

The White Dwarf Binary Pathways Survey IV: Three close white dwarf binaries with G-type secondary stars

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ABSTRACT

Constraints from surveys of post common envelope binaries (PCEBs) consisting of a white dwarf plus an M-dwarf companion have led to significant progress in our understanding of the formation of close white dwarf binary stars with low-mass companions. The white dwarf binary pathways project aims at extending these previous surveys to larger secondary masses, i.e. secondary stars of spectral type AFGK. Here we present the discovery and observational characterization of three PCEBs with G-type secondary stars and orbital periods between 1.2 and 2.5 days. Using our own tools as well as MESA we estimate the evolutionary history of the binary stars and predict their future. We find a large range of possible evolutionary histories for all three systems and identify no indications for differences in common envelope evolution compared to PCEBs with lower mass secondary stars. Despite their similarities in orbital period and secondary spectral type, we estimate that the future of the three systems are very different: TYC 4962-1205-1 is a progenitor of a cataclysmic variable system with an evolved donor star, TYC 4700-815-1 will run into dynamically unstable mass transfer that will cause the two stars to merge, and TYC 1380-957-1 may appear as super soft source before becoming a rather typical cataclysmic variable star.

Key words: binaries: close – stars: white dwarfs – stars: evolution – techniques: radial velocities

1 INTRODUCTION

Close white dwarf binary stars produce some of the most important incidents in modern astronomy. In particular, thermonuclear explosions such as Type Ia Supernovae (SN Ia) are one possible outcome of close white dwarf binary star evolution. For decades two main SN Ia formation channels were considered, the single (Webbink 1984) and double (Iben & Tutukov 1984) degenerate channels. Both

channels lead to thermonuclear explosions of white dwarfs close to the Chandrasekhar mass limit.

In recent years it has become evident that reaching a mass close to the Chandrasekhar limit might not be a strictly necessary criterion for SN Ia explosions, enlarging the possible evolutionary scenarios for single or double degenerate SN Ia detonations (Fink et al. 2007; Sim et al. 2010; Kromer et al. 2010; Brooks et al. 2016; Guillochon et al. 2010; van Kerkwijk et al. 2010; Pakmor et al. 2013). Furthermore, although classical SN Ia lying on the Phillips relation (Nugent et al. 1995) still make up the majority of observed thermonuclear SNe explosions, a variety of peculiar SNe related

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to the thermonuclear explosion of a white dwarf have been discovered recently. Among others, these atypical thermonuclear SNe include explosions with low ejecta velocities (e.g. SN Iax), possibly produced by accretion of a white dwarf from a He-star that may potentially leave a burned white dwarf remnant (Foley et al. 2013), and calcium rich SNe which might be long delay-time thermonuclear explosions of a white dwarf in a binary system that was dynamically ejected from its host galaxy (Foley 2015). The growing variety of potential SN Ia progenitor systems and the increasing number of different flavours of thermonuclear supernovae (see Jha et al. 2019, for a review), all resulting from close binaries that contain at least one white dwarf driven close to the Chandrasekhar mass limit, imply that a detailed understanding of white dwarf binary star evolution is required if we want to know what are the main progenitor systems of thermonuclear SN and how frequently a given SN Ia progenitor configuration is produced.

The standard hypothesis for close white dwarf binary formation is that the initial main sequence binaries, consisting of two stars with masses exceeding $1 M_{\odot}$, evolve through a common envelope phase (e.g. Webbink 1984; Zorotovic et al. 2010) generated by the initially more massive star filling its Roche-lobe as a giant star. The close binaries emerging from common envelope evolution consist of a white dwarf and a main sequence companion star with an orbital period usually shorter than a few days. The future evolution of these post common envelope binaries (PCEBs) depends essentially on their orbital period but also on the evolutionary status of the secondary star and the mass ratio. Broadly speaking, for periods longer than 1-2 days, a second phase of mass transfer will likely be initiated when the secondary has evolved into a giant star which will lead in most cases to a second common envelope phase and potentially to a double degenerate binary or a merger (e.g. Webbink 1984; Olling et al. 2015). For shorter periods, angular momentum loss can drive the system into mass transfer when the secondary star is still on the main sequence which, depending on the mass ratio, can generate thermal time scale mass transfer and stable nuclear burning on the surface of the white dwarf (Shen & Bildsten 2007). These systems are called Super Soft X-ray Sources (SSS) and are considered possible single degenerate SN Ia progenitors as the white dwarf can grow in mass (e.g. Di Stefano 2010).

Despite the general understanding of close white dwarf binary formation outlined above, the predictions of present day binary population models for close white dwarf binaries are very uncertain. Usually the outcome of common envelope evolution is determined using a simple parameterized energy conservation equation and, unfortunately, the key parameter, the common envelope efficiency, is so far only well constrained for low-mass ($\lesssim 0.5 M_{\odot}$) main sequence secondary stars (Zorotovic et al. 2010; Nebot Gómez-Morán et al. 2011). These low mass PCEBs are important for population studies of cataclysmic variables (CVs) but most likely not relevant for studying evolutionary channels towards SN Ia explosions because, as a consequence of nova eruptions, the white dwarfs in CVs do not significantly grow in mass (Yaron et al. 2005; Wijnen et al. 2015; Schreiber et al. 2016). Providing observational constraints on the formation and evolution of close white dwarf binaries with more massive secondary stars is therefore of utmost importance for a better understanding of binary pathways that may lead to thermonuclear SN.

We are conducting a survey of close binaries consisting of main sequence stars of spectral type AFGK plus a white dwarf (WD+AFGK from now on). Our target selection is based on optical surveys that are combined with the Galaxy Evolution Explorer DR6+7 photometry (GALEX; Bianchi et al. 2014), to identify ob-

jects with ultraviolet (UV) excess indicative for the presence of a white dwarf. Close binaries are then identified and characterized through radial velocity measurements. So far, we have identified the first pre-SSS binary system (Parsons et al. 2015), confirmed that our target selection is efficient (Parsons et al. 2016, hereafter paper I), obtained first measurements of the fractions of close binaries among our targets (Rebassa-Mansergas et al. 2017, hereafter paper II), and discussed in detail the limited contamination of our target sample through hierarchical triple systems with a white dwarf component (Lagos et al. 2020, hereafter paper III).

Here we present the identification and analysis of three close white dwarf binaries with early type secondary stars. We measured their orbital periods, estimated the stellar masses based on high resolution spectroscopy and used numerical codes to describe the history and future of the three systems. We found that these three close binaries have most likely formed through common envelope evolution and that their evolutionary history can be reconstructed with common envelope parameters similar to those used for PCEBs with M-dwarf companions. Concerning the future of the three systems we obtained that, despite having relatively similar stellar and binary parameters, their final fates are rather different.

2 OBSERVATIONS AND DATA ANALYSIS

White dwarf plus AFGK binaries are very difficult to be identified based on optical data alone because the main sequence star completely outshines any contribution from the white dwarf. Here we present a brief explanation of how we overcame this problem and describe the identification process of the three objects discussed in this work.

2.1 Target selection

Whereas we relied in Paper I and Paper II on the selection of WD+AFGK binary candidates from the RADial Velocity Experiment (RAVE; Kordopatis et al. 2013) and The Large Sky Area Multi-Object Fibre Spectroscopic Telescope (LAMOST; Deng et al. 2012) databases, in the present work, we have taken advantage of the *Tycho-Gaia* Astrometric Solution (TGAS; Lindegren et al. 2016) for the identification of suitable candidates in a magnitude limited sample.

