

# High efficiency black silicon Interdigitated Back Contacted solar cells on p- and n-type c-Si substrates

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## Abstract

This work demonstrates the high potential of  $\text{Al}_2\text{O}_3$  passivated black silicon in high efficiency Interdigitated Back Contacted (IBC) solar cells by reducing surface reflectance without jeopardizing surface passivation. Very low reflectance values, below 0.7% in the 300-1000 nm wavelength range, together with striking surface recombination velocities values of 17 and 5 cm/s on p- and n-type crystalline silicon substrates respectively, are reached. The simultaneous fulfillment of both requirements, low reflectance and low surface recombination, paves the way for the fabrication of high efficiency interdigitated back contacted Si solar cells using black silicon at their front surface. Outstanding photovoltaic efficiencies over 22% have been achieved both in p- and n-type, 9  $\text{cm}^2$  cells. 3D simulations suggest that efficiencies up to 24% can be obtained in the future with minor modifications in the baseline fabrication process.

**Keywords:** black silicon, crystalline silicon, IBC solar cell, high efficiency, ALD  $\text{Al}_2\text{O}_3$ , passivation

## 1. Introduction

Interdigitated Back Contacted (IBC) crystalline Silicon solar cells allow excellent photovoltaic conversion efficiencies [1-3], in particular a new world record efficiency of 25.6% in large area n-type device has been recently obtained [4]. In IBC cells, both base and emitter regions with their corresponding electrodes are alternated in a strip-like fashion at the backside of the cell, whereas light impinges on the front device surface. Some advantages of this structure are evident, the absence of metal shadow losses and the facility to assemble solar cells in modules with coplanar connection. However, if one wants to exploit the full benefit of this structure, high minority-carrier bulk lifetime, outstanding front and rear surface passivation [5, 6], and low front reflectance are mandatory.

Very low reflectance values in the whole solar spectrum wavelength range can be achieved by nanostructuring front cell surface through so-called black silicon etching (b-Si) [7]. Another advantage of b-Si is to maintain excellent low reflectance values for a broad range of light incidence angles, which offers a potential for daily and yearly photovoltaic energy production increase [8]. However, as a result of black silicon formation the surface area greatly increases and consequently the recombination at the surface becomes a critical issue. It is well known that  $\text{Al}_2\text{O}_3$  films achieve outstanding passivation on planar and textured n- and p-type c-Si due to both chemical and field effect passivation [9, 10]. Regarding surface passivation of b-Si, very promising results using atomic layer deposited ALD  $\text{Al}_2\text{O}_3$  have been reported recently on both types of substrate polarities [11-13]. Additionally, ALD  $\text{Al}_2\text{O}_3$  material is also a high optical bandgap material (~6.5 eV) [10] with scarcely light absorption in the whole solar spectrum wavelength range [14], paving the way for applying nanostructures to the front surface of the IBC cell architecture (b-Si IBC cell from now on).

In a recent work, we have demonstrated high efficiency b-Si IBC solar cells (>22%) on p-type FZ c-Si substrates ( $9 \text{ cm}^2$  cell area) using b-Si passivated with ALD  $\text{Al}_2\text{O}_3$  at the front surface [8]. However the negative charge within the  $\text{Al}_2\text{O}_3$  layer accumulates or inverts the silicon surface on p- and n-type material respectively; thus it remains unclear if the overall surface passivation of b-Si on n-type will reach the same excellent values

as on p-type b-Si surface in finished IBC's solar cells. In fact, no efficiencies over 20% have been yet reported for n-type b-Si IBC's [15-17].

It is especially relevant to explore n-type b-Si IBC cells because it is the material of choice of IBC's high efficiency solar cell manufacturers for a number of reasons: on one hand because, unlike p-type substrates, this material do not exhibit light induced degradation (LID) by the formation of boron-oxygen complexes [18]. On the other hand, n-type substrates are less sensitive to the degradation of charge carrier lifetime due to recombination-active metal impurities (e.g. iron contamination) exhibiting higher stable lifetimes than p-type c-Si [3].

In this work we extend our previous results on p-type b-Si IBC cells [8] to n-type substrates. Surface passivation of b-Si nanostructures and photovoltaic b-Si IBC cell parameters are directly compared and analyzed for both types of substrates. 3D simulations will help to detect the main efficiency loss mechanisms giving clues to improve both p- and n-type b-Si IBC cells.

## **2. Black silicon etching. Reflectance and surface passivation**

### **2.1 Experimental**

In order to assess the reflectance of b-Si and its surface passivation we use conventional randomly textured surfaces as a reference. High quality defect free crystalline silicon FZ <100> oriented  $2.7 \pm 0.2 \Omega\text{cm}$  4 inches wafers ( $285 \pm 15 \mu\text{m}$  thick) were used in the experiments.

Black silicon was etched through a cryogenic inductively coupled plasma reactive-ion etching process (ICP-RIE) at -120 °C using SF<sub>6</sub> and O<sub>2</sub> as etching gases. Reference cells were textured by anisotropic etching in tetramethylammonium hydroxide (TMAH), isopropanol (IPA) and DI water (320 ml:350 ml:3300 ml) at 80°C resulting in random pyramids. After the etch, the wafers were RCA cleaned and HF-dipped before the atomic layer deposition of 20 or 90 nm thick Al<sub>2</sub>O<sub>3</sub> layers at 200°C for the b-Si and reference samples respectively. In the ALD process, TMA was used as the aluminum

source and water as the oxidant. Surface passivation has been monitored using symmetrically treated samples. Reflectance has been characterized on samples etched only on the front side with a back reflector scheme of thermally grown dry SiO<sub>2</sub> (110 nm) and aluminum Al (2 μm) deposited by e-beam. A final forming gas annealing at 400°C for 10 min was performed in all cases.

