \times 10^{-5}$  S/m. The relative permittivity used in the simulations was that of the iso-paraffin. The relative permeability ( $\mu_r = 83.6$ ) resulted from a calibration based on fitting the simulation results to the measured field in the probe position under an applied DC field of 16 kA/m. The solver selected for the simulations was the iterative FGMRES.

**Table 4** Properties of the ferrofluid, air and AA 2024-T3 at 293 K, as used in COMSOL simulations.

Material	Electrical conductivity (S/m)	Relative permittivity (-)	Relative permeability (-)
Magnetite	$9.61 \times 10^8$ [17]	15 [18], 33.7–81 [19]	1.4–2.0 [18]
Iso-paraffin	Insulator	1.9 [20]	–
Ferrofluid	0 (measured) <sup>a</sup>	2	1.5–6 [21], 3–96 [22]
Air	0 (COMSOL database) <sup>a</sup>	1 (COMSOL database)	1 (COMSOL database)
AA 2024-T3	$1.60 \times 10^7 - 1.89 \times 10^7$ [23]	1.44 [24], 8.0 [25]	1.00002 [26]

<sup>a</sup> The electrical conductivity measured for the commercial ferrofluid was 0 S/m, and the conductivity for air in COMSOL's materials database is 0 S/m. This is reported to cause problems in the solver. Instead, it is recommended to use a very small conductivity value, and so 10 S/m was used.

## 3. THEORY & CALCULATIONS

The ferrofluid in the crack is modelled as a magnetic dipole with semi-length  $a$  and radius  $R$ , such that it has the length of the crack and a volume equal to the volume of ferrofluid placed in the crack. The origin of the reference system used in this work is the dipole centre. The  $z$  axis is the revolution axis of the dipole, parallel to the crack and revolution axes of the coils. The  $y$  axis is perpendicular to the  $z$  axis, pointing opposite to gravity (see Online Resource 1).

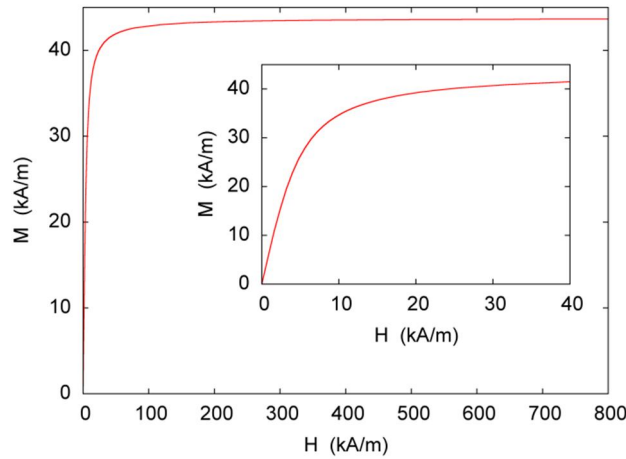
### 3.1 Equilibrium (DC) magnetization of the ferrofluid

Ferrofluids can be magnetized by external DC magnetic fields [3]. This phenomenon increases with  $H$ , up to reaching  $M_{sat}$  [27]. The following hypotheses are assumed:

1. The ferrofluid is mono-disperse, i.e., the particles are all identical in properties, composition, dimensions and shape (assumed spherical, with diameter 10 nm).

2. For being conservative, the generic ferrofluids are considered dilute colloidal suspensions with volume fraction of 7%, and  $M_{sat}$  is the smallest in the literature for the corresponding type of particles (see Table 2).
3. Each magnetic particle is a single magnetic domain. This is coherent with the critical diameters found for the studied ferrofluids (see Table 2).
4. The magnetization  $M$  is homogeneous within the ferrofluid.
5. The magnetic particles in suspension are isotropic and non-interacting.

If a ferrofluid is a collection of individual, non-interacting and mono-disperse magnetic dipoles, a theory by Langevin [3] states that, under a field applied in the  $z$  axis, the ferrofluid non-dimensional magnetization along that axis is  $\tilde{M}_z = M/M_{sat} = \coth(\alpha) - 1/\alpha$ , with the Langevin parameter  $\alpha = m\mu_0 H/k_B T$ , where  $m$  is the magnetic moment of the particles ( $15.5 \times 10^{-19} \text{ A}\cdot\text{m}^2$ , for magnetite [28]),  $\mu_0$  is the free space permeability constant,  $k_B$  is the Boltzmann constant, and  $T$  is the temperature. Fig. 2 shows the magnetization curve computed for the commercial ferrofluid at 293 K. If the applied field is suppressed,  $M$  relaxes to a new equilibrium state following an exponential decay being  $\tilde{M}_z = \frac{1}{3}\alpha \exp(-t/\tau)$ , where  $\tau$  is the relaxation time [3].