For the purpose of characterising the excess emission with respect to normal UV colours of AFGK stars, we cross-matched the TGAS database with DR6+7 GALEX photometry, identifying ~ 13000 reliable matches with $FUV \leq 19$ and TGAS parallax uncertainty $\sigma_{\varpi}/\varpi \leq 0.15$. We complemented them with optical and infrared photometry from The AAVSO Photometric All-Sky Survey DR9 (APASS; Henden et al. 2015), Two Micron All-Sky Survey (2MASS; Skrutskie et al. 2006), and Wide-Field Infrared Survey Explorer (WISE; Wright et al. 2010). We determined photometric effective temperatures, T_{eff} , by fitting the optical-to-infrared spectral energy distributions (SED) with a grid of synthetic spectra (PHOENIX; Husser et al. 2013). The result of the SED fitting delivered ≈ 1000 WD+AFGK binary candidates, displaying moderate to significant UV-excess with respect to normal colours of main-sequence and giant stars of below-Solar metallicity ($Z = -1$). We later revised the target selection by also including parallaxes and broad-band photometry from the second data release of *Gaia* (DR2; Gaia Collaboration et al. 2018a), bringing the target list down to ≈ 800 candidates. The identification of suitable candidates is illustrated in Fig. 1. A more detailed description of the photometric

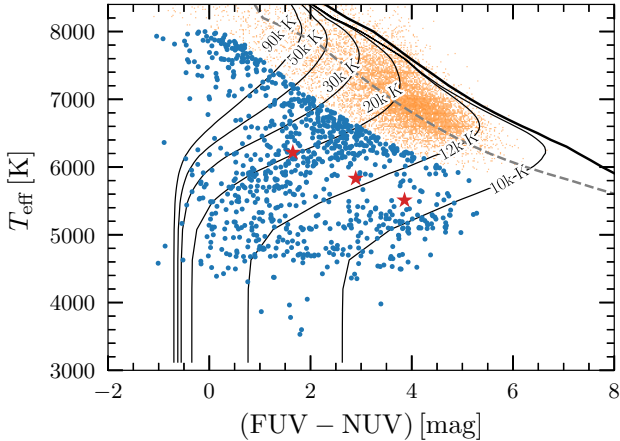


Figure 1. *GALEX* FUV – NUV colours vs photometric T_{eff} of $\sim 13\,000$ TGAS stars (orange density cloud). The WD+AFGK binary candidates are shown as blue circles, while the three binaries studied in this paper are represented by red star symbols. The synthetic colours of *PHOENIX* (Husser et al. 2013) models with $\log g = 4.5$ and $Z = 0$, or $\log g = 3.5$ and $Z = -1$, are displayed as solid black and dashed grey curves, respectively. The colours of WD+AFGK companions, computed for a representative set of white dwarf T_{eff} , are determined by combining $\log g = 8$, Koester (2010) models with *PHOENIX* models, which we scaled via the WD mass-radius relation (Fontaine et al. 2001) and the main-sequence masses and radii (Choi et al. 2016), respectively.

Table 1. *GALEX* magnitudes with APASS, 2MASS and WISE fluxes of the three objects discussed in this work

Filter	λ [Å]	TYC 4700-815-1	TYC 1380-957-1	TYC 4962-1205-1
FUV	1516	14.91±0.01	18.31±0.08	18.71±0.07
NUV	2267	13.26 ±0.00	15.45±0.01	14.51±0.01
V	5394.3	$9.36(0.06)\times 10^{-13}$	$2.06(0.09)\times 10^{-13}$	$9.69(0.04)\times 10^{-13}$
J	12350	$2.05(0.04)\times 10^{-13}$	$5.14(0.01)\times 10^{-14}$	$2.80(0.05)\times 10^{-13}$
H	16620	$9.25(0.04)\times 10^{-14}$	$2.36(0.67)\times 10^{-14}$	$1.52(0.06)\times 10^{-13}$
Ks	21590	$3.67(0.71)\times 10^{-14}$	$1.00(0.20)\times 10^{-14}$	$6.12(0.02)\times 10^{-14}$
W1	33526	$7.47(0.16)\times 10^{-15}$	$1.93(0.004)\times 10^{-15}$	$1.25(0.36)\times 10^{-14}$
W2	46028	$2.14(0.039)\times 10^{-15}$	$5.50(1.01)\times 10^{-16}$	$3.66(0.006)\times 10^{-15}$

V, J, H, Ks, W1, W2 fluxes are in $\text{erg cm}^{-2}\text{s}^{-1}\text{Å}^{-1}$

classification of the candidates and first spectroscopic follow-up observations of the sample, are presented in Ren et al. (2020).

Here we present follow-up observations and a detailed analysis of three systems with short orbital periods: TYC 4700-815-1, TYC 1380-957-1 and TYC 4962-1205-1. The *GALEX* magnitudes with APASS, 2MASS and WISE fluxes of these three systems are listed in Table 1. Hereafter, we will refer to the three systems as TYC 4700, TYC 1380 and TYC 4962 respectively.

2.2 High resolution spectroscopy

The high-resolution spectroscopic observations were performed using five different instruments installed at different telescopes. We used the Ultraviolet and Visual Echelle Spectrograph (UVES, Dekker et al. 2000) on the Very Large Telescope in Cerro Paranal and the Fiber-fed Extended Range Optical Spectrograph (FEROS, Kaufer et al. 1999) at the 2.2 m telescope of the Max Planck Gesellschaft (MPG) at La Silla Observatory, both in Chile. We furthermore obtained data with the High Resolution Spectrograph

(HRS, Jiang 2002) at Xinglong 2.16 m telescope (XL216) located in China, the Middle-Resolution fiber-fed Echelle Spectrograph (MRES) at the 2.4 m Thai National Telescope (TNT) in Thailand and the Echelle Spectrograph from REOSC for the Sierra San Pedro Martir Observatory (ESPRESSO, Levine & Chakarabarty 1995) at 2.12 m telescope of the Observatorio Astronómico Nacional at San Pedro Mártir (OAN-SPM)¹, México. Nearly all spectra were obtained during gray nights and relatively poor weather conditions. However, all with sufficient signal to noise ratios.

The cross-dispersed echelle spectrograph UVES has a spectral resolution of 58 000 for a 0.7 arcsec slit. The two-arms cover the wavelength range of 3000–5000 Å (Blue) and 4200–11 000 Å (Red), centered at 3900 and 5640 Å respectively. Standard reduction was performed making use of the specialized pipeline EsoReflux workflows (Freudling et al. 2013).

FEROS on the 2.2 m Telescope has a resolution of $R \approx 48\,000$ and covers the wavelength range from ~ 3500 – 9200 Å. Data obtained with the FEROS spectrograph were extracted and analysed with an automated pipeline developed to process spectra coming from different instruments in an homogeneous and robust manner, the CERES code (Jordán et al. 2014; Brahm et al. 2017). After performing typical image reductions, spectra were optimally extracted following Marsh (1989). The instrumental drift in wavelength through the night was corrected with a secondary fiber observing a Th-Ar lamp.

The HRS Spectrograph at Xinglong 2.16 m telescope provides echelle spectra of a 49 800 resolving power for a fixed slit width of 0.19 mm and covers a wavelength range of ~ 3650 – $10\,000$ Å. Th-Ar arc spectra were taken at the beginning and end of each night. Additional follow-up spectroscopy with MRES at the TNT was obtained using a slit width of 1.4 arcsec. MRES provides spectra of a resolving power of 15 000 covering the 3900–8800 Å wavelength range. Arc spectra were taken at the beginning of each night and we used sky lines to account for the flexure of the spectrograph.

The ESPRESSO Spectrograph of OAN-SPM provides spectra covering 3500–7105 Å with a spectral resolving power of $R \approx 18\,000$. A Th-Ar lamp before each exposure was used for wavelength calibration. The MRES, HRS and ESPRESSO spectra were reduced using the *echelle* package in IRAF². Standard procedures, including bias subtraction, cosmic-ray removal, and wavelength calibration were carried out using the corresponding tasks in IRAF.

2.3 Radial Velocity Measurements

Radial velocities (RVs) from UVES and ESPRESSO spectra were computed using the cross-correlation technique against a binary mask corresponding to a G-type star. The uncertainties in radial velocity were computed using scaling relations (Jordán et al. 2014) with the signal-to-noise ratio and width of the cross-correlation peak, which were calibrated with Monte Carlo simulations. FEROS spectra were analyzed with the CERES code which also calculates radial velocities using cross-correlation (for more details, see Jordán et al. 2014; Brahm et al. 2017). RVs from the spectra obtained with HRS and MRES telescopes were obtained by fitting the normalised CaII absorption triplet (at 8498.02, 8542.09, and 8662.14 Å) and

¹ <http://www.astrossp.unam.mx>

² IRAF is distributed by the National Optical Astronomy Observatories, which are operated by the Association of Universities for Research in Astronomy, Inc., under cooperative agreement with the National Science Foundation.

