

1                   **Dynamic mechanical relaxation behavior**  
2                   **of binary metallic glasses**

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4           M.N. Liu<sup>a</sup>, Q. Hao<sup>a</sup>, J. Dong<sup>b</sup>, B.A. Sun<sup>b,c</sup>, S.D. Feng<sup>d</sup>, D. Crespo<sup>e</sup>, J.C. Qiao<sup>a,\*</sup>

5           <sup>a</sup>*School of Mechanics, Civil Engineering and Architecture, Northwestern Polytechnical*  
6           *University, Xi'an 710072, China*

7           <sup>b</sup>*Institute of Physics, Chinese Academy of Sciences, Beijing 100190, China*

8           <sup>c</sup>*Songshan Lake Materials Laboratory, Dongguan, Guangdong 523808, China*

9           <sup>d</sup>*State Key Laboratory of Metastable Materials Science and Technology, Yanshan*  
10           *University, Qinhuangdao 066004, China*

11           <sup>e</sup>*Departament de Física, Barcelona Research Center in Multiscale Science and*  
12           *Technology & Institut de Tècniques Energètiques, Universitat Politècnica de*  
13           *Catalunya, 08930-Barcelona, Spain*

14  
15           \*Corresponding author:

16           Dr. J.C. Qiao   E-mail address: qjczy@nwpu.edu.cn

1 **Abstract:** Secondary  $\beta$  relaxation is a fundamental issue to understand the diffusion,  
2 plasticity and glass transition behavior of metallic glasses. Binary metallic glasses  
3 provide a simple glassy system to probe the dynamic mechanical processes. In the  
4 current research, dynamic mechanical behavior of the typical binary metallic glasses  
5 was investigated by dynamic mechanical analysis (DMA). The dynamic mechanical  
6 relaxation behavior of  $\text{Cu}_{50-x}\text{Ti}_{50+x}$  ( $x= 0, 7$  and  $9$ ) and  $\text{Cu}_x\text{Zr}_{100-x}$  ( $x= 30, 50, 56, 61.8$   
7 and  $64$ ) binary metallic glass system was studied. It is found that the  $\beta$  relaxation  
8 becomes pronounced with the increase of copper, which is analyzed in the framework  
9 of the principle of mixing enthalpy. However, the  $\beta$  relaxation of CuZr or CuTi binary  
10 metallic glasses is much more modest than that of  $\text{Y}_{65}\text{Co}_{35}$ ,  $\text{Dy}_{65}\text{Co}_{35}$  and  $\text{Y}_{67}\text{Cu}_{33}$ .

11 **Keywords:** Binary metallic glass; Mechanical relaxation; Secondary relaxation;  
12 Structural heterogeneity

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26 **1. Introduction and background**

1 Metallic glasses (MGs) have drawn exhaustive interests in the past decades due to  
2 their unique physical and mechanical properties, such as high strength [1], excellent  
3 corrosion resistance [2], super thermo-plastic ability [3, 4] and superior catalytic  
4 capabilities [5]. These physical or mechanical properties are closely related to their  
5 heterogeneous microstructure. Different than conventional crystalline solids, there are  
6 no recognizable defects (i.e., dislocations and grain boundaries) in metallic glasses. It  
7 is very important to establish the correlation between the microstructure and  
8 mechanical/physical properties of MGs [6, 7]. The link between mechanical relaxation  
9 processes, in particular, the secondary  $\beta$  relaxation and mechanical properties (i.e.  
10 plasticity) has been established [8-10]. Therefore, MGs are considered to be the model  
11 system for probing the  $\beta$  relaxation of glassy materials.

12 It is believed that glassy materials experience two relaxation processes. The main  
13  $\alpha$  relaxation emerges at high temperature and becomes "frozen" when the temperature  
14 lies below the glass transition temperature  $T_g$ . On the other hand, the secondary  $\beta$   
15 relaxation remains active below the  $T_g$ , decoupled from the main  $\alpha$  relaxation. The  
16 primary  $\alpha$  relaxation relies on the cooperative movement of atoms, while the secondary  
17  $\beta$  relaxation is associated to the local motion of atoms. In essence, the main  $\alpha$  relaxation  
18 of non-crystalline solids is linked to the dynamic glass transition process. **According to**  
19 **the recent investigations [11-15], this is well in accordance with the cooperative shear**  
20 **model based on the concept of a potential energy landscape. While the slow  $\beta$  process**  
21 **corresponds to individual shear transformation zone (STZ) transitions, the  $\alpha$  relaxation**  
22 **can be related to their percolation, leading to macroscopic plastic flow. Nevertheless,**  
23 **of more significant importance is the identification of the universal fast  $\beta'$  secondary**  
24 **relaxation in addition to the slow  $\beta$  and  $\alpha$  relaxation for MGs [15, 16]. Mechanical**  
25 **spectroscopy data shows that the two secondary relaxations in MGs have quite different**  
26 **dynamic characteristics [13, 17]. The peak of the loss modulus of fast  $\beta'$  relaxation has**  
27 **much lower magnitude but larger broadness than that of the slow  $\beta$  relaxation. In the**  
28 **past two decades, a large corpus of experimental data has disclosed that  $\beta$  relaxation is**  
29 **key to describe the properties of glass-forming materials [18]. While there is only one**

1 relaxation process at high temperature for liquid, the secondary relaxation process in  
2 the glass is commonly associated with the breakdown of the Stokes-Einstein relation.  
3 A correlation between atomic mobility and the local structure has been proved in  
4 supercooled liquids [19].