**Fig. 2** Magnetization of the ferrofluid  $M$  vs. applied field strength  $H$ , at 293 K, by Langevin's theory [3]

### 3. 2 Magnetic field of the ferrofluid

For applied DC fields or absence of applied field, crack detection would be based on the local variations of  $\vec{B}$  due to the ferrofluid in the crack. A magnetic field can be calculated with Maxwell's equations [29]:  $\vec{B}(\vec{r}) = \vec{B}_I(\vec{r}) + \vec{B}_{II}(\vec{r}) = \mu_0 \vec{M}(\vec{r}) - \mu_0 \nabla \varphi^*(\vec{r})$ , where  $\vec{r}$  is the position vector of the point where the field is evaluated,  $\vec{M}$  is the magnetization, and  $\varphi^*$  is a

scalar potential. Inside the dipole,  $\bar{B}_I$  depends on the ferrofluid magnetization. Outside the dipole,  $\bar{B}_I$  is associated to the electromagnetic noise. In this work, this term is null in the  $z$  axis due to the offset applied to the magnetometer prior to testing. Thus, the magnetic flux density associated to the ferrofluid  $\bar{B}_{II}$  is [29]:

$$\text{Eq. 1} \quad \bar{B}_{II}(\vec{r}) = -\mu_0 \nabla \varphi^*(\vec{r}) = -\mu_0 \nabla \frac{1}{4\pi} \int_{\Omega} \bar{M}(\vec{r}') \cdot \frac{\vec{r}-\vec{r}'}{|\vec{r}-\vec{r}'|^3} dV'$$

where  $\Omega$  is the volume of magnetized material (i.e., the volume of ferrofluid) and  $\vec{r}'$  is the position vector of a differential magnetic dipole. If the ferrofluid in the crack is modelled as a cylindrical dipole with semi-length  $a$  and radius  $R$ , when the applied field is aligned with the dipole longitudinal axis, i.e., the  $z$  axis, Eq. 1 becomes:

Eq. 2

$$\bar{B}_{II}(\vec{r}) = -\frac{1}{4\pi} \mu_0 M \pi R^2 \left[ \left( \frac{x}{(x^2+y^2+(z+a)^2)^{3/2}} - \frac{x}{(x^2+y^2+(z-a)^2)^{3/2}} \right) \bar{i} + \left( \frac{y}{(x^2+y^2+(z+a)^2)^{3/2}} - \frac{y}{(x^2+y^2+(z-a)^2)^{3/2}} \right) \bar{j} + \left( \frac{z+a}{(x^2+y^2+(z+a)^2)^{3/2}} - \frac{z-a}{(x^2+y^2+(z-a)^2)^{3/2}} \right) \bar{k} \right]$$

To explore the response if the applied field is perpendicular to the crack/dipole longitudinal axis, the dipole is rotated 90° to align it with the  $x$  axis. Then, Eq. 1 becomes:

$$\text{Eq. 3} \quad \bar{B}_{II}(\vec{r}) = -\frac{1}{4\pi} \mu_0 M \pi R^2 \left[ \left( \frac{z}{((x+a)^2+y^2+z^2)^{1/2}(y^2+z^2)} - \frac{z}{((x-a)^2+y^2+z^2)^{1/2}(y^2+z^2)} + \frac{z(x-a)^2}{((x-a)^2+y^2+z^2)^{3/2}(y^2+z^2)} - \frac{z(x+a)^2}{((x+a)^2+y^2+z^2)^{3/2}(y^2+z^2)} \right) \bar{i} + \left( \frac{yz(x-a)}{((x-a)^2+y^2+z^2)^{3/2}(y^2+z^2)} - \frac{yz(x+a)}{((x+a)^2+y^2+z^2)^{3/2}(y^2+z^2)} - \frac{2yz(x+a)}{((x+a)^2+y^2+z^2)^{1/2}(y^2+z^2)^2} + \frac{2yz(x-a)}{((x-a)^2+y^2+z^2)^{1/2}(y^2+z^2)^2} \right) \bar{j} + \left( \frac{x+a}{((x+a)^2+y^2+z^2)^{1/2}(y^2+z^2)} + \frac{z^2(x-a)}{((x-a)^2+y^2+z^2)^{3/2}(y^2+z^2)} + \frac{2z^2(x-a)}{((x-a)^2+y^2+z^2)^{1/2}(y^2+z^2)^2} - \frac{x-a}{((x-a)^2+y^2+z^2)^{1/2}(y^2+z^2)} - \frac{z^2(x+a)}{((x+a)^2+y^2+z^2)^{3/2}(y^2+z^2)} - \frac{2z^2(x+a)}{((x+a)^2+y^2+z^2)^{1/2}(y^2+z^2)^2} \right) \bar{k} \right]$$

### 3.3 Magnetic field of the ferrofluid in the xy plane

After preliminary computations with Eq. 2, the field associated to the ferrofluid was observed to be very significant in the  $z$  axis close to the dipole tip (see Online Resource 2), but it is not possible to take measurements there. Also, the field vanishes dramatically with the distance to the  $z$  axis, and the direction of the field lines changes significantly in short distances. Thus, it