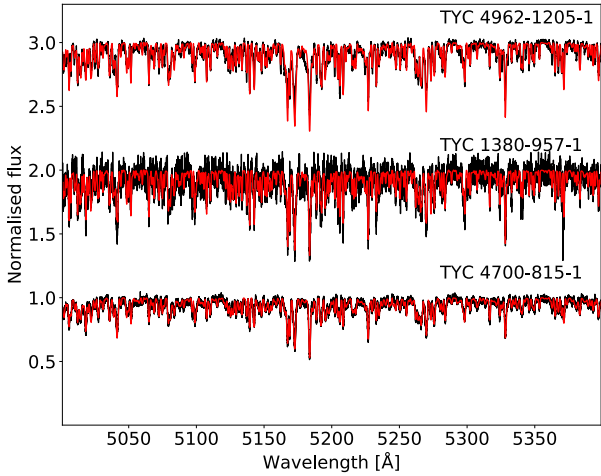


Figure 2. A small range of the observed high resolution spectra (black) from UVES for TYC 4700 and FEROS for TYC 1380 and TYC 4962. Spectral fits to the three G type stars using synthetic spectra within iSpec (red). The best-fit results were combined with SED fitting and *Gaia* DR2 parallaxes to obtain our final parameter estimates.

Nai respectively, with a combination of a second order polynomial and a triple-Gaussian profile of fixed separations, as described in Paper II. Only when the CaII absorption triplet was too noisy to get a reliable RV, the Nai doublet at ~ 5889.95 and 5895.92 Å was used. In the latter case we used a second order polynomial and a double-Gaussian profile of fixed separation. The RV uncertainty is obtained by summing the fitted error and a systematic error of 0.5 km s^{-1} in quadrature (more details in Ren et al. (2020)).

All RVs obtained for the three objects are provided in the appendix in the tables A1, A2, A3.

3 STELLAR AND BINARY PARAMETERS

3.1 Characterizing the secondary star

We determined the parameters of the main-sequence stars in our binaries by fitting our high-resolution spectra with model spectra and combining the obtained results with a fit to their spectral energy distributions (SEDs) scaled to *Gaia* DR2 parallaxes. This way, the temperature and surface gravities of the secondary stars are constrained from the spectral fit. The radius of the secondary star is estimated by scaling a theoretical model to match the SEDs fluxes at the distance implied by the *Gaia* parallaxes.

For each target we selected the spectra with the highest signal-to-noise ratio and normalised the continuum of the spectra. Based on the initial parameter estimates of the stars from the CERES pipeline, synthetic spectra were generated with iSpec (Blanco-Cuaresma et al. 2014). We then used the synthetic spectral fitting technique within iSpec to measure the effective temperatures, surface gravities, metallicities [Fe/H] and rotational broadenings for our targets. The limb darkening coefficient was kept fixed at a value of 0.6 and the spectral resolution was fixed at the measured value for each observation based on fits to sky lines. A small section of the spectra and the best fit models are shown in Fig. 2.

SEDs of our targets including *V* band data from the APASS (DR9, Henden et al. 2015), *J*, *H* and *K_s* band data from 2MASS

(Cutri et al. 2003) and *W1* and *W2* band data from *WISE* (Cutri & et al. 2012). We fitted the SEDs with BT-NextGen theoretical spectra (Allard et al. 2012) using a custom Python code. The input parameters of the code are the effective temperature, surface gravity and radius of the star and the parallax and reddening of the system. The theoretical spectra were scaled by a factor $(R/D)^2$, where R is the radius of the star derived from the SED we create of the secondary star, D is the distance. The distributions of our model parameters were found using the Markov chain Monte Carlo (MCMC) method (Press et al. 2007) implemented using the python package EMCEE (Foreman-Mackey et al. 2013), where the likelihood of accepting a model was based on the χ^2 of the fit and additional prior probabilities on the effective temperature and surface gravity from the iSpec fit, the parallax from *Gaia* DR2 (Gaia Collaboration et al. 2018b) and the reddening from 3D interstellar dust maps (Lallement et al. 2019). A short initial MCMC chain was used to determine the approximate parameter values, which were then used as the starting values in a longer "production" chain, to determine the final values and their uncertainties. The production chain used 50 walkers, each with 10 000 points. The first 1000 points were classified as "burn-in" and were removed from the final results. Stellar masses were then computed using the fitted surface gravities and radii. The best fit values and their uncertainties are listed in Table 2 and plotted in Fig. 3, along with the parameter distributions.

3.2 Orbital periods

The most critical parameter constraining the past and future of a given PCEB is its orbital period. We used the least-squares spectral analysis method based on the classical Lomb-Scargle periodogram (Lomb 1976; Scargle 1982) that was implemented in the astroML python library (VanderPlas et al. 2012) and *The Joker* (Price-Whelan et al. 2017, 2020), a specialized Monte Carlo sampler created to find converged posterior samplings for Keplerian orbital parameters. *The Joker* is especially useful for non-uniform data and allows to also identify eccentric orbits. As we have previously identified systems with eccentric orbits that most likely represent hierarchical triples with white dwarf components (Paper III), double checking circular orbits derived from Lomb-Scargle periodograms with *The Joker* was required. *The Joker* found the period corresponding to the highest peak in the periodograms shown in Fig. 5 and the eccentricity in all three cases to be insignificant. The highest peaks in the periodograms shown in Fig. 5 also clearly provided the best fit to the data. All other aliases with similar power in the periodogram disagreed drastically with some radial velocity measurements and the χ^2 values resulting from sine fits using the periods corresponding to the second highest peaks in the periodograms are at least a factor 33 higher in all cases and can therefore be excluded. The measured orbital periods for TYC 4700, TYC 1300 and TYC 4962 are 2.4667, 1.6127 and 1.2798 days respectively. The final sine fits to our radial velocity data are shown in Fig. 4. Radial velocities from different telescopes are plotted in different colors.

3.3 The white dwarf mass

The masses of white dwarfs in binaries are notoriously difficult to measure and usually UV spectra e.g taken with the Hubble Space Telescope (*HST*) are required to get reliable estimates (e.g. Paper I). For the three systems presented here, *HST* spectra are not available and the presence of a white dwarf is only indicated by excess UV emission detected with *GALEX*. However, based on the parameters