As can be seen in Fig. 1, the black silicon etching results in a morphology consisting of pillars densely and randomly distributed in the silicon surface with heights and base widths of less than 1 μm and 200 nm respectively. Only 1 μm of silicon was etched before the actual b-Si formation. Therefore this technique is very well suited to ultrathin crystalline Silicon solar cells (< 10 μm) [19] where conventional surface texturing is not possible due to relative high silicon consumption in the etching (~20 μm), as well as an alternative to acidic etching on mc-Si substrates [11].

<Fig.1>

## 2.2 Reflectance measurements

Reflectance measurements were made using a UV-VIS-NIR spectrometer (Shimadzu 3600) with integrating sphere. Fig. 2 shows the reflectance  $R$  of two samples, one of them b-Si etched and the other textured with random pyramids. The spectral photon flux density  $\phi_{AM1.5}$  corresponding to the AM1.5G solar spectrum is also depicted in the graph for comparison. Black silicon nanostructures exhibit a very low reflectance, below 0.7% in the 300-1000 nm wavelength range, clearly lower than the reference with random pyramids in the UV and blue part of the spectrum.

<Fig.2>

Using these reflectance results, one could estimate in both cases the maximum short-circuit current density  $J_{sc,max}$  of a hypothetical solar cell assuming zero transmittance, i.e ideal back reflector, and an ideal internal quantum efficiency as follows

$$J_{sc,max} \equiv q \int_{300 \text{ nm}}^{1130 \text{ nm}} \phi_{AM1.5}(\lambda) \times (1 - R(\lambda)) d\lambda \quad (1)$$

where  $q$  is the elementary charge and the integral is evaluated in the useful absorption wavelength range in c-Si under AM1.5G solar spectrum, i.e. 300-1130 nm. For b-Si surface  $J_{sc,max}$  reaches  $43.5 \text{ mA/cm}^2$ ,  $0.5 \text{ mA/cm}^2$  more than its texturized counterpart ( $43.0 \text{ mA/cm}^2$ ), or  $1 \text{ mA/cm}^2$  less compared with an ideal case of  $44.5 \text{ mA/cm}^2$  assuming a perfect antireflection surface, i.e.  $R=0\%$  in the whole studied wavelength range.

### 2.3 Surface passivation

In order to determine the surface passivation quality of our samples we performed minority-carrier lifetime measurements as a function of the excess carrier density  $\Delta n$  through the quasi-steady-state photoconductance (QSS-PC) decay method using the WCT-120 apparatus (Sinton Instruments). Effective surface recombination velocity  $S_{eff}$  is a good indicator of surface passivation. This parameter can be calculated easily from lifetime measurements in symmetrical samples using

$$\frac{1}{\tau_{eff}} - \frac{1}{\tau_{in}} \cong \frac{2S_{eff}}{w} \quad (2)$$

where  $\tau_{eff}$  is the measured effective lifetime,  $w$  the current wafer thickness, and  $\tau_{in}$  is the intrinsic bulk lifetime (due to Auger and radiative recombination mechanisms) according to the parameterization by Richter et al. [20]. Trap assisted Shockley-Read-Hall SRH bulk recombination mechanism can be neglected in the study due to the high quality of the wafers.

As it can be seen in fig. 3a, reference samples achieve lifetimes over 3 ms at low injection levels independently of the doping polarity indicating a very efficient surface passivation with  $S_{\text{eff}}$  values below 5 cm/s at  $\Delta n=10^{15} \text{ cm}^{-3}$  (see Fig. 3b). Surprisingly b-Si samples achieve also very low  $S_{\text{eff}}$  values, well below 20 cm/s (<5 cm/s for n-type substrate), despite the very unfavorable aspect ratio of the nanostructures compared with random pyramids resulting in a greatly increased total surface value.

<**Fig.3**>

Table I summarizes  $S_{\text{eff}}$  results in all analyzed samples. Especially remarkable surface passivation is obtained in the b-Si n-type sample yielding a  $S_{\text{eff}}$  of 4.8 cm/s, practically the same as the reference.

<**Table I**>

Reflectance and surface passivation results achieved in our precursors demonstrate the high potential of b-Si nanostructures, especially in n-type substrates, to be used for high efficiency solar cell fabrication. Regarding IBC solar cells, front surface  $S_{\text{eff}}$  values as those shown in table I would allow to obtain photovoltaic efficiencies up to 24.0% using a conventional fabrication baseline process as it will be discussed in section 4.

### 3. IBC solar cell fabrication details

Black silicon and reference random pyramids textured IBC solar cells, were fabricated using p- and n-type <100> FZ c-Si  $275\pm10 \mu\text{m}$  thick 4 inches wafers with  $2\pm0.5 \Omega\text{cm}$  resistivity. Cross sections of p- and n-type b-Si IBC cells are shown in Fig. 4a and 4b respectively.

<Fig.4>

Fabrication of IBC devices consist of high-temperature steps (dopant diffusions and thermal oxidations), photolithography and front surface passivation with ALD Al<sub>2</sub>O<sub>3</sub> on b-Si or textured surfaces. Phosphorus and boron diffusions are the same independently of the substrate polarity, consequently base and emitter areas are defined by interchanging lithography masks.