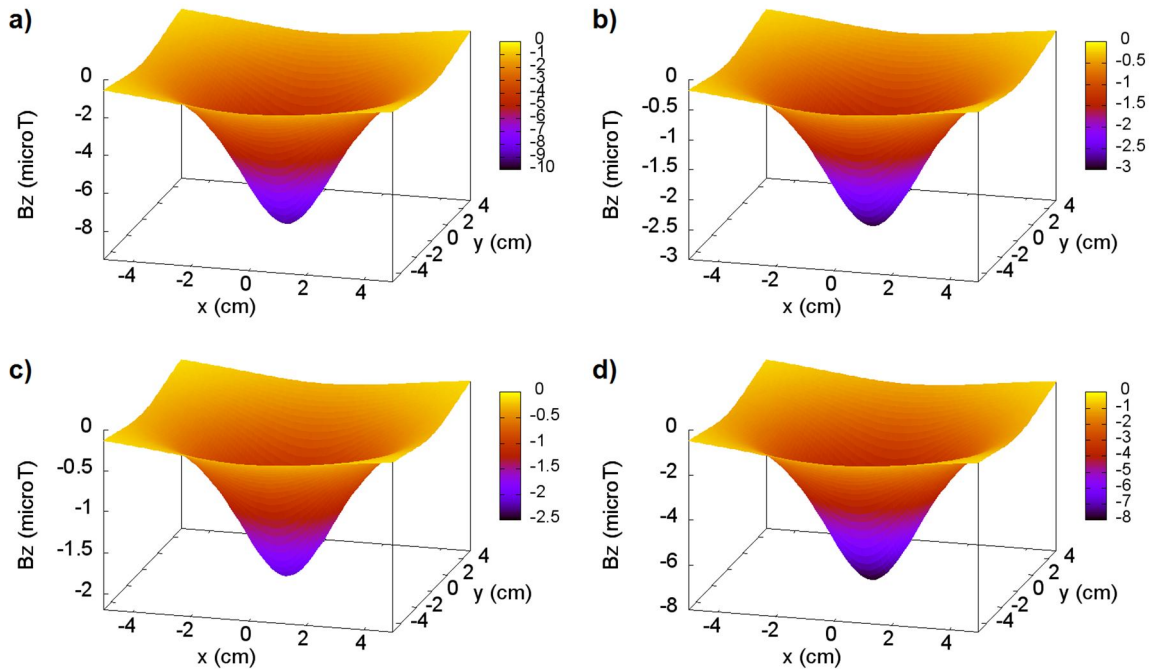


is very complex to establish the most appropriate position and orientation of the probe if it is to be placed in the vicinity of the tip of the crack/dipole but separated from the  $z$  axis (the magnetometer is only able to measure the field in one direction). Conversely, in the  $xy$  plane the field is expected to have component only in the  $z$  axis,  $B_z$ , which facilitates taking measurements. For these reasons, the study was focused on the  $xy$  plane, where the signal when the applied field is aligned with the dipole longitudinal axis, as obtained from Eq. 2, is:

$$\text{Eq. 4} \quad |\bar{B}_{II}(\bar{r})|_{z=0} = B_z \bar{k} = -\frac{1}{4\pi} \mu_0 M \pi R^2 \left[ \frac{2a}{(x^2+y^2+a^2)^{3/2}} \right] \bar{k}$$

For the reference crack, modelled by a dipole with  $a$  of 30 mm and  $R$  of 0.51 mm (see Online Resource 1), Fig. 3 shows  $B_z$  as obtained from Eq. 4 for the generic ferrofluids and the commercial ferrofluid, all at their  $M_{sat}$ . On the other side, if the applied field is perpendicular to the dipole longitudinal axis, from Eq. 3, the signal in the  $xy$  plane is:

$$\text{Eq. 5} \quad |\bar{B}_{II}(\bar{r})|_{z=0} = B_z \bar{k} = -\frac{1}{4\pi} \mu_0 M \pi R^2 \left[ \frac{x+a}{((x+a)^2+y^2)^{1/2}y^2} - \frac{x-a}{((x-a)^2+y^2)^{1/2}y^2} \right] \bar{k}$$



**Fig. 3** For the reference crack (modelled by a dipole with  $a$  of 30 mm and  $R$  of 0.51 mm), analytical  $B_z$  vs.  $x$  and  $y$  coordinates, in the  $xy$  plane, for ferrofluids made of a) iron ( $\alpha$ -Fe) nanoparticles, b) magnetite ( $\text{Fe}_3\text{O}_4$ ) nanoparticles, c) maghemite ( $\gamma$ - $\text{Fe}_2\text{O}_3$ ) nanoparticles, and d) for the commercial ferrofluid, all at their  $M_{sat}$