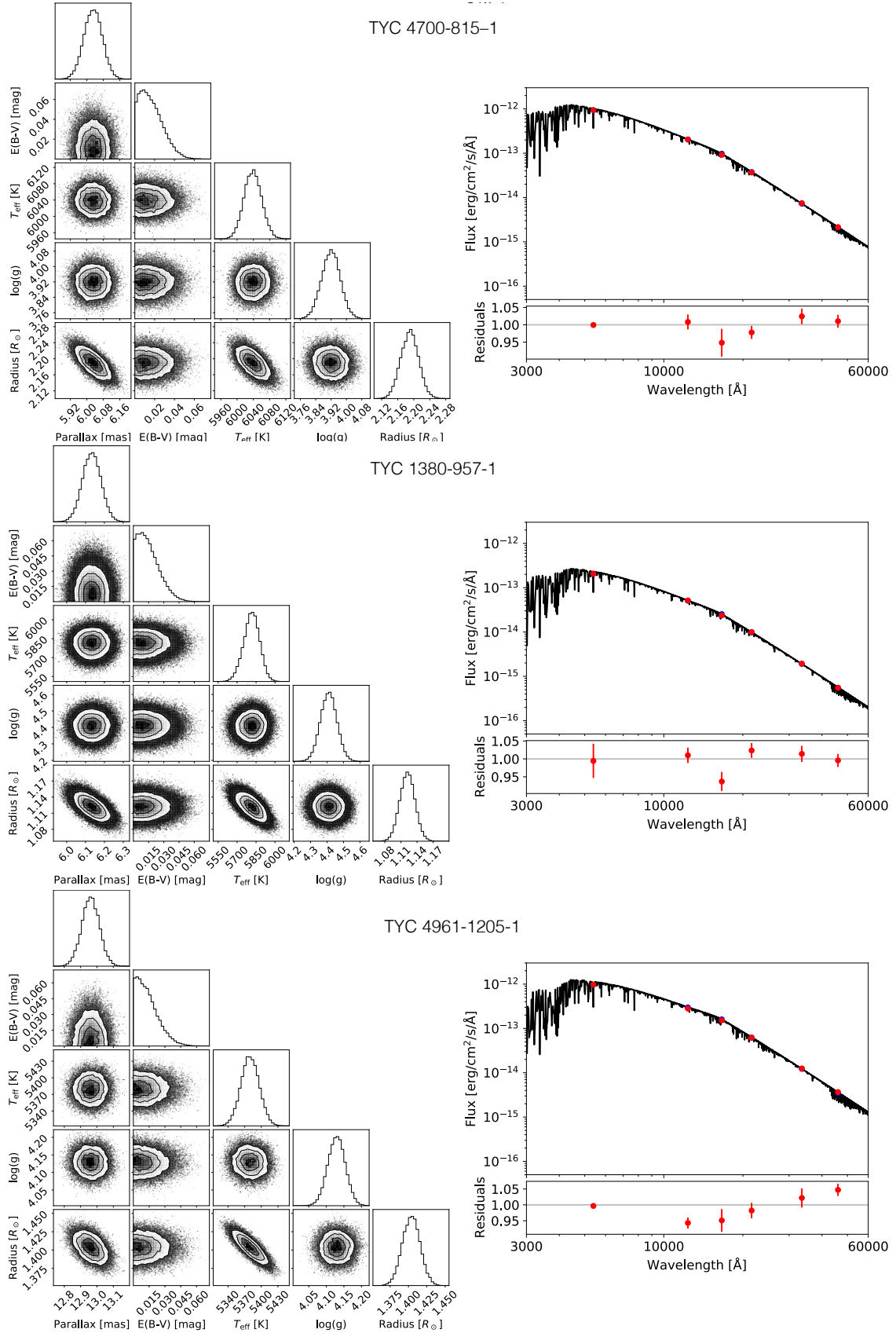


Figure 3. Spectral energy distributions created from AAVSO Photometric All Sky Survey, J , H and K_s band data from 2MASS, and $W1$ and $W2$ band data from WISE were fitted with BT-NextGen theoretical spectra. We used as input parameters the previously estimated effective temperature, surface gravity and radius of the star and the parallax and reddening of the system. The Markov chain Monte Carlo method was used to derive posterior distributions from which we derived our final parameters and the corresponding uncertainties.

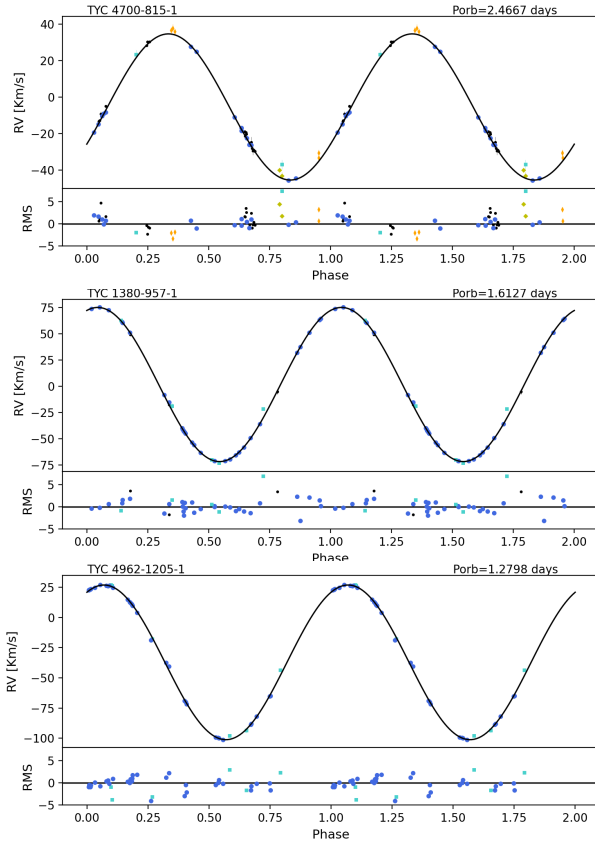


Figure 4. Phase-folded radial velocity curve for the main-sequence stars with the corresponding residuals to the best fit. The colors indicate which instrument was used to obtain the measurement: FEROS (blue big dots), ESPRESSO (cyan squares), HRS (black small dots), MRES (orange thin diamonds), and UVES (green diamonds).

estimated for the secondary star we can derive rough estimates for the white dwarf mass. A lower limit for the white dwarf mass can be obtained based on the measured secondary mass and orbital period and assuming the maximum inclination of 90 degrees. Assuming furthermore that the orbit of the secondary is synchronised, a reasonable assumption for the short orbital periods we measured, we can estimate a range of possible white dwarf masses based on the $v \sin(i)$ and radius measured for the secondary star using the binary mass function.

Both estimates, the minimum white dwarf mass and the one based on $v \sin(i)$ and synchronised rotation are shown in Table 2. We note that especially the latter estimate must be taken with a grain of salt, as the secondary star parameters were obtained by fitting single main sequence star model spectra to observed spectra of secondary stars in close binaries. This might imply that our $v \sin(i)$ measurements are affected by systematic errors that most likely exceed the purely statistical errors we provide in Table 2.

4 DISCUSSION

We have identified three close binary stars consisting of a G-type star and most likely a white dwarf whose presence is indicated by a significant UV excess, and measured the stellar and binary parameters of the three systems. In what follows we investigate possible

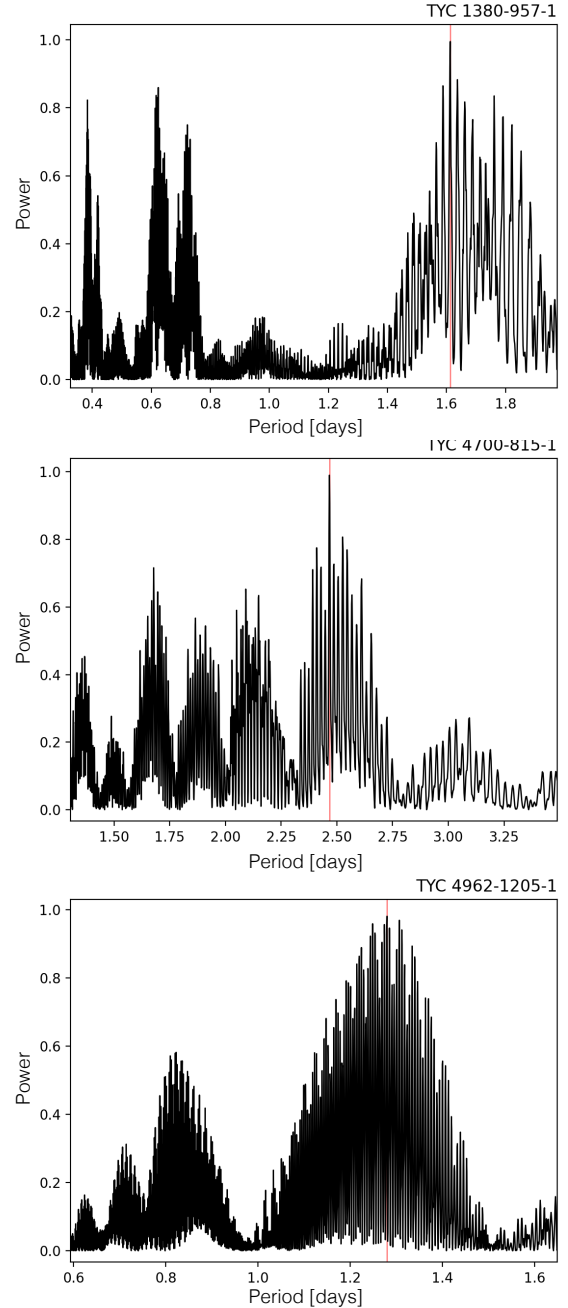


Figure 5. Lomb-Scargle periodograms of TYC 1380 (top panel), TYC 4700 (middle) and TYC 4962 (bottom). The red lines highlight the highest peak which corresponds to the orbital period of the systems.

implications for white dwarf binary formation and evolution by reconstructing their past and predicting their future. Finally, we will discuss how certain we can be that our interpretation that the UV excess comes from a white dwarf is the most likely one.