5 Binary metallic glasses are excellent model alloys to study the mechanical  
6 relaxations. Theoretical research [20] and simulations [21] have proven that the  $\beta$   
7 relaxation is closely linked to the diffusion of atoms in MGs. It has been proved that  
8 the activation energy of the  $\beta$  relaxation  $E_\beta$  is related to the diffusion activation energy  
9 in the Fe/Co/Ni MG system [22]. In addition, it was further confirmed that the self-  
10 diffusion of Fe/Co/Ni plays an important role in the  $\beta$  relaxation process. Previous  
11 investigations have proved that the glass transition ( $\alpha$  relaxation) behavior is mainly  
12 dominated by the slowing down of diffusion of large particles, while the diffusion of  
13 small particles remains in a deep glassy state [23-25]. These results rise one unsolved  
14 question: What is the link between the diffusion and the relaxation modes in glassy  
15 materials?

16 In the current work, the mechanical relaxation behavior of the binary MGs, i.e.,  
17  $\text{Cu}_x\text{Zr}_{100-x}$  ( $x= 30, 50, 56, 61.8$  and  $64$ ),  $\text{Cu}_{50-x}\text{Ti}_{50+x}$  ( $x= 0, 7$  and  $9$ ) and rare earth based  
18 (i.e.  $\text{Y}_{65}\text{Co}_{35}$ ,  $\text{Dy}_{65}\text{Co}_{35}$  and  $\text{Y}_{67}\text{Cu}_{33}$ ) was probed by dynamic mechanical analysis  
19 (DMA). The results suggested that the slow  $\beta$  relaxation . The current research provides  
20 a new way to understand the origin of the slow  $\beta$  relaxation.