The fabrication process flow for b-Si IBC cells is pointed out in fig. 5a, following the next main stages, namely: i) Boron diffusion at 1000°C soak temperature (both sides) resulting in a sheet resistance  $R_{sh}$  of ~25 Ω/sq. ii) Isotropic etching (~2 μm etching depth) to define p+ base or emitter regions on p- or n-type substrates respectively. iii) Dry oxidation to grow a thermal SiO<sub>2</sub> to mask subsequent phosphorus diffusions. iv) n++ phosphorus diffusion at 900°C ( $R_{sh}$  ~25 Ω/sq) to define a selective emitter in p-type substrates or a high doped base contact in n-type wafers. v) Light n+ phosphorus diffusion at 830°C ( $R_{sh}$  ~100 Ω/sq). vi) Dry thermal oxidation at 1080°C to grow 110 nm thick SiO<sub>2</sub> to be used as a back reflector layer. Additionally, this film serves to mask anisotropic etching of reference textured IBC cells. Sheet resistances at the end of the drive-in stage were about 10, 70 and 100 Ω/sq in n++, n+ and p+ doped regions respectively. vii) and viii) b-Si etching and ALD Al<sub>2</sub>O<sub>3</sub> (20 nm thick) as explained in section 2.1. Reference IBC cells were textured with random pyramids using anisotropic etching. ix) After opening contact windows (250 μm contact pitch) at the middle of each base and emitter finger, Aluminum was evaporated by e-beam both at the rear (~2 μm) and front sides (~0.5 μm). This last metal layer clearly defines the active area (see Fig. 5c). Metal is patterned using standard photolithography to define both interdigitated electrodes at the backside (see Fig. 3b) and light windows at the front surface. x) Finally a forming gas annealing at T=400°C during 10 min is performed to recover e-beam damage [21] and activate Al<sub>2</sub>O<sub>3</sub> surface passivation [22].

<Fig.5>

#### 4. Results and discussion

Four wafers were processed, two p-type and two n-type, two using b-Si and two randomly textured reference IBC cells. In each wafer four  $3\text{ cm} \times 3\text{ cm}$  topologies were designed with different emitter coverage  $f_e$ , (i.e. emitter to total active area ratio), covering a  $f_e$  range from 67% to 86%. Base and emitter busbars were placed outside active area in all devices.

Solar cells were measured using a solar simulator (ORIEL 94021A) and a 4142B (Agilent HP) I-V characteristics tracer under standard test conditions STC (AM1.5G 1  $\text{kW/m}^2$  solar spectrum  $T=25^\circ\text{C}$ ). External quantum efficiencies curves were obtained with QX10 instrument using a  $2\text{ cm} \times 2\text{ cm}$  light beam and 0.1 suns bias light. The best b-Si cell and its reference counterpart were also measured ( $I$ - $V$  and  $EQE$  curves) in an independent laboratory (Fraunhofer CalLab).

Measured photovoltaic parameters as a function of  $f_e$  for the b-Si and reference IBC solar cells are shown in Fig. 6. 3D simulation results using Sentaurus Device software (Synopsys® TCAD tool) of b-Si structures are also included in the study (see Appendix section to simulation details) considering both an aluminum thickness of  $d_{\text{Al}}=2\text{ }\mu\text{m}$  in the backside metallization grid and the ideal case neglecting ohmic losses in fingers and busbars, i.e  $d_{\text{Al}} \rightarrow \infty$ .

<Fig.6>

As can be seen in Fig 6a and Fig. 6b, if  $f_e$  is tuned adequately, we reach efficiencies up to 22.1% ( $f_e=80\%$ ) and 22.0% ( $f_e=75\%$ ) in p- and n-type fabricated b-Si cells respectively. Table II summarizes photovoltaic parameters of the best b-Si cells, and its reference counterparts.

<Table II>

Excellent experimental short current density  $J_{sc}$  values about 41 mA/cm<sup>2</sup>, in both substrate polarities, indicate that a good front surface passivation is achieved in our nanostructures in agreement with high lifetime results of our test samples (see section 2). 3D simulations -see Fig. 6 and results enclosed in parenthesis in table II- envisage efficiencies up to 24.0%, with  $J_{sc}$ , open circuit voltages,  $V_{oc}$ 's, and fill factors,  $FF$ 's, of  $42.0 \pm 0.5$  mA/cm<sup>2</sup>,  $675 \pm 15$  mV and  $83.0 \pm 1\%$  ( $d_{Al} \rightarrow \infty$ ) respectively, independent of base doping. Experimental results are very close to the simulated values. However,  $FF$  is the parameter we might improve in future runs by increasing aluminum thickness  $d_{Al}$  or using a two level metallization scheme, i.e.  $d_{Al} \rightarrow \infty$ . Specific series resistance due to fingers and busbars is estimated to be about 0.6-0.8 Ωcm<sup>2</sup> in our current devices. By doubling Al thickness we can achieve specific resistances well below 0.4 Ωcm<sup>2</sup> increasing  $FF$  beyond 82%. Additionally,  $V_{oc}$ 's values in n-type IBC cells are clearly lower than in p-type cells possibly due to technological issues leading to poorly passivated surface at p+ boron regions (emitter or back surface field regions on n- and p-type devices respectively). This fact might explain the general trend observed in the  $V_{oc}$  curves in Fig 6c,d (experimental and simulated results), i.e. the larger relative size of the p+ region the lower  $V_{oc}$  value. Hence, best fitting between simulations and experimental results in the  $f_e$  50-90% range requires higher surface recombination at p+ than at n+ surfaces ( $S_{pas,p+} = 5 \times 10^4$  cm/s compared with  $S_{pas,n+} 1 \times 10^3$  cm/s). As a consequence optimum emitter coverage  $f_e$  has to be lower in n-type devices than in p-type ones. In fact, simulations foresee  $f_e < 60\%$  as the best choice when n-type substrates are used.

Measured external quantum efficiency  $EQE$  curves for the best b-Si cells on n- and p-type IBC devices with the corresponding reference counterparts are shown in Fig. 7. Simulated  $EQE$  results for b-Si cells using the front surface recombination velocity values measured on our test samples (see table I) are also included (solid lines) in the study.

<Fig.7>

Agreement between simulated and measured  $EQE$  values indicates that the excellent front surface passivation is preserved in our finished devices, leading to an  $EQE$  over 95% in the 400-1000 nm  $\lambda$  range in all cases in agreement with  $J_{sc}$  electrical results. Experimental  $EQE$  curves in b-Si and reference n-type devices are quite similar at  $\lambda \sim 600$  nm (see Fig. 7b), confirming that the b-Si surface reaches the same low level of surface recombination as textured one.