4.1 The evolutionary history

The main purpose of the white dwarf binary pathway project is to progress with our understanding of the formation and evolution of close white dwarf binaries with secondary stars in the mass range $\sim 0.5\text{--}1.5 M_{\odot}$. Previous surveys of close white dwarf main

Table 2. Stellar Parameters

Parameter	TYC 4700-815-1	TYC 1380-957-1	TYC 4962-1205-1
Orbital Period [days]	2.4667 ± 0.0087	1.6127 ± 0.0109	1.2798 ± 0.0026
RV amplitude [Km s ⁻¹]	40.067 ± 0.083	73.32 ± 0.025	64.057 ± 0.014
V _y [Km s ⁻¹]	-5.46 ± 0.044	1.70 ± 0.020	-37.19 ± 0.012
v sin(i) [Km s ⁻¹]	45.3 ± 4.86	30.00 ± 2.84	37.30 ± 3.11
Inclination [deg]	63-90	49-71	37-47
Distance [pc]	165.84 ± 1.20	163.07 ± 1.29	77.141 ± 0.321
Sec. T _{eff} [K]	6038.7 ± 31.7	5815.5 ± 64.7	5380 ± 30
Sec. Mass [M _⊙]	1.454 ± 0.171	1.181 ± 0.145	0.969 ± 0.058
Sec. Radius [R _⊙]	2.190 ± 0.023	1.122 ± 0.014	1.404 ± 0.012
Sec S Type	G0IV	G2V	G6V
Sec log g [dex]	3.920 ± 0.049	4.410 ± 0.050	4.129 ± 0.025
Sec. Metallicity [Z/H]	-0.04 ± 0.13	-0.03 ± 0.07	-0.42 ± 0.07
Sec. Luminosity [L _⊙]	5.713 ± 0.093	1.290 ± 0.04	1.479 ± 0.014
Minimum WD Mass [M _⊙]	0.38	0.59	0.4
WD Mass(i) [M _⊙]	0.38-0.44	0.64-0.85	0.59-0.77

sequence binaries nearly exclusively identified systems with M-dwarf companions. These surveys were fundamental to constrain the models for common envelope evolution. The classical energy conservation equation that relates the energy required to unbind the envelope and the released orbital energy can be written as

$$E_{\text{bind}} = \alpha_{\text{CE}} \Delta E_{\text{orb}} \quad (1)$$

where α_{CE} is the common envelope efficiency. This equation is usually used to estimate the outcome of common envelope evolution. Typically a simple expression for the binding energy is used, i.e.

$$E_{\text{bind}} = -\frac{GM_1 M_{1,e}}{\lambda R_1} \quad (2)$$

where M_1 , $M_{1,e}$ and R_1 are the total mass, envelope mass and radius of the primary star at the onset of the common envelope phase, respectively. The binding energy parameter λ strongly depends on the mass and the radius of the white dwarf progenitor when it fills the Roche-lobe and dynamically unstable mass transfer starts. A more general form of the binding energy equation takes into account possible contributions from recombination energy (U_{rec}) in the envelope ejection process and a corresponding second common envelope parameter α_{rec} , i.e.

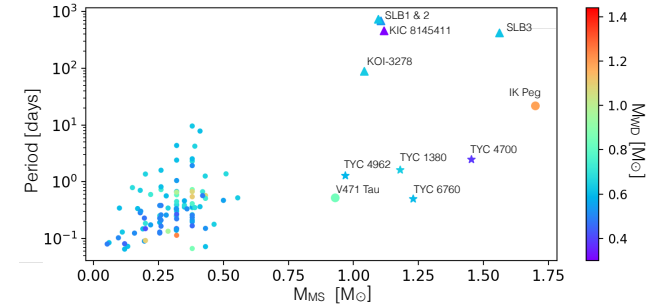
$$E_{\text{bind}} = \int_{M_{1,c}}^{M_1} -\frac{Gm}{r(m)} dm + \alpha_{\text{rec}} \int_{M_{1,c}}^{M_1} U_{\text{rec}}(m). \quad (3)$$

Combining the above equations with stellar evolution tracks, it is possible to estimate the progenitor system parameters for a given PCEB and a given common envelope efficiency. This technique has been pioneered by Nelemans & Tout (2005). For low mass secondary stars, Zorotovic et al. (2010) showed that all known systems can be reconstructed with $\alpha_{\text{CE}} = 0.2-0.3$. Their conclusion that, at least for low-mass secondary stars, the outcome of common envelope evolution can be reproduced by assuming a low value of α_{CE} has later been confirmed by various comparisons of the predictions of binary population synthesis with the observed sample (e.g. Camacho et al. 2014; Toonen & Nelemans 2013; Zorotovic et al. 2014b, 2011).

One of the key questions the white dwarf binary pathway project shall answer is whether the outcome of common envelope evolution for systems with more massive secondary stars can be described with the same formalism and the same efficiency parameter that is now well established for systems with M-dwarf secondaries. Eight systems with secondary stars earlier than M and with orbital periods that clearly imply that mass transfer must have occurred in the progenitor systems, potentially common envelope evolution, have been previously discovered: IK Peg, KOI-3278, SLB1, SLB2, SLB3, and KIC 8145411 with much longer orbital periods of ~ 21 ,

Table 3. Binary parameters obtained by reconstructing common envelope evolution where $M_{1,i}$ and $P_{\text{orb},i}$ denote the initial primary mass and orbital period. The first row provides the parameters for $\alpha_{\text{CE}} = 0.2-0.3$ while in the second row the efficiency was kept as a free parameter.

Object	α_{CE}	WD Mass [M _⊙]	$M_{1,i}$ [M _⊙]	$P_{\text{orb},i}$ [days]	Age[Gyr]
TYC 4700	0.2–0.3	0.40–0.44	1.30–1.53	203–386	2.78–4.81
	0.18–1.0	0.38–0.44	1.29–1.90	109–395	
TYC 1380	0.2–0.3	0.64–0.85	2.54–3.94	706–2222	0.22–0.77
	0.10–1.0	0.64–0.85	2.36–4.06	126–2222	
TYC 4962	0.2–0.3	0.59–0.77	2.28–3.36	617–1917	0.35–1.04
	0.07–1.0	0.59–0.77	1.75–3.46	103–1917	


Figure 6. Period distribution of PCEBs including the three systems we present in this work. A large number of PCEBs with M-dwarf companion (small dots) have been characterized and their orbital period distribution peaks at a few hours (Nebot Gómez-Morán et al. 2011). White dwarf binaries with G-type secondary star (the four systems from our pathway survey are represented with stars) cover a much wider range of periods which extends from half a day to several 100 days. While the short period systems are certainly PCEBs, the origin of the systems with longer periods is less clear. Big circles are previously known PCEBs and triangles are Self-Lensing Binaries.

88, 683, 728, 419, and 455 days respectively (Wonnacott et al. 1993; Kruse & Agol 2015; Kawahara et al. 2018; Masuda et al. 2019) as well as the shorter orbital period systems V471 Tau (O’Brien et al. 2001) and TYC 6760-497-1 (Parsons et al. 2015), both with very short periods of just 12 hours. Two additional potential members of the class of detached PCEBs with secondary stars earlier than M are V1082 Sgr and GPX-TF16E-48 (Tovmassian et al. 2018; Krushinsky et al. 2020). However, for now we exclude these two systems from our list of PCEBs as in the first case the mass transfer rate seems unusually high for a detached system and as in the second case the secondary star parameters appear to be still somewhat uncertain.

The orbital periods of the currently known potential PCEBs, including the three systems presented here, are shown in Fig. 6 as a function of secondary mass. For those systems with orbital periods exceeding 20 days, orbital energy as the only energy source during common envelope evolution is not sufficient to explain their long orbital periods. As shown by Zorotovic et al. (2014a) such systems require additional energy, most likely recombination energy, to contribute to expelling the envelope. Otherwise, these systems can not be explained in the context of common envelope evolution.

With the three systems presented in this work, our dedicated survey for PCEBs with massive secondary stars has now found four systems with short orbital periods. The evolutionary history of the previously identified PCEB TYC 6760-497-1 can be reconstructed assuming common envelope evolution without contributions from recombination energy and using the same efficiency that has been shown to work best for PCEBs with M-dwarf secondaries.

We applied the algorithm developed by Zorotovic et al. (2010, 2011, 2014a) which combines the stellar evolution code *BSE* (Hurley et al. 2002) with the above listed common envelope energy equations to reconstruct the evolutionary history of the three new PCEBs with G-type secondary star that we observationally characterized in the previous sections.

In short, the CE reconstruction works as follows. We calculated a grid of stellar evolution tracks with initial masses up to eight solar masses (with a step size of $0.01M_{\odot}$). As the core mass of the giant star at the onset of CE evolution must have been equal to the white dwarf mass in the currently observed post-CE system, we searched the grid for giant stars with a core mass equal to the estimated range of possible white dwarf masses. As the radius and the mass of each of the identified potential white dwarf progenitor stars is also given by the stellar model, we can determine the orbital period at the onset of CE evolution for all these candidates using the mass of the companion and Roche-geometry. Knowing the stellar masses, the orbital period, and the envelope mass at the onset of CE evolution as well as the current configuration, allows then to determine whether sufficient orbital energy was provided to unbind the envelope of the giant and how efficient the use of orbital energy must have been.