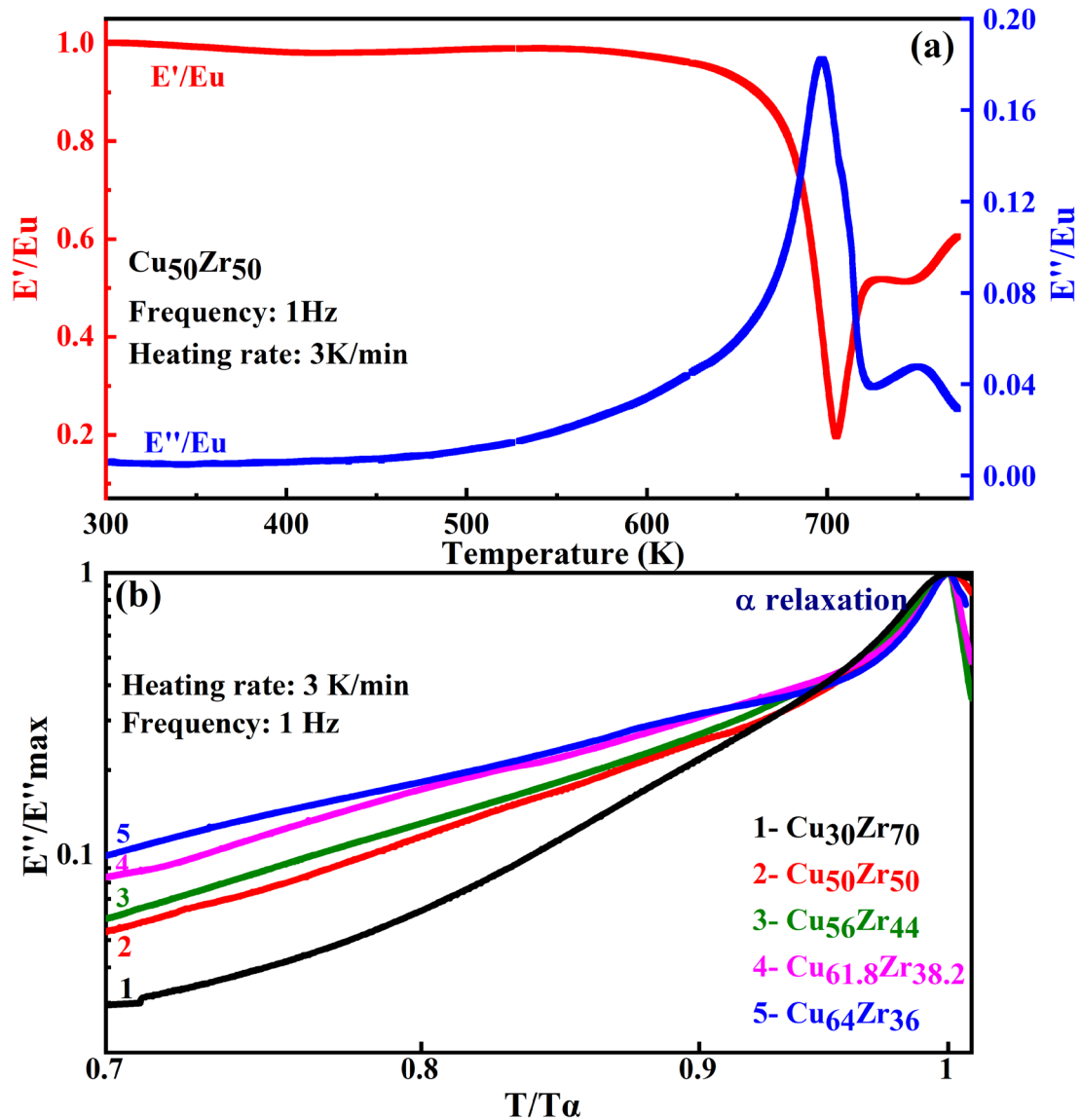
## 21 **2. Experimental procedure**

22 The ingots of the model alloys were prepared by arc melting in a Ti-gettered Argon  
23 atmosphere. The ingots were re-melted at least 5 times to keep the chemical  
24 homogeneity. Ribbons were obtained by single-roll melt-spinning in an Argon  
25 atmosphere with a tangential velocity of about 40 m/s. The thickness of the ribbons is  
26 around 30-40  $\mu\text{m}$ . Dynamical mechanical behavior of the MGs was carried out on a  
27 commercial DMA (TA Q800) in film tensile mode in a nitrogen-flushed atmosphere.

## 28 **3. Results and discussions**

1        **Fig.1 (a)** shows the temperature dependence of the normalized storage modulus  
2 (E'/E<sub>u</sub>) and the loss modulus (E''/E<sub>u</sub>) of a Cu<sub>50</sub>Zr<sub>50</sub> binary MG (driving frequency is 1  
3 Hz, the heating rate is 3 K/min). E<sub>u</sub> corresponds to the value of unrelaxed modulus at  
4 ambient temperature. It can be seen that in the low temperature range the model alloy  
5 is staying in amorphous state. The storage modulus E' remains approximately constant  
6 with a high value, while the loss modulus E'' is very low. Consequently, elastic  
7 deformation dominates the mechanical behavior in this temperature domain. By  
8 increasing the temperature, the storage modulus decreases drastically while the loss  
9 modulus reaches its maximum value around 695 K, which is associated with the  
10 primary α relaxation. It is well recognized that the main α relaxation corresponds to the  
11 evolution from the out-of-equilibrium state to the supercooled liquid in glassy materials  
12 [17]. When the temperature is above 700 K, both E' and E'' increase exceedingly as  
13 temperature increases, as a consequence of crystallization. It should be stressed that  
14 pronounced secondary β relaxation has not been observed in the Cu<sub>50</sub>Zr<sub>50</sub> binary MG.

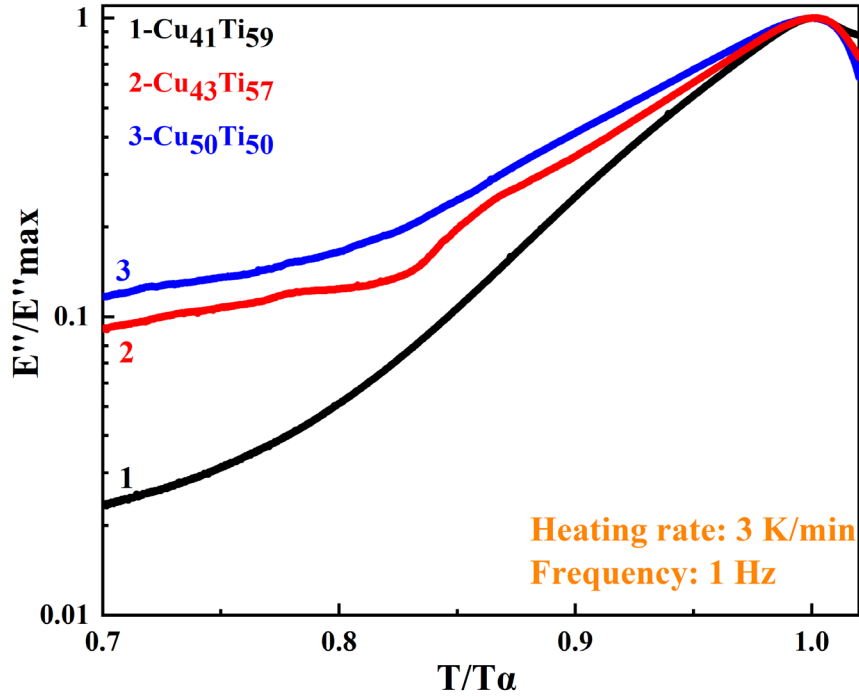
15        It is well known that the secondary β relaxation is linked to the localized  
16 movement of atoms while α relaxation is related to cooperative motion of atoms [26-  
17 28], and appears in form of shoulder or excess wing in the dynamic mechanical-loss  
18 spectra [29]. **Fig.1 (b)** presents the evolution of the normalized loss modulus with  
19 temperature in typical Cu<sub>x</sub>Zr<sub>100-x</sub> (x=30, 50, 56, 61.8 and 64) binary MGs. The data are  
20 normalized to the values of temperature and loss modulus at the peak of α relaxation,  
21 namely T<sub>α</sub> and E''<sub>max</sub>. As already found in the literature, the secondary β relaxation in  
22 CuZr allows displays as an excess wing rather than a discernible shoulder or peak.  
23 Interestingly, the secondary relaxation process is sensitive to the chemical composition.  
24 By increasing the content of zirconium, the excess wing due to β relaxation of the binary  
25 MGs was almost removed.