Experimental  $J_{sc}$  values differ from the theoretical  $J_{sc,max} \sim 43.5$  mA/cm<sup>2</sup> (see section 2.2) by the slight drop about 5% in  $EQE$  in the 300-1000 nm wavelength range compared with an ideal  $EQE=100\%$ . Notice that every 1% drop in  $EQE$  of our b-Si IBC cells decreases  $J_{sc}$  about 0.4 mA/cm<sup>2</sup>. The drop in photocurrent is partly related to internal quantum efficiency,  $IQE$ , losses by recombination either at the front surface or at non-doped and doped base (back surface field BSF) regions at the backside. However an unexpected high drop of  $J_{sc}$  occurs for  $f_e > 75\%$  in our n-type devices, as compared to the general trend of simulated  $J_{sc}$  values with  $f_e$  (see inset of Fig 6d). In order to clarify the origin of these  $J_{sc}$  losses, we have studied the effect of emitter coverage on  $EQE$  curves (see inset in Fig. 7b). Surprisingly in n-type cells  $EQE$  decreases with increasing  $f_e$  while the contrary occurs with much less overall variation in p-type devices. This anomalous behavior of our n-type devices might be related to lattice damage during boron diffusion [23, 24] resulting in a lower lifetime zone surrounding p+ regions.

## 5. Conclusions

This work demonstrates the viability of applying black silicon etching together with ALD Al<sub>2</sub>O<sub>3</sub> surface passivation in both n- and p-type IBC solar cells. Very low reflectance values below 0.7% in the 300-1000 nm wavelength range with striking surface recombination velocities about 17 and 5 cm/s on p- and n-type substrates, respectively, are achieved. The last result, very similar value to the one reached on a conventional surface with random pyramids passivated with 90 nm of ALD Al<sub>2</sub>O<sub>3</sub>, confirms the high potential of b-Si on high efficiency n-type IBC solar cells. Outstanding photovoltaic efficiencies around 22% are achieved in the fabricated b-Si IBC cells independently of the substrate polarity. Even higher efficiencies up to 24.0% are expected in following runs by reducing series resistance of the rear metallization

grid and improving both surface passivation and bulk quality of our diffused boron p+ regions.

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## Appendix. b-Si IBC solar cell 3D simulation details

Black silicon b-Si IBC solar cells have been simulated as a repetition of a simple domain, so-called unit cell, reducing simulation time and hardware computational requirements. In fig. 8 the simulated unit cell is shown. Notice that edge effects are neglected in the simulations.

<Fig.8>

The simulated unit cell is simplified as follows: i) all contacts are square shaped with a width  $2 \times r_c = 60 \mu\text{m}$  and placed at the middle of base and emitter fingers with a pitch,  $p$ , of  $250 \mu\text{m}$ - resulting in a locally contacted structure. ii) The wafer thickness  $w$  was set to  $w=275 \mu\text{m}$ . iii)  $S_F$  and  $S_R$  are the effective surface recombination velocities at the front and rear non-doped surfaces respectively.  $S_F$  value was adjusted accordingly to the values of table I, whereas  $S_R$  was always set to  $10 \text{ cm/s}$ . iv) p+ boron and n+ phosphorus diffusions are defined in a strip-like fashion with widths  $L_{p+}$  and  $L_{n+}$  respectively with a gap between regions of  $gap=40 \mu\text{m}$ . Additionally, we include a square n++ phosphorus doped region surrounding the corresponding contact width  $2 \times L_{n++}=120 \mu\text{m}$ . By varying

$L_{p+}$  or  $L_{n+}$  from a minimum value of 105  $\mu\text{m}$ , we can adjust emitter coverage  $f_e$  on n-type or p-type substrates respectively. Gaussian doping profiles are considered for the n+, n++ and p+ doped regions with junction depths of 1.4, 2.5 and 5  $\mu\text{m}$  respectively. Additionally, the peak doping  $N_s$  for n+ and p+ diffusions was set  $1 \times 10^{19} \text{ cm}^{-3}$  in both cases, and  $7 \times 10^{19} \text{ cm}^{-3}$  for the n++ region. Doping profile parameters are consistent with simulations of our baseline fabrication process using SUPREM (Stanford University Process Engineering Modeling) software. v) Surface recombination velocity of n+, n++ at non-contacted regions  $S_{\text{pas},n+}$  and  $S_{\text{pas},n++}$  were adjusted to  $1 \times 10^3$  and  $7 \times 10^3 \text{ cm/s}$  respectively accordingly Cuevas et al. parameterization [25].  $S_{\text{pas},p+}$  was set to  $5 \times 10^4 \text{ cm/s}$ . This last value is representative of a boron p+ surface with a peak doping of  $\sim 10^{19} \text{ cm}^{-3}$  and passivated with an annealed thermal SiO<sub>2</sub> grown in chloride atmosphere once surface passivation is stabilized in time [26]. vi) Surface recombination velocity at the contacts was set to the thermal velocity limit ( $S_{c,p+}=S_{c,n++}=5 \times 10^6 \text{ cm/s}$ ) as boundary condition. vii) Bulk recombination is limited by intrinsic mechanisms (Auger and radiative recombination), which is modelled using the parameterization given in [19]. Additionally trap-assisted Shockley-Read-Hall SRH recombination is neglected in the study (the related lifetime is set to 1 s). viii) The nanostructured front surface is replaced by a flat polished surface with total optical transmission ( $T=100\%$ ) and considering simultaneously a modified AM1.5G spectral irradiance  $S_{\text{mod}}=S(\lambda) \times (1-R(\lambda))$  as incident light. Here,  $S(\lambda)$  is the spectral irradiance of the standard AM1.5G solar spectrum and  $R(\lambda)$  is the measured reflectance of the b-Si IBC precursor shown in Fig. 2. Additionally, an internal front and rear reflectance of  $R_{\text{int\_front}}=100\%$  and  $\langle R_{\text{int\_rear}} \rangle$  respectively are introduced in the simulations. The last parameter is adjusted for each emitter coverage taking into account the metallization gap between base and emitter electrodes  $gap_{\text{met}}=80 \mu\text{m}$ , an internal reflection of 96% for metallized non-contacted areas, and 85% for contacted and non-metallized areas. The former value corresponds to an ideal back reflector using an optimized Al/SiO<sub>2</sub> stack. The last value arises from the best fitting between simulated and measured EQE curves in the 1050-1150 nm wavelength range for all emitter coverage  $f_e$  range (60-86%). ix) Flat bands have been considered at the passivated non-doped semiconductor surfaces, i.e. neither dielectric charge nor metal-insulator-semiconductor interface are considered in the simulations. Therefore, there are no inversion or accumulation layers below passivated surfaces. x) Physical models - e.g. carrier statistics, band gap narrowing, and carrier mobilities - are