## 4. RESULTS & DISCUSSION

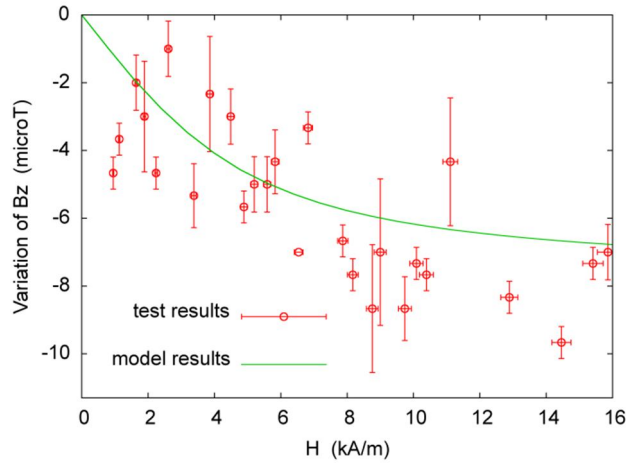
Measurements under DC.1 test conditions provided the background magnetic field in the laboratory. Then, the magnetometer reading was set to zero by applying appropriate offset. Measurements under DC.2 test conditions served for establishing a baseline for identifying variations of  $B_z$  due to presence of ferrofluid in the crack before supplying DC power to the coils (i.e., before magnetizing the ferrofluid). These measurements also confirmed that the samples do not modify the background field, as expected since AA 2024-T3 is diamagnetic. Analogously, measurements under DC.3 test conditions served for establishing a baseline for identifying variations of  $B_z$  while supplying DC power to the coils (i.e., while the ferrofluid is being magnetized by application of an external DC field). Measurements under DC.4 and DC.5 test conditions allowed computing the variations of  $B_z$  in two different hypothetical versions of the NDT technique:

- **Tech-DC.I: Variation of  $B_z$  due to presence of ferrofluid in the crack when the ferrofluid has not been previously magnetized:** For the reference crack, this variation (the difference between measurements obtained in cases DC.4 and DC.2) could not be determined, since the fields are below the sensor resolution.
- **Tech-DC.II: Variation of  $B_z$  due to presence of ferrofluid in the crack when the ferrofluid is being magnetized:** For the reference crack, Fig. 4 shows this variation (the difference between measurements obtained in cases DC.5 and DC.3), and also the model results derived from the theoretical development in Section 3, where the ferrofluid magnetization increases with  $H$ . Although the test results show high dispersion, probably due to errors in measuring the crack size and the volume of ferrofluid deposited from test to test, the average error of the model results is  $-2.2 \pm 51.2$  % and both follow the same trend.

### 4.1 Comparison of test results for Tech-DC.I and Tech-DC.II

For the version Tech-DC.I, the local variation of  $B_z$  due to the ferrofluid when it has not been previously magnetized is below the minimum resolution of the AlphaLab magnetometer. As expected, the signal is much more significant for the version Tech-DC.II, and increases with  $H$  since the ferrofluid equilibrium magnetization also increases [27]. The variation for the highest tested field is  $-7 \mu\text{T}$ . Although according to the theory it should be the maximum signal, it is not the case (the highest measured variation is  $-9.7 \mu\text{T}$ , for  $H$  of  $14.5 \text{ kA/m}$ ), probably due to experimental error. The maximum signal-to-noise ratio is  $-69 \text{ dB}$ ,

corresponding to a variation of  $-5 \mu\text{T}$ , achieved for the smallest  $H$ . An NDT technique based on the version Tech-DC.II would require the operator to scan at least two times the inspection surface: first to clean the surface and spread the ferrofluid (a likely drawback is that the surface has to be clean and the cracks must not be polluted, like for PT), and second to apply the external field while measuring the response. Thus, the inspection equipment should be able to simultaneously generate a field to magnetize the ferrofluid (preferably up to  $M_{sat}$ , from the basis of signal detectability), and to measure the local variations of  $\bar{B}$ . Crack detection capabilities can be enhanced using higher sensitivity sensors (magnetometers with resolutions down to 0.1 nT are common) or ferrofluids with higher  $M_{sat}$ . Finally, a 3-components Hall probe would be more appropriate, since the operator in an NDT inspection does not know the crack orientation and the direction in which the induced field is higher.

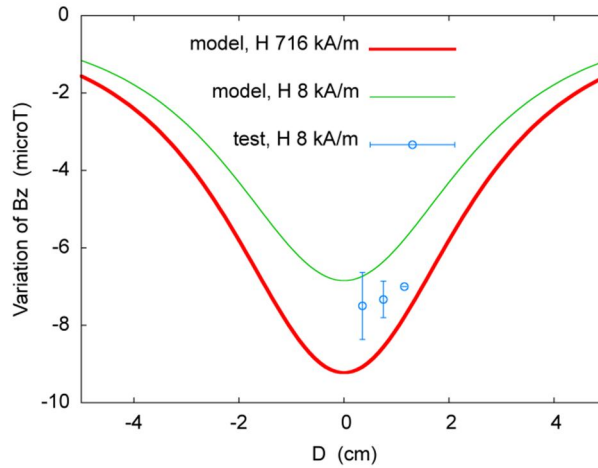


**Fig. 4** For version Tech-DC.II at 293 K for the reference crack, analytical and measured variation of  $B_z$  (difference between measurements obtained in cases DC.5 and DC.3) vs. applied field strength  $H$  (case DC.3), for the commercial ferrofluid, with the Hall probe located at 3.5 mm from the dipole axis

#### 4.2 Comparison of model results with test results for Tech-DC.II

The model results in the  $xy$  plane (see Eq. 4 and Fig. 3) show that, as expected, the absolute value of the variation of  $B_z$  due to the ferrofluid in the crack decreases with the distance to the dipole axis in the  $xy$  plane  $D = \sqrt{x^2 + y^2}$ . For sample #2 (see crack dimensions in Table 5), tests were made with the Hall probe located at increasing distance from the dipole axis (namely, at 3.5, 7.5 and 11.5 mm), with  $H$  of 8 kA/m. Fig. 5 shows the measured signal, i.e., the variation of  $B_z$  due to presence of ferrofluid in the crack when the ferrofluid is being magnetized (the difference between measurements in cases DC.5 and DC.3), compared to model results. Model results for  $H$  slightly above 700 kA/m are also shown, as a

representative condition at which the ferrofluid has reached  $M_{sat}$ . The tests confirm that the signal decreases with  $D$ , but apparently at a slightly slower rate to that shown by the model (the error of the model increases with  $D$ , and in average it is  $-13.7 \pm 3.0 \%$ ). As expected, the sensor should always be placed as close as possible to the inspection surface.