1  
 2 **Fig.1 (a)** Evolution of normalized storage modulus and loss modulus with the  
 3 temperature of  $Cu_{50}Zr_{50}$  metallic glass (heating rate is 3 K/min; driving frequency is 1  
 4 Hz). **(b)** The normalized loss modulus  $E''/E''_{max}$  as a function of temperature in  
 5  $Cu_xZr_{100-x}$  ( $x = 30, 50, 56, 61.8$  and  $64$ ) MGs. Measurement frequency is 1 Hz and  
 6 heating rate is 3 K/min.  $E''_{max}$  is the peak temperature of the main  $\alpha$  relaxation.

7 In order to further understand the effect of the chemical composition in binary

1 MGs, mechanical relaxation of  $\text{Cu}_{50-x}\text{Ti}_{50+x}$  ( $x=0, 7$  and  $9$ ) binary MGs was investigated.  
 2 **Fig.2** shows the temperature dependence of the normalized loss modulus with  
 3 temperature of  $\text{Cu}_{50-x}\text{Ti}_{50+x}$  ( $x=0, 7$  and  $9$ ) binary MGs. According to the mechanical-  
 4 loss spectra, the slow  $\beta$  relaxation was suppressed by increasing the content of titanium.



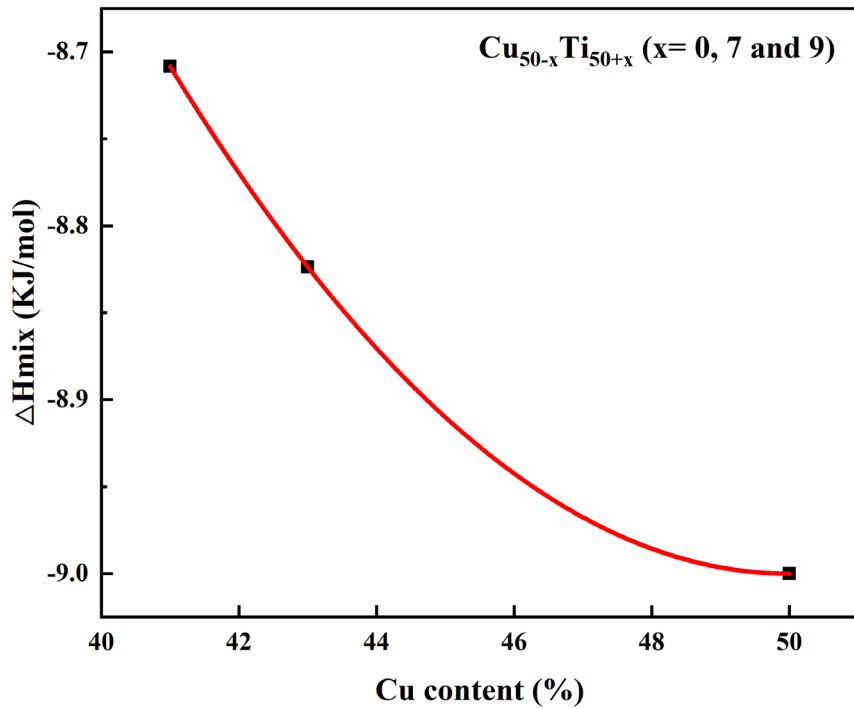
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 6 **Fig.2** Normalized loss modulus  $E''/E''_{\max}$  as a function of temperature in  $\text{Cu}_{50-x}\text{Ti}_{50+x}$   
 7 ( $x=0, 7$  and  $9$ ) MGs. Testing frequency is 1 Hz and heating rate is 3 K/min.