the same as considered in [27]. xi) Ohmic losses related to the backside metallization grid (fingers and busbars) are modelled by a lumped series resistance  $R_s$  in  $\Omega$  units whose value is:

$$R_s \cong \frac{R_{sh,Al}}{3} \cdot \left( \frac{1}{f_{m,b}} + \frac{1}{f_{m,e}} + \frac{1}{f_{m,bus}} \right) \cdot \frac{A}{A_i} \quad (3)$$

where  $A$  and  $A_i$  is the total cell and simulated unit cell area respectively,  $f_{m,b}$ ,  $f_{m,e}$ ,  $f_{m,bus}$  are the metallization factors due to base and emitter fingers and busbars respectively.  $R_{sh,Al}$  is the metal sheet resistance which is evaluated taking into account the aluminum resistivity  $\rho_{Al} 2.8 \times 10^{-6} \Omega\text{cm}$  and the metal layer thickness  $d_{Al}$  as  $R_{sh,Al} = \rho_{Al}/d_{Al}$ . Notice that  $f_{m,b}$  and  $f_{m,e}$  changes accordingly with emitter coverage considering current  $L_{p+}$ ,  $L_{n+}$  values and the gap between electrodes  $gap_{met}$  as follows:

$$f_{m,b} = \frac{L_{p+} - E_{gap,met}/2}{L_{p+} + L_{n+} + E_{gap,met}} \quad (4)$$

$$f_{m,e} = \frac{L_{n+} - E_{gap,met}/2}{L_{p+} + L_{n+} + E_{gap,met}} \quad (5)$$

$f_{m,bus}$  was always set to 0.133 independent of  $f_e$ . Additionally, the contact resistance between metal and high doped contacted regions is neglected in the simulations.

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**Table I.**  $S_{\text{eff}}$  values evaluated at  $\Delta n=10^{15} \text{ cm}^{-3}$  for both p- and n-type substrates considering random pyramids textured or black silicon morphologies.

|                         | <i>p-type</i><br>b-Si | <i>p-type</i><br>random | <i>n-type</i><br>b-Si | <i>n-type</i><br>random |
|-------------------------|-----------------------|-------------------------|-----------------------|-------------------------|
| $S_{\text{eff}}$ (cm/s) | 16.7                  | 3.2                     | 4.8                   | 4.8                     |

**Table II. Best results under STC conditions (AM1.5G 1 kW/m<sup>2</sup> solar spectrum T=25°C) for b-Si and their reference cells on n-type and p-type substrates (9 cm<sup>2</sup> cell area). Values enclosed in parenthesis correspond to simulated results with  $d_{Al} \rightarrow \infty$ .**

| Cell                     | $f_e$ (%) | $J_{sc}$ (mA/cm <sup>2</sup> ) | $V_{oc}$ (mV)    | FF (%)             | $\eta$ (%)         |
|--------------------------|-----------|--------------------------------|------------------|--------------------|--------------------|
| <b>b-Si (p-type) (*)</b> | <b>80</b> | <b>41.0 (41.4)</b>             | <b>678 (684)</b> | <b>79.5 (83.0)</b> | <b>22.1 (23.5)</b> |
| Ref. (p-type) (*)        | 80        | 41.4                           | 679              | 78.9               | 22.2               |
| <b>b-Si (n-type)</b>     | <b>75</b> | <b>40.9 (41.5)</b>             | <b>662 (661)</b> | <b>81.3 (83.9)</b> | <b>22.0 (23.0)</b> |
| Ref. (n-type)            | 75        | 40.3                           | 645              | 80.7               | 21.0               |

(\*) These cells were measured in an independent laboratory (Fraunhofer ISE CalLab).

## Figure Captions

**Figure 1.** a), c) FIB/SEM images of silicon surface after b-Si and random pyramids etching respectively. b), d) A close view showing b-Si and random pyramid surface morphologies respectively. Notice that fig a) and d) have the same scale.

**Figure 2.** Front reflectance measured in two p-type precursors with b-Si and random textured surfaces passivated with ALD  $\text{Al}_2\text{O}_3$ . Samples include a back reflector scheme of  $\text{SiO}_2$  (110 nm)/Al (2  $\mu\text{m}$ ). Spectral photon flux density considering AM1.5G solar spectrum  $\phi_{\text{AM}1.5}$  is also depicted in the graph.

**Figure 3.** a) Effective lifetime and b) effective surface recombination velocity measurements vs. excess carrier density  $\Delta n$  on b-Si (open symbols) and random (filled symbols) p-type and n-type substrates. p-type data are the same as reported in [8].