**Fig. 5** For version Tech-DC.II at 293 K for sample #2, analytical and measured variation of  $B_z$  (difference between measurements obtained in cases DC.5 and DC.3) vs. distance to dipole axis in the  $xy$  plane  $D$ , for the commercial ferrofluid, for applied field strength  $H$  of 8 and 716 kA/m (ferrofluid at  $M_{sat}$ )

To further validate the model, tests with surface cracks of different dimensions were made, with the Hall probe located back in the reference position, with  $H$  of 8 and 16 kA/m. Table 5 shows the measured signals compared to model results. Finally, Table 6 shows the model results and test results (again, the difference between readings in cases DC.5 and DC.3) for a crack perpendicular to the applied field, with the Hall probe placed at 5.5 mm from the dipole axis. The signals measured for the crack oriented in the direction of the applied external field and perpendicular to it are virtually identical in spite of the distance to the crack being larger in the second case. On the other side, the signal predicted by the model is significantly higher if the defect lays perpendicular to the applied field. Thus, the model seems not so appropriate to estimate the signal for the latter condition. However, the test results suggest that not knowing the direction of the defect when applying the external field may not be very relevant to the performance of the proposed NDT method.

**Table 5** For version Tech-DC.II at 293 K, analytical and measured variation of  $B_z$  for surface cracks of different dimensions for applied field strength  $H$  of 8 kA/m (test 1) and 16 kA/m (test 2), with the Hall probe located at 3.5 mm from the dipole axis. The relative error of the model results is provided.

Sample	Crack dimensions	Dipole	Measured $B_z$ variation	Analytical $B_z$ variation (error)
#2	$l = 66.40$ mm	$a = 33.20$ mm	test 1: $-7.5 \pm 0.9$ $\mu$ T	test 1: $-6.7$ $\mu$ T (-10.2 %)
	$w = 2.85$ mm	$R = 0.609$ mm	test 2: $-10.6 \pm 2.3$ $\mu$ T	test 2: $-7.9$ $\mu$ T (-25.3 %)
	$d = 0.52$ mm			
#3	$l = 50.88$ mm	$a = 25.44$ mm	test 1: $-10.3 \pm 0.5$ $\mu$ T	test 1: $-9.8$ $\mu$ T (-5.0 %)
	$w = 1.97$ mm	$R = 0.567$ mm	test 2: $-13.4 \pm 3.3$ $\mu$ T	test 2: $-11.5$ $\mu$ T (-13.9 %)
	$d = 0.65$ mm			
#4	$l = 34.86$ mm	$a = 17.43$ mm	test 1: $-18.0 \pm 8.6$ $\mu$ T	test 1: $-18.1$ $\mu$ T (0.4 %)
	$w = 2.21$ mm	$R = 0.535$ mm	test 2: $-26.4 \pm 3.4$ $\mu$ T	test 2: $-21.2$ $\mu$ T (-19.7 %)
	$d = 0.52$ mm			
#5	$l = 12.75$ mm	$a = 6.375$ mm	test 1: $-35.0 \pm 0.0$ $\mu$ T	test 1: $-31.0$ $\mu$ T (-11.5 %)
	$w = 0.95$ mm	$R = 0.377$ mm	test 2: $-47.0 \pm 0.8$ $\mu$ T	test 2: $-36.4$ $\mu$ T (-22.6 %)
	$d = 0.60$ mm			

**Table 6** For version Tech-DC.II at 293 K, analytical and measured variation of  $B_z$  for a crack perpendicular to the applied field with strength  $H$  of 8 kA/m (test 1p) and 16 kA/m (test 2p). The relative error of the model results is provided.

Sample	Crack dimensions	Dipole	Measured $B_z$ variation	Analytical $B_z$ variation (error)
#5	$l = 12.75$ mm	$a = 6.375$ mm	test 1p: $-35.0 \pm 0.0$ $\mu$ T	test 1p: $-72.6$ $\mu$ T (107.4 %)
	$w = 0.95$ mm	$R = 0.377$ mm	test 2p: $-47.7 \pm 0.9$ $\mu$ T	test 2p: $-85.3$ $\mu$ T (78.9 %)
	$d = 0.60$ mm			

The average error of the model results in Table 5 is  $-13.5 \pm 8.2$  %, and it falls to  $-5.3 \pm 44.3$  % when taking into account the results in Fig. 4 and Fig. 5 (the transversal tests are not considered). Aside from the limitations of the model, the discrepancies with the test results may be due to the many sources of experimental error, e.g., a miss-alignment of the Hall probe with the  $z$  axis, the probe position error, the thermal energy which induces particle vibration, and even particle interaction, assumed inexistent. But, overall, it seems that the model can provide useful predictions in spite of its simplicity and the lack of knowledge on some important features of the ferrofluid that the manufacturer did not provide.