8 The behavior of  $\beta$  relaxation in MGs can be explained in the framework of mixing  
 9 enthalpy principle [30]. A negative enthalpy of mixing for a pair of atoms exhibits an  
 10 evident  $\beta$  relaxation, while a positive or a significantly different mixing enthalpy will  
 11 suppress  $\beta$  relaxation. The average mixing heat  $\Delta H^{chem}$  is given by

$$12 \quad \Delta H_{mix} = 4 \sum \Delta H_{ABmix} c_A c_B \quad (1)$$

13 where  $\Delta H_{ABmix}$  represents the enthalpy of mixing between elements A and B.  $c_A, c_B$   
 14 indicate the molar fractions of elements A and B in the alloy, respectively [31]. The data  
 15 of mixing enthalpy of constituent atomic pairs are derived from Ref. [32]. **Fig.3** plots

1 the mean chemical affinity  $\Delta H_{mix}$  of the  $\text{Cu}_{50-x}\text{Ti}_{50+x}$  ( $x=0, 7$  and  $9$ ) MGs. Cu-Ti  
2 atomic pair has a relatively negative enthalpy of mixing of  $-9$  kJ/mol. In the case of the  
3  $\text{Cu}_{50-x}\text{Ti}_{50+x}$  ( $x=0, 7$  and  $9$ ) MGs,  $\text{Cu}_{50}\text{Ti}_{50}$  has the largest negative enthalpy of mixing,  
4 which corresponds to the most prominent  $\beta$  relaxation peak. The experiments in the  
5 current research are in good agreement with the prediction of mixing enthalpy principle.



6

7 **Fig.3** Mixing enthalpy  $\Delta H_{mix}$  for the  $\text{Cu}_{50-x}\text{Ti}_{50+x}$  ( $x=0, 7$  and  $9$ ) MGs.

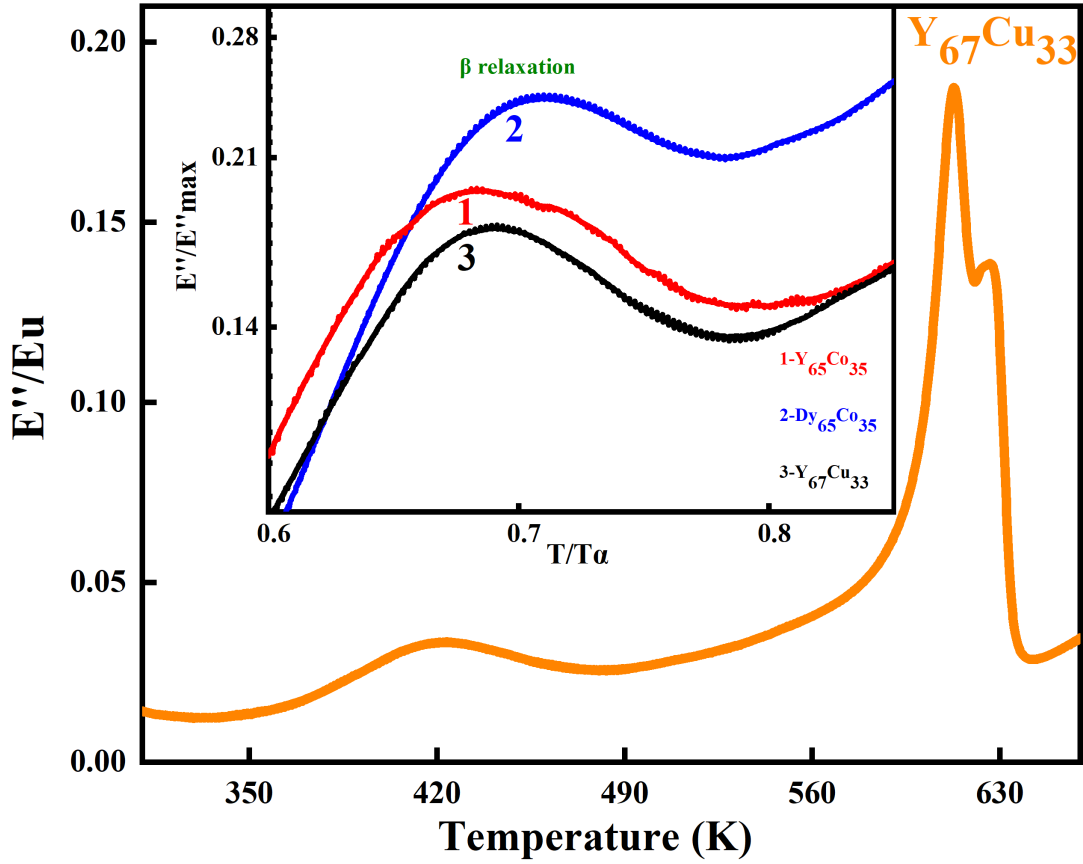
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9 Recent investigation on a LaNiAl MG suggests that the slow  $\beta$  relaxation is  
10 sensitive to the chemical compositions [33]. It has been found that binary La-Ni MGs  
11 shows an distinct slow  $\beta$  relaxation peak while La-Al and La-Cu binary MGs show  
12 excess wing process. Since both La and Y are rare earth (RE) elements, whether the  
13 unusual dynamics are general to similar MGs based on other RE elements. The question  
14 which remains open is whether the pronounced  $\beta$  relaxation peak is general for RE-  
15 Ni/Co binary MGs. The binary RE-Ni/Co MGs may serve as model materials to  
16 investigate the  $\beta$  relaxation in MGs due to its simple composition. In general, only a