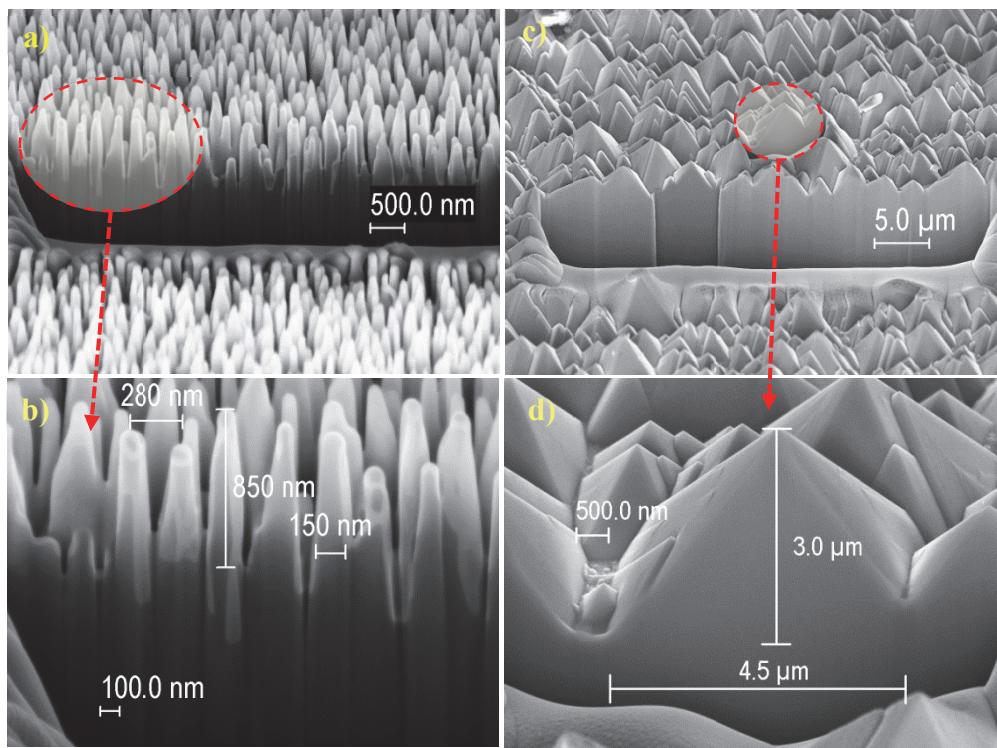
**Figure 4.** Cross section sketches of a) p- and b) n-type b-Si IBC solar cells.

**Figure 5.** a) Fabrication process flow for n- and p-type b-Si IBC cells. Images of a manufactured 4 inches wafer at the end of the fabrication process showing b) rear and c) front surfaces. Each wafer contains four  $9 \text{ cm}^2$  devices with different emitter coverage 67%, 75%, 80% and 86%.

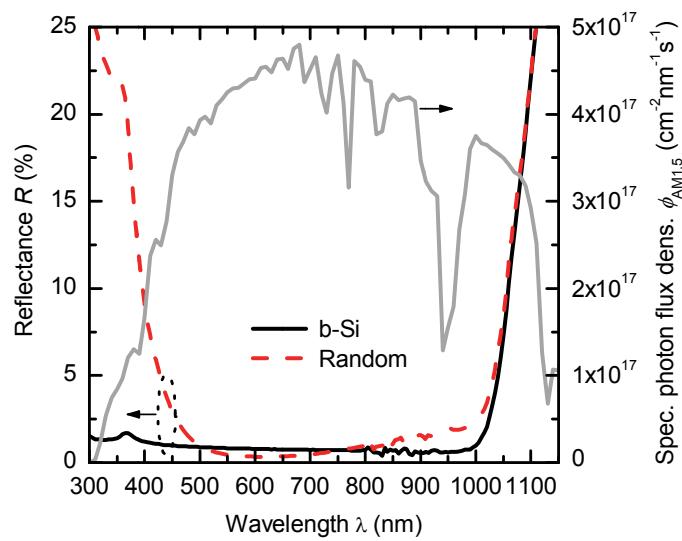
**Figure 6.** a), b) Measured values of efficiency  $\eta$ , c), d)  $V_{\text{oc}}$  and e), f)  $FF$  vs. emitter coverage  $f_e$  of b-Si (open symbols) and reference (filled symbols) IBC solar cells on p-(left) and n-type substrates (right). Dashed lines correspond to simulated efficiencies taking into account aluminum thickness of the backside metallization grid,  $d_{\text{Al}}$ . Short-circuit current density  $J_{\text{sc}}$  is also shown in the insets of e) and f). Note, that the b-Si p-type cell with  $f_e=75\%$  failed, therefore the corresponding values are missing in the graphs.

**Figure 7.** a) External quantum efficiencies  $EQE$ 's of best p-type b-Si IBC solar cell ( $f_e=80\%$ ) and its reference counterpart.  $EQE$  data of both curves were measured in an independent laboratory (Fraunhofer ISE CalLab). b)  $EQE$ 's of best n-type b-Si cell ( $f_e=75\%$ ) and its reference counterpart. Experimental  $EQE$  curves for several  $f_e$ 's are also shown in the insets. Solid lines in a) and b) correspond to 3D simulations of b-Si cells considering the front surface recombination velocities of our precursors of Table I.