Now that the model has been validated and the order of magnitude of the error is known, for comparison purposes, the variations of  $B_z$  associated to minimum detectable surface cracks for various classical NDT techniques are computed in the reference position of the Hall Probe (see Table 7). In all cases, the signals would be detectable with the available instrumentation. Tests with cracks with these dimensions have not been realized because it is not possible to control with the necessary accuracy the size of the cracks generated by the metal saw.

**Table 7** For version Tech-DC.II at 293 K, analytical variation of  $B_z$  for the minimum detectable surface cracks for various NDT techniques, for applied field strength  $H$  of 8 kA/m (test 1) and 16 kA/m (test 2).

NDT technique	Crack dimensions	Dipole	Analytical $B_z$ variation
Assumed initial damage in fail-safe	$l = 1.27$ mm	$a = 0.635$ mm	test 1: $-18.7 \mu\text{T}$
	$w = 0.51$ mm <sup>a</sup>	$R = 0.255$ mm	test 2: $-21.9 \mu\text{T}$
	$d = 0.51$ mm <sup>a</sup>		
Assumed initial damage in fail-safe (slow-flaw growth)	$l = 3.18$ mm	$a = 1.59$ mm	test 1: $-37.0 \mu\text{T}$
	$w = 0.51$ mm <sup>a</sup>	$R = 0.255$ mm	test 2: $-43.5 \mu\text{T}$
	$d = 0.51$ mm <sup>a</sup>		
Eddy current testing 1 & general visual inspection 1	$l = 5.08$ mm	$a = 2.54$ mm	test 1: $-41.6 \mu\text{T}$
	$w = 0.51$ mm	$R = 0.255$ mm	test 2: $-48.8 \mu\text{T}$
	$d = 0.51$ mm		
Eddy current testing 2	$l = 2.54$ mm	$a = 1.27$ mm	test 1: $-81.1 \mu\text{T}$
	$w = 0.51$ mm	$R = 0.402$ mm	test 2: $-95.3 \mu\text{T}$
	$d = 1.27$ mm		
Dye penetrant testing 1	$l = 6.36$ mm	$a = 3.18$ mm	test 1: $-49.9 \mu\text{T}$
	$w = 0.51$ mm <sup>a</sup>	$R = 0.286$ mm	test 2: $-58.7 \mu\text{T}$
	$d = 0.64$ mm		
Dye penetrant testing 2	$l = 3.82$ mm	$a = 1.91$ mm	test 1: $-149.3 \mu\text{T}$
	$w = 0.51$ mm <sup>a</sup>	$R = 0.493$ mm	test 2: $-175.5 \mu\text{T}$
	$d = 1.91$ mm		
Magnetic particle testing 1	$l = 9.56$ mm	$a = 4.78$ mm	test 1: $-57.9 \mu\text{T}$
	$w = 0.51$ mm <sup>a</sup>	$R = 0.352$ mm	test 2: $-68.0 \mu\text{T}$
	$d = 0.97$ mm		
Magnetic particle testing 2	$l = 6.36$ mm	$a = 3.18$ mm	test 1: $-149.0 \mu\text{T}$
	$w = 0.51$ mm <sup>a</sup>	$R = 0.493$ mm	test 2: $-175.1 \mu\text{T}$
	$d = 1.91$ mm		
General visual inspection 2	$l = 12.70$ mm	$a = 6.35$ mm	test 1: $-22.0 \mu\text{T}$
	$w = 0.51$ mm <sup>a</sup>	$R = 0.255$ mm	test 2: $-25.9 \mu\text{T}$
	$d = 0.51$ mm <sup>a</sup>		

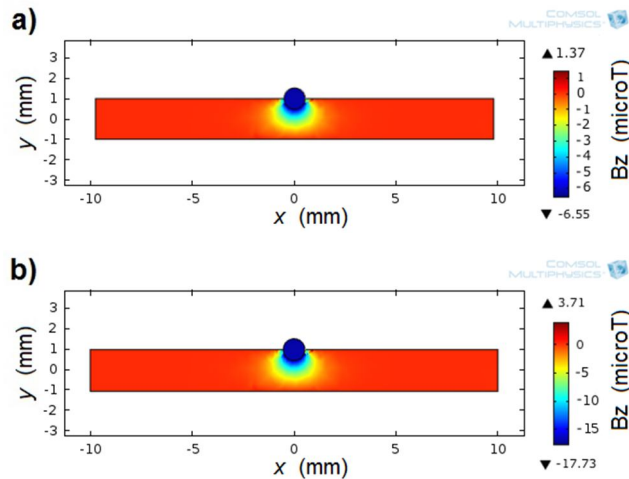
<sup>a</sup> In those cases in which there is no information about the width/depth of the corresponding minimum detectable surface crack, a value of 0.51 mm has been used in the simulations, since it is the minimum width/depth value for the minimum detectable surface cracks for the techniques considered.