Using this algorithm and leaving the common envelope efficiency as a free parameter we found possible progenitor systems for a relatively large range of values for α_{CE} (see Table 3). While the low white dwarf mass found for TYC 4700 implies that the progenitor filled its Roche-lobe on the first giant branch, for the other two systems two scenarios are possible. Either the progenitor filled its Roche-lobe on the early AGB (low values of α_{CE}) or as a thermally pulsating AGB star ($\alpha_{CE} \gtrsim 0.4$).

With now five out of eleven of all known white dwarf binaries with secondaries earlier than M having short orbital periods and not requiring extra energy during common envelope evolution, it seems that perhaps a population of PCEBs with such secondaries exists that forms identical to PCEBs with low mass secondaries. A crucial question that remains to be answered is what the origin of the longer orbital period systems is. If it is common envelope evolution as well, then one would need to explain why such long orbital period systems are exclusively found among PCEBs with more massive secondaries. Alternatively, at least for some of these longer period systems, it might not have been common envelope evolution that led to their formation. However, firm conclusions cannot be drawn from these relatively few known systems. Therefore, it is of great importance for our understanding of white dwarf binary formation to continue measuring orbital periods of close WD+AFGK binaries.

4.2 The future of the three systems

PCEBs with more massive secondary stars are crucial for our understanding of close white dwarf binary formation. Their total mass exceeds the Chandrasekhar limit and therefore, in contrast to most white dwarf plus M-dwarf binaries, some of these systems might be potentially SNIa progenitors. The future of these PCEBs with solar type secondary stars, however, is more diverse than that of PCEBs with late-type secondary stars, which virtually all become CVs sooner or later (in some cases it may take several Hubble times). The future of PCEBs with G type secondary stars depends on the orbital period at the end of common envelope evolution, on the angular momentum loss mechanism that drives the system to shorter orbital periods, on the mass ratio of the binary and on the evolutionary status of the secondary star. The richness and variety of evolutionary futures of close WD+AFGK binaries is illustrated by the three systems characterized in this work. All three systems

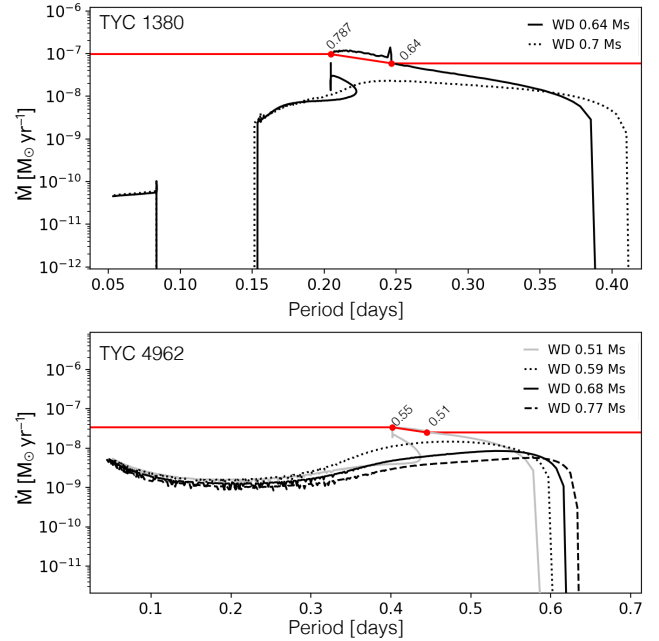


Figure 7. The future evolution of TYC 1380 and TYC 4962 for two and four different white dwarf masses respectively. While TYC 1380 will evolve into a normal CV after perhaps passing through a period of high mass transfer (above the red line) which would convert the system into a super soft source for ≈ 1.7 Myr. TYC 4962 will become a CV with an evolved donor star and therefore, in contrast to CVs with unevolved donors, not detach in the orbital period range of 2 – 3 hr. This system will only appear as a super soft source if the white dwarf mass is smaller than the value estimated from $v \sin(i)$ (grey line, bottom panel).

contain a G-type secondary star and a white dwarf with a mass typical for these objects and a short orbital period between 1.2 and 2.5 days. Despite these similarities, they face very different futures.

In order to estimate the future of these systems in more detail, we performed dedicated simulations using MESA (Paxton et al. 2011) and the same set-up and the same criteria for stable mass transfer as in Parsons et al. (2015). For the current parameters we adopt the parameters listed in Table 2. The white dwarf mass estimates derived from the observations are quite uncertain. We therefore experimented with different white dwarf masses covering the range of masses and the minimum mass given in Table 2. We discuss in more details simulations for different white dwarf masses which have been chosen to illustrate how the future evolution depends on this parameter.

4.2.1 TYC 1380-957-1: potentially the second known super soft source progenitor

Simulations for TYC 1380 were performed assuming $M_{MS} = 1.18 M_{\odot}$ and different white dwarf masses covering the range constrained by our observations, i.e. we used 0.59, 0.64, 0.65, 0.68, 0.7, 0.73, 0.85 M_{\odot} . The current orbital period of the system is $P_{orb} = 1.6127$ days. Two evolutionary tracks are shown in the upper panel of Fig. 7. We find that the future of the system depends sensitively on the white dwarf mass. For the white dwarf masses exceeding $\approx 0.64 M_{\odot}$ the mass transfer remains dynamically and thermally stable and the system evolves into a CV. As an example for this evolution Figure 7 shows the track calculated with a white

dwarf mass of $0.7 M_{\odot}$. In this case, in 1.2 Gyr from now, when the orbital period has reduced to 9.84 hrs, mass transfer starts but never reaches values large enough to ignite stable nuclear burning on the surface of the white dwarf. Nova eruptions should therefore occur and the mass of the white dwarf remains nearly constant (non-conservative mass transfer). Thus, if the real white dwarf mass is $\geq 0.65 M_{\odot}$, TYC 1380 would simply be a pre-CV with a large secondary mass.

However, if the white dwarf mass is $0.64 M_{\odot}$ or smaller, the mass ratio becomes large enough to trigger thermal time scale mass transfer and stable hydrogen burning on the surface of the white dwarf. This future seems, however, somewhat less likely because thermal time scale mass transfer will only occur when the white dwarf mass is at the lower limit of the mass range derived from the $v \sin(i)$ measurement. The corresponding evolutionary track is also shown in Fig. 7. In ~ 1.1 Gyr from now the system will appear as a super soft source. During this phase we assume that 90 per cent of transferred mass remain on the white dwarf which therefore increases in mass. While the assumption of mass growth by 90 per cent of the accreted mass is arbitrary, we performed several numerical experiments with different mass growth fractions and found that the outcome of the evolution of this system does not sensitively depend on this parameter. When the white dwarf mass reaches $\sim 0.79 M_{\odot}$ at an orbital period of ~ 4.8 hrs, the mass ratio becomes small enough for stable angular momentum loss driven mass transfer and the system becomes a CV. Thus, while TYC 1380 will in any case become a CV, it might evolve through a short super soft X-ray binary phase if the white dwarf mass is equal to or smaller than $\sim 0.64 M_{\odot}$.

4.2.2 TYC 4700-815-1: unstable mass transfer and stellar merger

The future of TYC 4700 is straight forward to predict. Given its current orbital period of $P_{\text{orb}} = 2.4667$ days, the mass ratio of $3.3 < q < 3.8$ and the secondary mass and radius of $1.45 M_{\odot}$ and $2.2 R_{\odot}$ the system will run into dynamically unstable mass transfer when the secondary fills its Roche-lobe in 490 Myr from now when the orbital period will be 1.53 days. This is predicted by our MESA simulations in agreement with Ge et al. (2015, their figures 8 and 9). The orbital energy of the binary will not be sufficient to expel the envelope and therefore the two stars will merge and evolve into a giant star with the core of this giant being the white dwarf we observe today. This giant will then evolve into a single white dwarf as the final fate of TYC 4700.