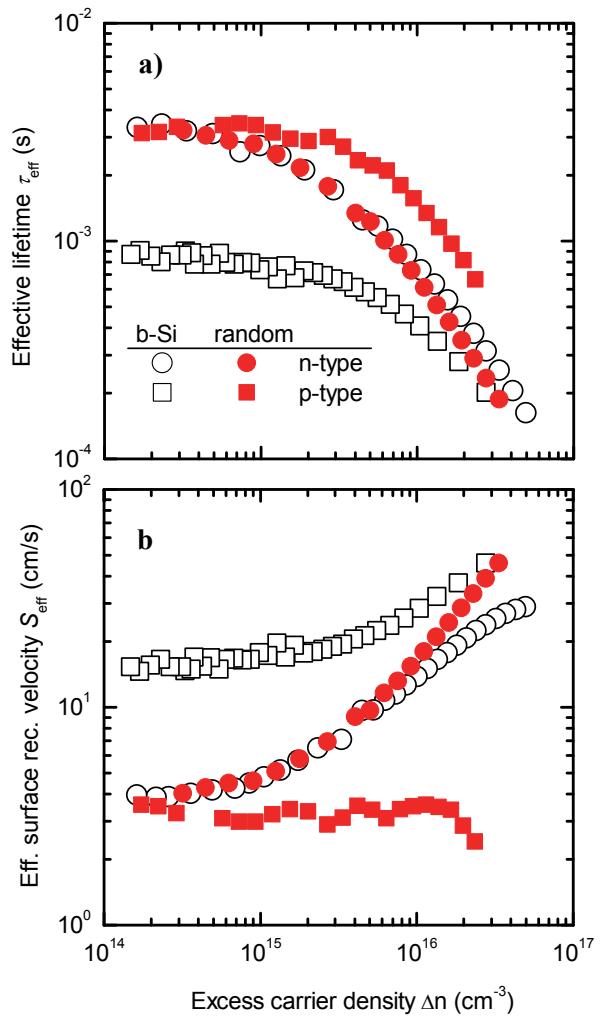
**Figure 8.** Sketch of the simulated structure: a) 2D cross section and b) bottom view.