From Eq. 4, for a given value of  $w$  and  $d$  (and thus radius of the dipole), it can be derived that, for a given distance of the Hall probe to the dipole axis in the  $xy$  plane,  $D$ , there is a dipole semi-length providing maximum signal:  $a = D/\sqrt{2}$ . For example, for a crack with  $w = d = 0.51$  mm, in the reference position of the probe ( $D$  of 3.5 mm), a maximum in the signal would be obtained for  $a$  of 2.47 mm. This trend can be observed in Table 7. Finally, the model has been used to estimate the minimum detectable surface crack for the proposed NDT technique. For instance, for  $H$  of 8 kA/m, for a crack with  $l$  of 0.001 mm and  $w = d = 0.51$  mm, the signal would be  $-15.4$  nT; for a crack with  $l$  of 1.27 mm and  $R = 0.005$  mm (corresponding to, e.g.,  $w = d = 0.01$  mm), the signal would be  $-7.2$  nT; and for a crack with  $l$  of 0.012 mm and  $R = 0.05$  mm (corresponding to, e.g.,  $w = d = 0.1$  mm), the signal would

be -7.1 nT. All these signals would be perfectly detectable using magnetometers with resolution down to 0.1 nT, which are commercially available and not uncommon. These model estimations suggest that the proposed NDT method has a promising performance.

### 4.3 Comparison with results from simulations with COMSOL

Fig. 6 a) and Fig. 6 b) show contour plots of the variations of  $B_z$  within the cross-section of the sample and the dipole modelling the ferrofluid in the reference crack, in the  $xy$  plane, at 293 K, for  $H$  of 16 and 716 kA/m, respectively. The latter is a representative condition at which the ferrofluid has reached  $M_{sat}$ . The purpose of these simulations is to compare the signal within the aluminium plate as provided by COMSOL with the signal computed using the theoretical development in Section 3, at equivalent distances but out of the plate. For instance, Table 8 shows the signal at half thickness of the plate, below the dipole centre, as given by COMSOL, and the signal 1 mm above the dipole axis as obtained from the model. It can be seen that the order of magnitude of these signals is similar. These results show the feasibility of detecting the surface flaws by means of sensor arrays embedded in the aluminium skin panels, in what could be a first step toward full self-diagnosis capabilities of the aircraft airframe as part of a structural health monitoring system.



**Fig. 6** Results from simulations with COMSOL: Contour plots of the variations of  $B_z$  within the cross-section of the sample and the dipole modelling the ferrofluid in the reference crack, in the  $xy$  plane, at 293 K, for applied field strength  $H$  of a) 16 kA/m and b) 716 kA/m (ferrofluid at  $M_{sat}$ )

**Table 8** For version Tech-DC.II at 293 K for the reference crack, variations of  $B_z$  obtained from the model and from the simulations with COMSOL for applied field strength  $H$  of 16 kA/m (test 1) and 716 kA/m (test 2) at 1.0 mm from the dipole axis.

Sample	Crack dimensions	Dipole	Numerical $B_z$ variation	Analytical $B_z$ variation
#1	$l = 60.00$ mm	$a = 30.00$ mm	test 1: $-3.1 \mu\text{T}$	test 1: $-6.9 \mu\text{T}$
	$w = 1.50$ mm	$R = 0.510$ mm	test 2: $-7.9 \mu\text{T}$	test 2: $-7.9 \mu\text{T}$
	$d = 0.70$ mm			

## 5. CONCLUSIONS

An innovative NDT technique is proposed for surface inspection of materials not necessarily magnetic or conductive, based on detection of local magnetic field variations due to ferrofluid deposited in the crack. A preliminary feasibility assessment is made, based on signal detectability without applied magnetic field, and under applied DC fields. For this purpose, the signals (local magnetic flux density variations) are quantified analytically, experimentally and numerically for cracks in plates of AA 2024. The main conclusions are:

- For the reference crack, filled with approximately  $50 \text{ mm}^3$  of ferrofluid, the magnetic field of the ferrofluid in absence of an applied field is below the sensor resolution. Detectable signals are obtained if the ferrofluid is being magnetized by an external field. The signals increase with the applied field strength  $H$ , e.g., reaching  $-7 \mu\text{T}$  at a distance of around 3.5 mm from the longitudinal axis of the reference crack, for  $H$  of 16 kA/m.
- The model agrees well with the tests: the average error excluding the results for a crack oriented perpendicular to the applied field is  $-5.3 \%$ , and both follow similar trends. For instance, the signal increases with  $H$  up to the saturation magnetization of the ferrofluid and, in a plane perpendicular to the crack longitudinal axis in the crack centre, decreases with the distance to the axis. Thus, it is concluded that the model can provide useful estimations of the signal.
- The signals measured for a crack oriented in the direction of the applied external field and perpendicular to it are virtually identical. This suggests that not knowing the direction of the defect when applying the external field may not be relevant to the performance of the proposed NDT method.
- The proposed NDT technique, requiring application of DC external fields to magnetize the ferrofluid to enhance the signal, seems promising: the model suggests that signals associated to cracks significantly smaller than surface cracks in a target application like aircraft skin panel inspection NASA STD-5009 are easily detectable with commercial magnetometers.