4.2.3 TYC 4962-1205-1: a progenitor of a CV with evolved donor

For TYC 4962 we adopted $M_{\text{MS}} = 0.969 M_{\odot}$ and three white dwarf masses corresponding to the range we estimated from the observations and one more closer to the minimum mass, i.e. we assumed $M_{\text{WD}} = 0.51, 0.59, 0.68$ and $0.77 M_{\odot}$. The current orbital period of the system is $P_{\text{orb}} = 1.2798$ days. The radius we estimated for the secondary star, $R_{\text{MS}} = 1.4 R_{\odot}$, implies that it has significantly evolved from the ZAMS. The starting model for our MESA simulations was obtained by evolving the secondary until it reaches the required radius of the evolved secondary star. According to this calculation the current age of the object is 6.5 Gyr and the secondary star has a temperature of 6172 K. Driven by angular momentum loss through magnetic braking the system will start mass transfer in 335 Myr at an orbital period of ~ 0.6 days somewhat depending on the white dwarf mass. Evolving the system further we found that,

independent of the mass of the white dwarf, the system evolves into a CV with a significantly evolved secondary star. In such systems the secondary does not become fully convective and therefore mass transfer is driven by magnetic braking until it reaches the orbital period minimum. The system does not evolve as detached binary through the orbital period gap (see bottom panel of Fig 7). This evolutionary path for CVs with evolved secondary stars has been predicted by Kolb & Baraffe (2000) and several such CVs have been found (e.g. Thorstensen et al. 2002; Rebassa-Mansergas et al. 2014). As our white dwarf mass estimates based on $v \sin(i)$ is somewhat uncertain, we also ran a simulation assuming a smaller white dwarf mass. If the white dwarf mass is as low as $0.51 M_{\odot}$, thermal time scale mass transfer will be triggered and stable hydrogen burning on the surface of the white dwarf will begin at a period of 0.44 days. This thermal time scale mass transfer will then continue for the next 582 Myr, until the white dwarf reaches the mass of $0.55 M_{\odot}$. The system will then become a cataclysmic variable with an evolved donor secondary star. Given the small white dwarf mass that is required for this path to occur, we consider this scenario rather unlikely. For completeness, the corresponding simulation is nevertheless shown in Fig. 7. In any case, TYC 4962 is the first known progenitor of a CV with evolved donor. The future evolution is driven by the nuclear evolution of the secondary star and therefore virtually independent on the white dwarf mass.

While it is clear that TYC 4962 will evolve into a CV with an evolved donor star, we note a minor inconsistency in our calculations. The current temperature of the secondary is much smaller than that of an evolved single star with the same mass. To match the observed mass and radius in the MESA calculations, we had to accept a temperature that clearly exceeds the observed value. In close binaries such as the systems presented here, the secondary star's rotation is synchronised with the orbital motion. The resulting fast rotation of the secondary star offers an explanation for the discrepancy between the MESA models and the measured temperature. Fast rotating stars are potentially very active which implies that they are probably larger and cooler than a single stars of the same mass. Therefore, also the evolutionary time the system needed to reach the radius we derived from the observations, was probably significantly less than the 6.5 Gyr we obtained with MESA for a single star.

4.3 Excluding contamination from active M-dwarf companions

In Parsons et al. (2015) and Paper I we presented *HST* spectra of ten of our RAVE candidates with the aim of confirming that the UV excess is, at least in most cases, indicative of a white dwarf. We found that this is indeed the case: nine out of the ten *HST* spectra confirmed the presence of a white dwarf and in one case the UV excess came from a hot subdwarf or pre-helium white dwarf. While this experiment clearly showed that our target selection works very well, for individual objects, we cannot per se exclude the possibility that an unseen flaring/active M-dwarf companion is producing the UV excess instead of a white dwarf.

However, flaring M-dwarfs are very unlikely to produce the *GALEX* fluxes of our three targets. Assuming that the optically unseen stellar component is an M-dwarf with $T_{\text{eff}} = 4000$ K and $\log g = 4.6$, according to the BT-NextGen spectral library (Allard et al. 2012) the magnitude difference between the M dwarf model and the observed NUV magnitude is 9.51, 11.7 and 8.81 for TYC 1380, TYC 4700 and TYC 4962, respectively. These values are

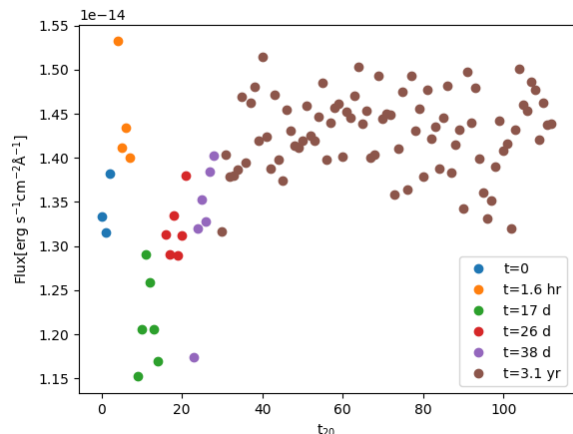


Figure 8. NUV light curve of TYC 1380. Each color represent one of the six exposures. The value of t represent the starting time of the exposure relative to the first exposure marked in blue, while t_{20} stands for 20 second bins.

far beyond the highest values of the observed Δm_{NUV} distribution found for flaring M-dwarfs (Welsh et al. 2007).

For one of the systems discussed in this paper, we could complement the above finding with the *GALEX* light curve. We used *GPHOTON* (Million et al. 2016) to produce an NUV light curve for TYC 1380 with 20 second binning in each exposure and an aperture radius of 0.004° and the resulting light curve is shown in Fig. 8. Taking into account that the fluxes represented by the green dots correspond to observations where the position of the system was close to the edge of the *GALEX* detector and are therefore likely compromised, Fig. 8 clearly excludes that the UV excess in TYC 1380 is the result of a flaring M-dwarf.

As a final test, we check whether a white dwarf model can produce the observed UV excess in TYC 4962 which is the system with the smallest UV excess of the three systems we discussed in this work. Using a chi-square fitting and interpolating the UV magnitudes of the cooling models³ of Holberg & Bergeron (2006), Kowalski & Saumon (2006) and Tremblay et al. (2011), we found the best-fit stellar parameters and distance of the white dwarf. Figure 9 shows that the UV excess is perfectly matched (p-value < 0.02) using the white dwarf model spectrum with $T_{\text{eff}} = 11\,000\text{ K}$, $\log g = 8$, $M = 0.6 M_{\odot}$, at the distance measured with *GAIA*. While this fit is highly degenerate – we only have two data points in the UV – it clearly shows that a white dwarf model can reproduce the UV excess.

We therefore conclude that the presence of a white dwarf in the three targets presented here is by far the most reasonable interpretation of the UV excess.

5 CONCLUSIONS

We presented the identification of three close binaries containing a G-type star and a white dwarf. The presence of the latter is indicated by a UV excess detected with *GALEX*. Based on measurements of the orbital and stellar parameters of these systems, we reconstructed their past and predicted their future. We find that all three systems

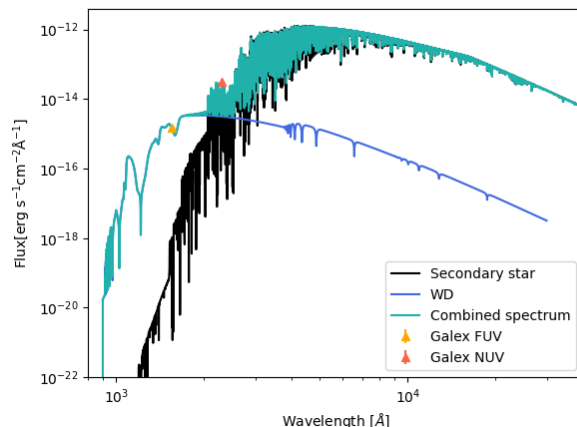


Figure 9. Best white dwarf model fit to the observed UV excess in TYC 4962. The synthetic spectrum of the white dwarf from Koester (2010) (shown in blue) corresponds to $T_{\text{eff}} = 11\,000\text{ K}$, $\log g = 8$, and $M = 0.6 M_{\odot}$. While this is clearly just one possible solution, it illustrates that the UV excess can be perfectly explained by the presence of a white dwarf.

have most likely formed through common envelope evolution. In contrast to previously known PCEBs with G-type companions, in the three systems presented here we do not find any indications for additional sources of energy acting during common envelope evolution. Despite having similar stellar masses and orbital periods, the future of all three systems is different. TYC 4962 will evolve into a cataclysmic variable (CV) star with an evolved donor, TYC 4700 will not survive the next phase of mass transfer and evolve into a single white dwarf, and TYC 1380 will perhaps appear as a super soft X-ray source before becoming a rather typical CV. This shows that the future evolution as well as the past of PCEBs with AFGK secondary stars is more complicated than those of PCEBs with M-dwarf companions. We therefore conclude that the identification and observational characterization of more PCEBs consisting of a white dwarf with an AFGK companion is likely to provide crucial constraints of white dwarf binary evolution theories with potential implications for our understanding of the progenitors of thermonuclear supernovae.