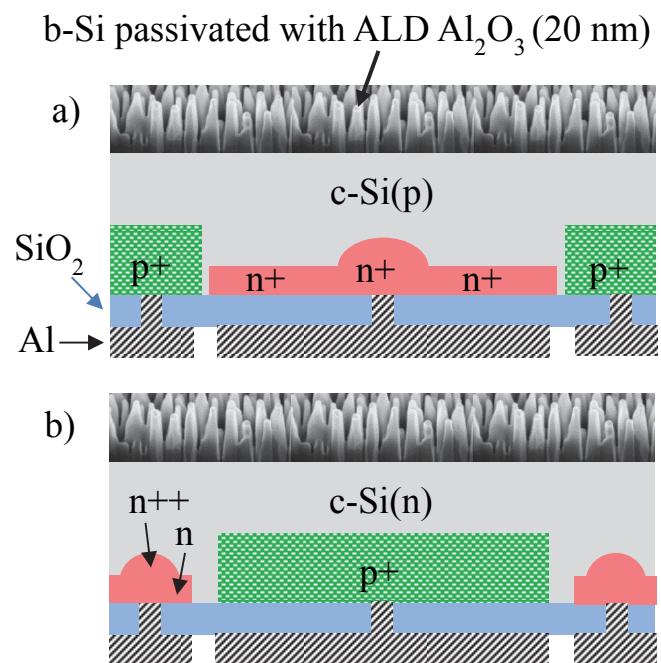
<Fig. 1>



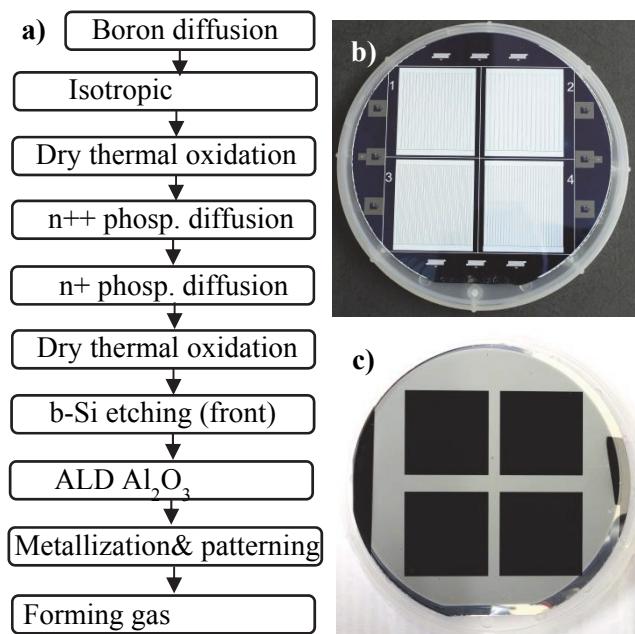
<Fig. 2>



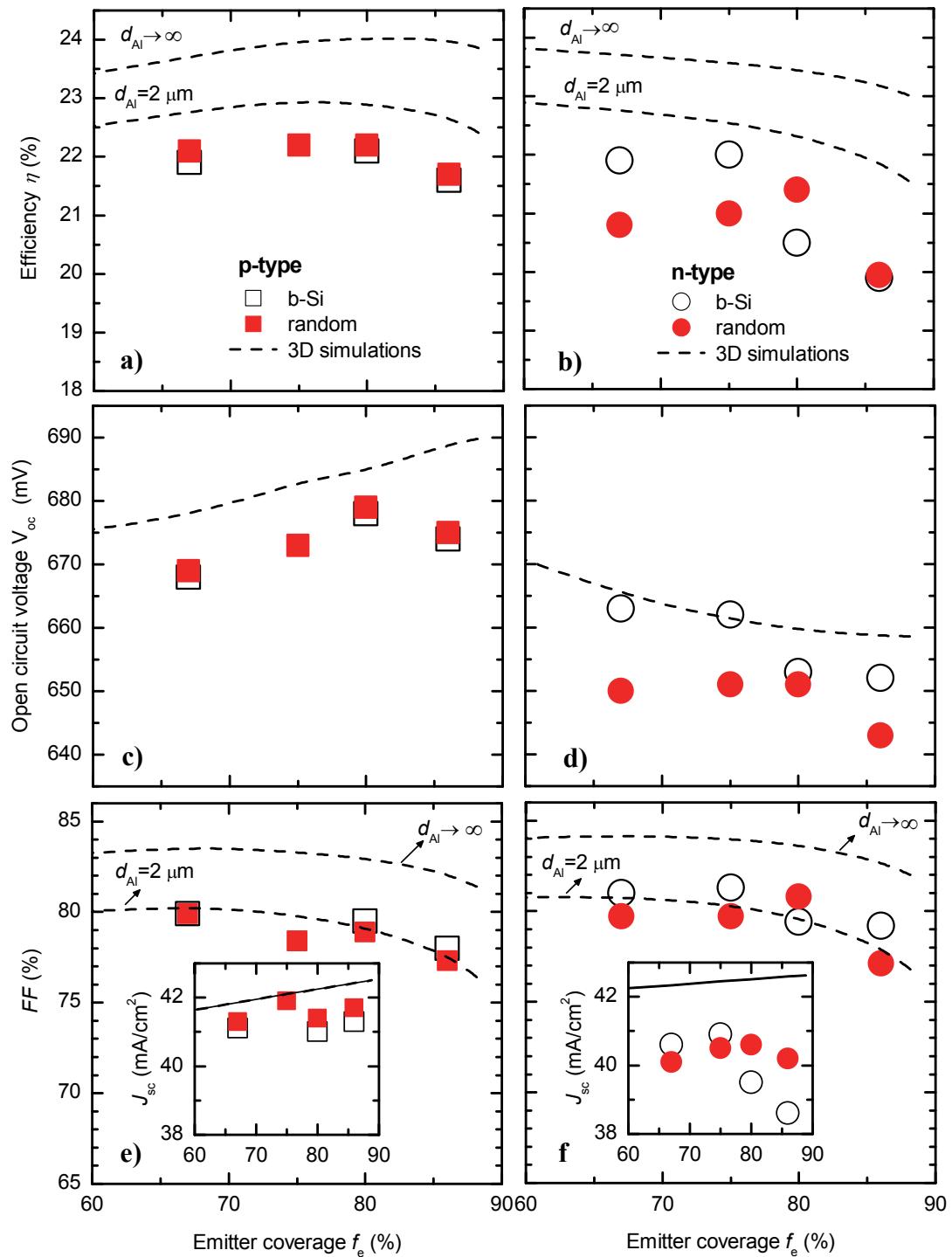
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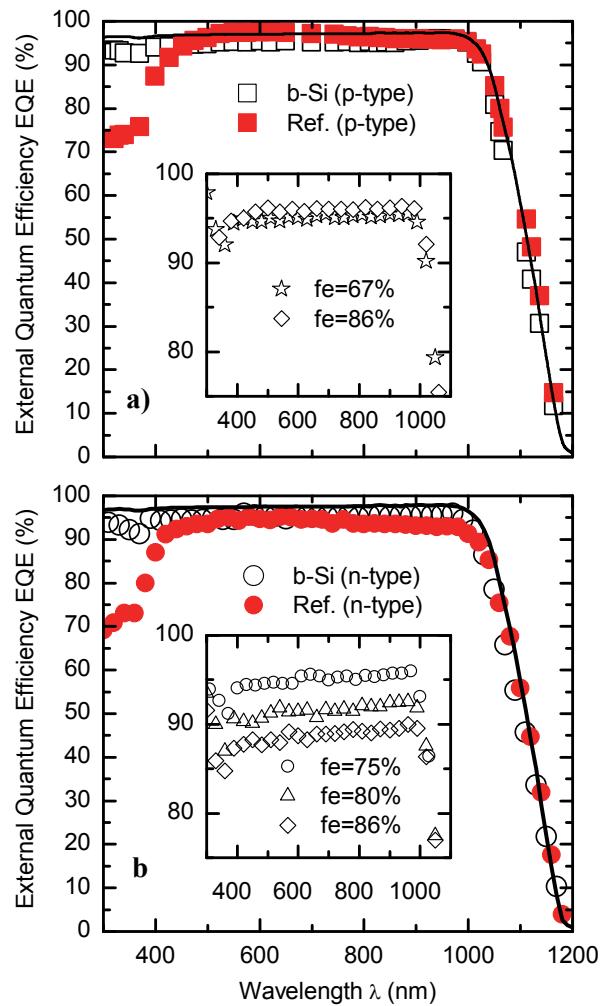
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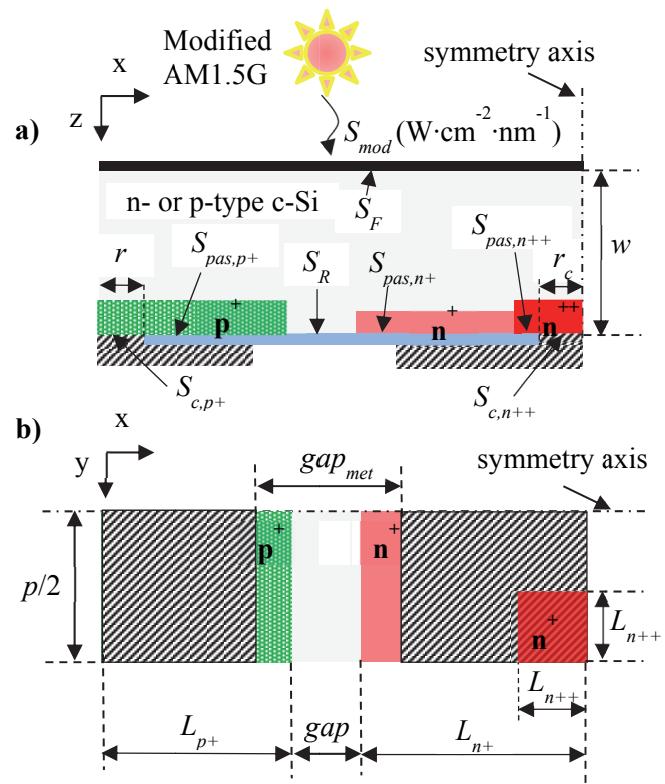
<Fig. 5>



<Fig. 6>



<Fig. 7>



■ Passivated b-Si ■ Rear isolation layer

▨ Base and emitter interdigitated electrodes

<Fig. 8>