1 few types of MGs exhibit significant  $\beta$  relaxation revealed by DMA measurements,  
2 such as La- [34] and Nd-based [28] MGs. For other MGs, i.e. Zr-[35], Cu-[36], Ti-[37],  
3 Mg-[38] and Ce-based [39] MGs show an excess wing in their loss modulus. Most  
4 binary MGs exhibit weak evidence of  $\beta$  relaxation in their dynamic mechanical  
5 spectroscopy. In contrast to other MGs,  $\text{Co}_{35}\text{Dy}_{65}$  and  $\text{Co}_{35}\text{Y}_{65}$  show a distinct  $\beta$   
6 relaxation peak in the mechanical relaxation measurements.

7 As shown in **Fig.4**, the temperature dependence loss modulus  $E''$  of a  $\text{Y}_{67}\text{Cu}_{33}$   
8 binary MG exhibits two distinct peaks : an  $\alpha$  relaxation peak of  $E''$  located at about  
9  $600\pm 5$  K and a pronounced  $\beta$  relaxation peak situated at about  $420\pm 5$  K. The  $\beta$   
10 relaxation peak of the loss modulus is associated to the heterogeneity of microstructure  
11 or the “defects” of MGs. The inset of **Fig.4** compares the  $\beta$  relaxation behavior of  
12  $\text{Dy}_{65}\text{Co}_{35}$ ,  $\text{Y}_{65}\text{Co}_{35}$  and  $\text{Y}_{67}\text{Cu}_{33}$  binary MGs. Only the curves with the testing frequency  
13 of 1Hz are exposed. The loss modulus  $E''$  curves have been scaled for comparison. The  
14 temperature is reduced by the peak temperature of  $\alpha$  relaxation  $T_\alpha$  and  $E''$  is normalized  
15 by  $E''_\alpha$ . From the normalized plot, it can be realized that the  $\beta$  relaxation of  $\text{Y}_{65}\text{Co}_{35}$  is  
16 stronger than that of  $\text{Y}_{67}\text{Cu}_{33}$ . In addition, the intensity  $\beta$  relaxation of  $\text{Dy}_{65}\text{Co}_{35}$  is the  
17 largest, which further confirms that the magnetic element Co can facilitate pronounced  
18  $\beta$  relaxation.



1

2 **Fig.4** Temperature dependence of the loss modulus  $E''$  of  $Y_{67}Cu_{33}$  at 1 Hz and a heating  
 3 rate is 3 K/min. The inset shows the temperature dependence of normalized loss  
 4 modulus in  $Y_{65}Co_{35}$  ,  $Dy_{65}Co_{35}$  and  $Y_{67}Cu_{33}$  binary MGs (heating rate is 3 K/min and  
 5 driving frequency is 1 Hz).

6

7 **Fig.5** illustrates the temperature dependence of  $E''$  at different frequencies (0.5, 1,  
 8 2, 4 or 8 Hz) of  $Dy_{65}Co_{35}$  and  $Y_{65}Co_{35}$  binary MGs. The  $T_{\beta}$  shifts to higher temperatures  
 9 by increasing the frequency. The activation energy for  $\beta$  relaxation ( $E_{\beta}$ ) is determined  
 10 by the  $\ln(f)$  vs  $1000/T_{\beta}$  plot. It can be fitted to an Arrhenius equation [40]:

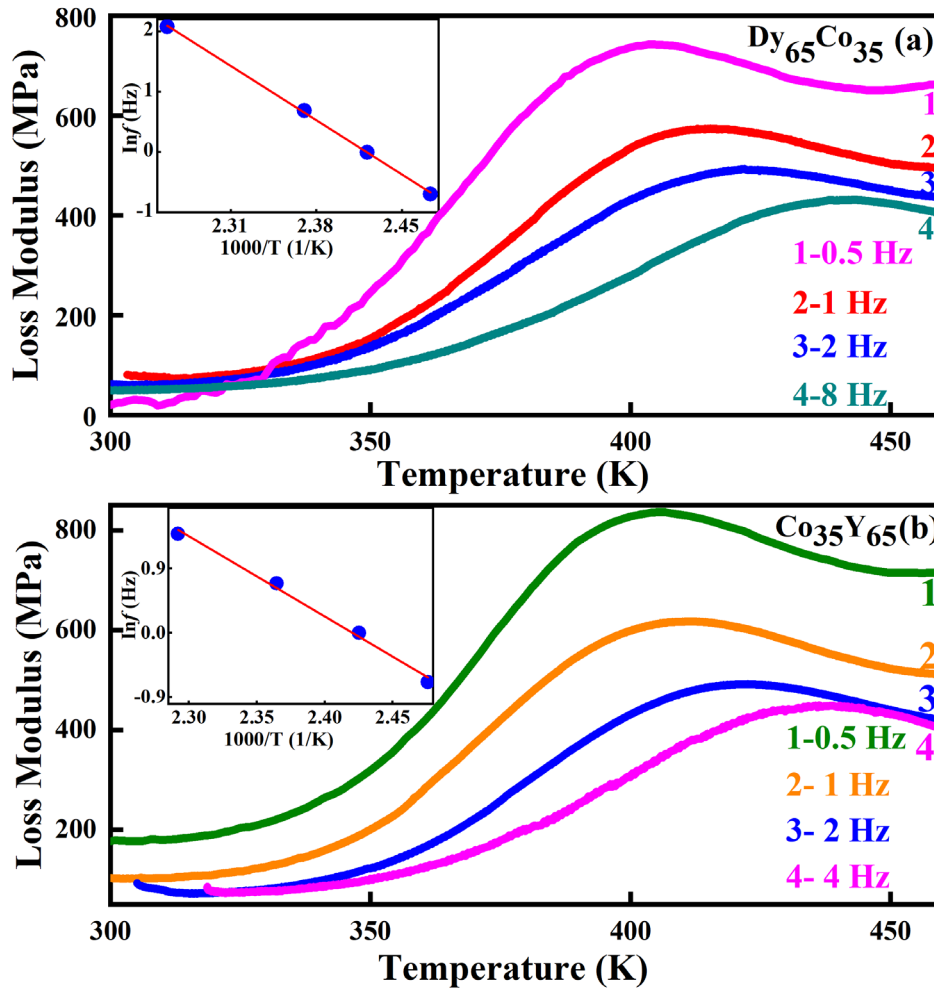
11

$$f = f_{\infty} \exp\left(-\frac{E_{\beta}}{RT}\right)$$

1        where  $f_{\infty}$  is the Arrhenius pre-factor,  $E_{\beta}$  is the activation energy of the  $\beta$   
2 relaxation and  $R$  is the gas constant.

3        The inset of **Fig.5 (a)** shows the Arrhenius fit of the measured data of Dy<sub>65</sub>Co<sub>35</sub>  
4 MG. The fitted activation energy  $E_{\beta}$  is 110 kJ/mol, or  $E_{\beta} = 23RT_g$  ( $R$  is the gas  
5 constant). In the same way, the inset of **Fig.5 (b)** shows the Arrhenius fit of the  
6 measured data of Y<sub>65</sub>Co<sub>35</sub>MGs. The fitted activation energy  $E_{\beta}$  is in this case 107  
7 kJ/mol, or  $E_{\beta} = 23RT_g$ . Both are in good agreement with the relationship reported  
8 between the  $E_{\beta}$  and  $T_g$  in different MGs [9].