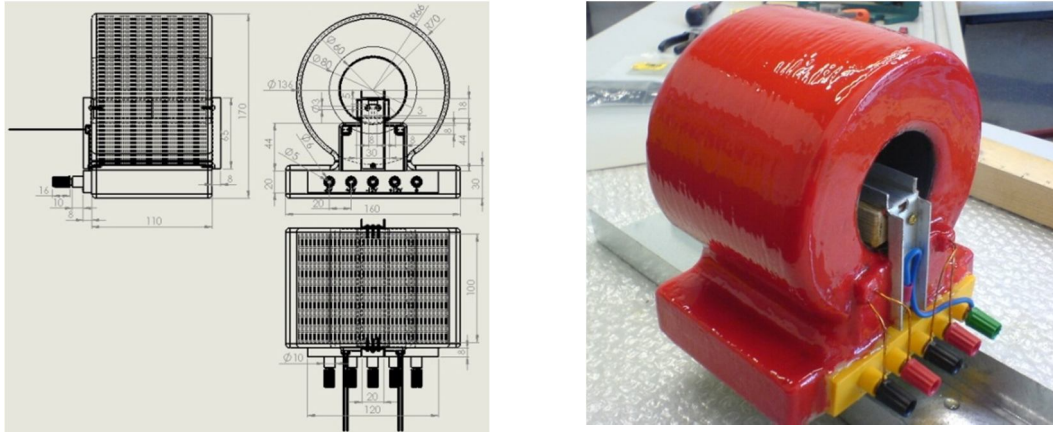


- Compared to PT, an advantage of the proposed NDT method is that it is quantitative and, therefore, can be used to estimate the size of the cracks.
- The use of non-commercial ferrofluids may provide some advantages like higher  $M_{sat}$ , leading to more intense signals (this can be achieved using  $\alpha$ -Fe particles), or lower coercivity/hysteresis effects, or the possibility of developing ferrofluids with tailored properties, like reduced viscosity and surface tension, making easier for the ferrofluid to sip into the surface cracks, etc.

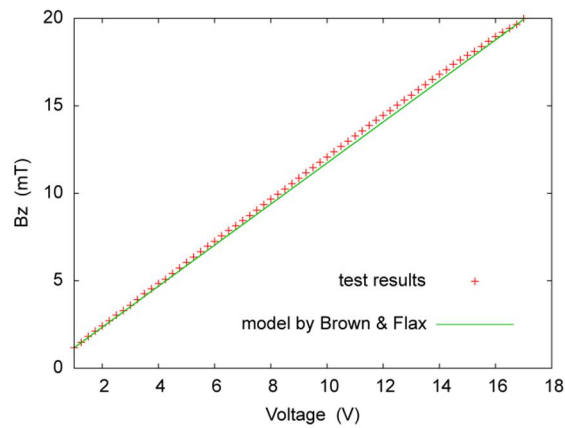
The ideas being considered for future work are: 1) to refine the research using, for instance, a 3-components Hall probe with higher sensitivity; 2) to correlate patterns in the local magnetic field variations with crack morphology; 3) to study the applicability of the technique to detect cracks in magnetic materials; 4) to study the effect of ferrofluid viscosity in crack penetration; 5) to study the feasibility of recycling classic eddy current equipment for implementing the proposed NDT technique; 6) perform PT, ET, MT (if possible, depending on the tested material) and tests with the proposed NDT method on small defects in the sub-mm range, to establish the limits of the NDT technique; and finally 7) to study the performance of the proposed technique upon application of AC fields. In AC, crack detection could be based on the phase lag between the field close to the crack and the applied field. This approach has inherent advantages: the phase lag, as opposed to  $\bar{B}$ , is independent of  $H$  and the quantity of ferrofluid in the crack, and increases significantly with the frequency of the applied AC field.

#### **APPENDIX A – Technical specifications of the custom-made bracket**

Coil #1 and coil #2 in the custom-made bracket are radially thick, multi-layered solenoids consisting of 1000 and 2800 turns, respectively, of solid Cu wire, 1 mm in diameter. Both coils are 100 mm long. Coil #1 has 30 (40) mm inner (outer) radius, while coil #2 has 40 (68) mm inner (outer) radius. Fig. 7 shows a photograph of the bracket and its lateral, frontal and top views, created with the commercial multiphysics software SolidWorks, from Dassault Systèmes SolidWorks Corp., Waltham, MA (USA). The DC magnetic field created by the coils when supplied with DC current can be estimated using a model by Brown and Flax [30]. In Fig. 8, the results from this model are compared with measurements using the AlphaLab magnetometer, for supplied DC voltages ranging from 1 to 17 V.



**Fig. 7** Lateral, frontal and top views created using commercial multiphysics software SolidWorks, from Dassault Systèmes SolidWorks Corp., Waltham, MA (USA) (left), and photograph of the custom-made bracket (right)



**Fig. 8**  $B_z$  for the DC field created by the coils vs. supplied DC voltage

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