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³ <http://www.astro.umontreal.ca/bergeron/CoolingModels/>

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DATA AVAILABILITY

The data underlying this article will be shared on reasonable request to the corresponding author.

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Table A1. TYC 4700-815-1 radial velocity measurements.

Instrument	Exptime [s]	S/N	BJD	RV [$Km\ s^{-1}$]	error [$Km\ s^{-1}$]
HRS	900	40	2458092.06920	-25.30	0.49
HRS	900	35	2458175.99240	-29.67	0.72
HRS	430	28	2458130.08178	-5.03	0.77
HRS	560	40	2458169.96268	28.34	0.50
HRS	560	40	2458169.97201	30.34	0.48
HRS	560	40	2458169.98144	30.07	0.65
HRS	560	40	2458169.99050	30.22	0.59
HRS	500	45	2458170.95383	-18.84	0.50
HRS	400	53	2458170.96179	-21.31	0.68
HRS	240	37	2458170.96850	-19.33	0.68
HRS	200	37	2458170.97399	-20.22	0.76
HRS	240	35	2458171.94807	-13.36	0.65
HRS	480	35	2458171.95351	-12.85	0.70
HRS	480	25	2458171.96850	-9.26	1.22
HRS	540	35	2458171.97273	-10.59	0.60
HRS	540	18	2458175.95720	-24.88	2.96
HRS	240	30	2458175.96460	-28.19	1.13
HRS	120	30	2458175.97270	29.59	2.14
HRS	240	35	2458175.98238	-28.99	0.70
MRES	120	21	2458184.03678	-30.60	1.81
MRES	240	28	2458184.04033	-33.22	1.34
MRES	60	25	2458185.01066	36.60	1.87
MRES	60	20	2458185.02971	37.64	2.27
MRES	150	24	2458185.04959	35.81	1.61
FEROS	120	39	2458486.59061272	-11.03	0.17
FEROS	120	37	2458486.66071053	-17.01	0.20
FEROS	120	60	2458486.66322596	-18.54	0.20
FEROS	120	48	2458487.63613462	-19.36	0.14
FEROS	120	57	2458487.68227370	-14.85	0.14
FEROS	120	53	2458488.61826496	27.69	0.15
FEROS	120	38	2458488.67449235	24.83	0.13
FEROS	120	23	2458489.60644339	-45.69	0.20
FEROS	120	53	2458489.67938950	-44.66	0.27
FEROS	120	42	2458491.64877842	-22.43	0.15
FEROS	120	34	2458491.67602759	-26.08	0.17
FEROS	120	41	2458491.69341871	-26.29	0.17
FEROS	120	40	2458492.65260996	-10.47	0.19
FEROS	120	36	2458492.67067858	-9.11	0.18
FEROS	120	30	2458492.68748925	-8.27	0.17
UVES	60	93	2458353.845644055	-40.01	1.36
UVES	60	37	2458390.870005791	-43.15	1.96
ESPRESSO	1200	20	2458802.8106283	-37.05	2.62
ESPRESSO	1200	19	2458803.8033296	23.41	2.39

Table A2. TYC 1380-957-1 radial velocity measurements.

Instrument	Exptime [s]	S/N	BJD	RV [$Km\ s^{-1}$]	error [$Km\ s^{-1}$]
HRS	1200	35	2458171.10569	49.36	1.0
HRS	1400	35	2458176.20292	-17.56	1.2
HRS	2000	26	2458243.03968	-4.99	2.2
ESPRESSO	1200	18	2458390.9747601	-69.82	1.19
ESPRESSO	1200	19	2458391.0227932	-72.73	0.92
ESPRESSO	1200	21	2458391.9867707	62.85	1.02
ESPRESSO	1200	22	2458450.9844952	-21.61	1.34
ESPRESSO	1200	20	2458451.9935934	-18.70	2.16
FEROS	120	18	2458486.68750550	32.12	0.13
FEROS	120	16	2458486.70998961	37.63	0.13
FEROS	120	17	2458487.66240839	-62.96	0.10
FEROS	120	16	2458487.75553172	-71.04	0.10
FEROS	120	31	2458488.66635571	72.53	0.10
FEROS	120	21	2458489.66670363	-35.87	0.17
FEROS	240	29	2458490.76647412	-39.85	0.09
FEROS	240	30	2458490.77041791	-40.92	0.09
FEROS	240	10	2458490.77430824	-42.00	0.09
FEROS	240	25	2458490.77819312	-42.96	0.09
FEROS	240	22	2458490.78256429	-43.78	0.09
FEROS	240	25	2458490.78825937	-44.96	0.09
FEROS	120	23	2458490.83072595	-53.54	0.13
FEROS	120	24	2458491.67514523	63.23	0.11
FEROS	120	24	2458491.68357584	64.50	0.11
FEROS	120	12	2458491.77918888	73.61	0.14
FEROS	120	16	2458492.66294326	-70.92	0.11
FEROS	120	17	2458492.69528291	-69.29	0.10
FEROS	120	19	2458492.73411412	-65.38	0.09
FEROS	120	13	2458492.75539783	-62.74	0.12
FEROS	120	13	2458492.78775605	-58.20	0.08
FEROS	180	15	2458600.50883313	-55.86	0.10
FEROS	180	21	2458601.49588304	75.18	0.09
FEROS	180	20	2458602.49497791	-49.29	0.07
FEROS	180	22	2458603.53605945	-8.12	0.09
FEROS	180	22	2458604.49805593	51.15	0.12
FEROS	180	19	2458605.48214311	-70.92	0.09
FEROS	240	24	2458606.48415264	61.11	0.08
FEROS	240	12	2458606.48754113	60.38	0.07

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APPENDIX A: RADIAL VELOCITIES MEASURED FOR ALL THREE SYSTEMS

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Table A3. TYC 4962-1205-1 radial velocity measurements.

Instrument	Exptime [s]	S/N	BJD	RV [$Km s^{-1}$]	error [$Km s^{-1}$]
FEROS	120	49	2458290.63890684	10.55	0.04
FEROS	120	54	2458290.64296182	9.74	0.04
FEROS	120	51	2458290.51895498	26.46	0.04
FEROS	300	69	2458290.61714122	14.96	0.04
FEROS	300	61	2458290.62804077	12.82	0.04
FEROS	120	61	2458290.63212027	11.97	0.06
FEROS	180	65	2458290.47164068	27.18	0.13
FEROS	180	64	2458292.54191363	-88.89	0.06
FEROS	180	66	2458292.54469778	-87.98	0.05
FEROS	180	67	2458292.64038933	-65.31	0.06
FEROS	180	65	2458292.64444679	-64.99	0.05
FEROS	300	6	2458293.48043724	-70.12	0.06
FEROS	300	56	2458293.48450716	-71.74	0.31
FEROS	300	39	2458293.64986508	-99.76	0.07
FEROS	300	55	2458293.65394596	-99.83	0.07
FEROS	300	95	2458294.47693482	11.14	0.06
FEROS	120	61	2458294.57760653	-18.89	0.07
FEROS	360	53	2458294.65813593	-37.52	0.05
FEROS	120	49	2458295.56410647	24.82	0.05
FEROS	300	59	2458295.62289638	26.42	0.05
FEROS	300	80	2458600.53527200	-40.52	0.05
FEROS	120	74	2458600.61655423	-69.08	0.05
FEROS	300	68	2458601.52255787	24.66	0.06
FEROS	300	48	2458601.64901651	4.00	0.06
FEROS	300	67	2458602.77673050	26.15	0.05
FEROS	120	68	2458603.55421825	-81.95	0.05
FEROS	300	69	2458604.61991781	-99.22	0.07
FEROS	300	65	2458604.65812713	-101.21	0.07
FEROS	300	62	2458606.51420823	22.62	0.05
FEROS	120	59	2458606.51687580	22.95	0.05
FEROS	120	64	2458606.51954985	23.15	0.05
FEROS	120	70	2458606.52224277	23.40	0.05
FEROS	120	68	2458606.52493824	23.68	0.05
ESPRESSO	1200	30	2458857.0716883	-43.47	0.91
ESPRESSO	1200	32	2458860.0296351	26.28	1.19
ESPRESSO	1200	28	2458862.0141624	-93.56	1.98
ESPRESSO	1200	40	2458915.0544167	27.21	1.18
ESPRESSO	1200	43	2458916.9562919	-98.02	1.56
ESPRESSO	1200	39	2458917.8290392	-17.92	1.95