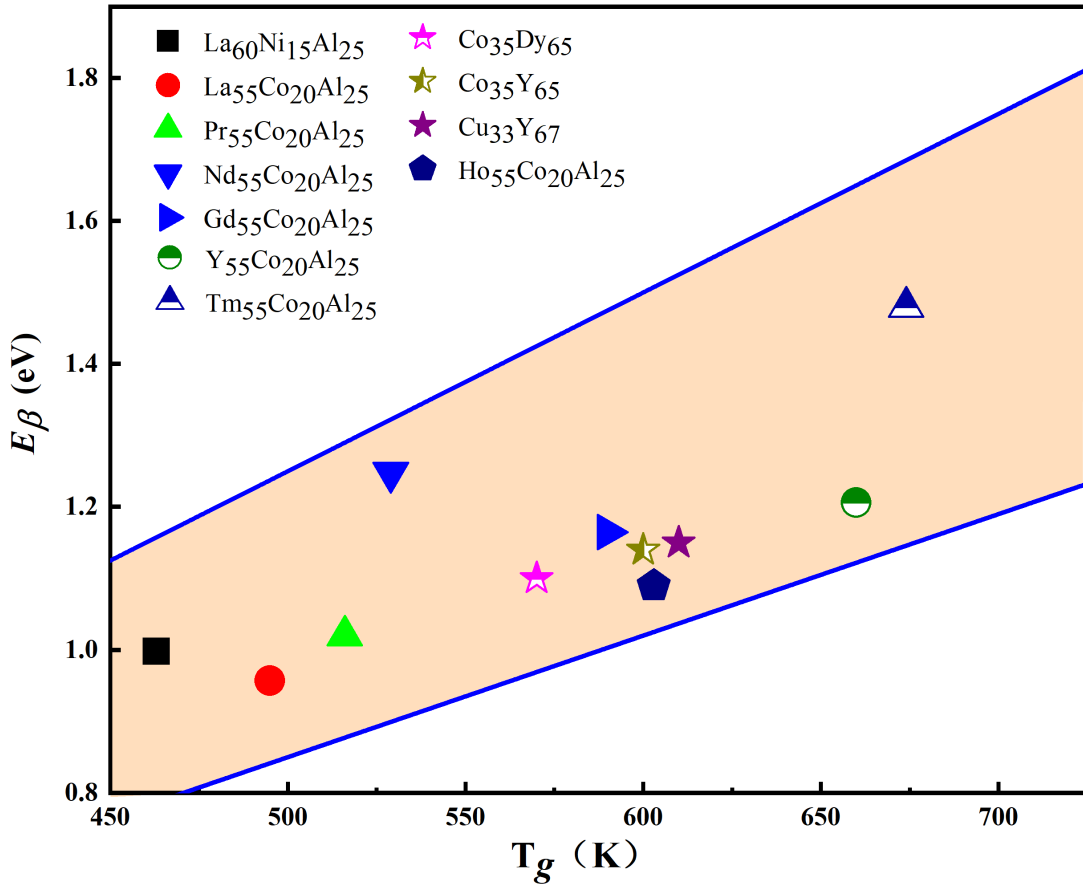
9        Compared to Dy<sub>65</sub>Co<sub>35</sub> MG, it is worth noting that Y<sub>65</sub>Co<sub>35</sub> MG presents an evident  
10 slow  $\beta$  relaxation at low temperature. Nevertheless, Y<sub>67</sub>Cu<sub>33</sub> MG don't show an  
11 apparent slow  $\beta$  relaxation. In fact, whether  $\beta$  relaxation is obvious will affect the  
12 mechanical properties of metallic glass. At room temperature, due to MGs are staying  
13 in the metastable state, namely, atoms are in "frozen" state. When the temperature  
14 increases (but still below the glass transition temperature  $T_g$ ), the "frozen" atoms soak  
15 up amounts of energy. The origin of  $\beta$  relaxation is not clear, but it is generally accepted  
16 that it is associated with string-like atomic rearrangements [41].



1  
2 **Fig.5 (a)** Temperature dependence of loss modulus of  $Dy_{65}Co_{35}$  metallic glass measured  
3 at various frequencies. Inset is the Arrhenius plot of the frequency vs peak temperature  
4 of the  $\beta$  relaxation. **(b)** Temperature dependence of loss modulus of  $Y_{65}Co_{35}$  metallic  
5 glass measured at different frequencies. Inset is the Arrhenius plot of the frequency vs  
6 peak temperature of the  $\beta$  relaxation.

7 **Fig.6** shows the correlation between the activation energy of  $\beta$  relaxation and the  
8 glass transition temperature  $T_g$  in typical MGs. It has been proven that activation energy  
9 of the slow  $\beta$  relaxation in glassy materials obeys the empirical relation  $E_{\beta}=(26 \pm 4)RT_g$   
10 ( $R$  is gas constant)[9]. As  $T_g$  is the characteristic temperature of  $\alpha$  relaxation, the  
11 relationship between  $E_{\beta}$  and  $T_g$  shows that  $\beta$  relaxation and  $\alpha$  relaxation are tightly  
12 connected. In other words,  $\alpha$  relaxation may consist of many percolating  $\beta$  relaxation

1 processes. In the case of  $\text{La}_{30}\text{Ce}_{30}\text{Al}_{15}\text{Co}_{25}$  [42] metallic glass, the activation energy of  
 2  $\alpha$  relaxation and  $\beta$  relaxation is 3.61 eV and 0.86 eV, respectively. It should be noted  
 3 that the apparent activation energy of  $\alpha$  relaxation is approximately four times that of  $\beta$   
 4 relaxation. The correlation between the activation energy of  $\text{Dy}_{65}\text{Co}_{35}/\text{Y}_{65}\text{Co}_{35}$  metallic  
 5 glasses and the glass transition temperature  $T_g$  obeys with the  $24RT_g$  relationship.  
 6



7  
 8 **Fig.6** Correlation between the activation energy of slow  $\beta$  relaxation  $E_\beta$  with the glass  
 9 transition temperature  $T_g$  in typical MGs. The data are obtained from Ref. [43].

10 **4. Conclusion**

11 In summary, mechanical relaxation processes of binary MGs were probed by  
 12 mechanical spectroscopy.  $\beta$  relaxation of the  $\text{Cu}_x\text{Zr}_{100-x}$  ( $x= 30, 50, 56, 61.8$  and  $64$ )  
 13 binary MG systems was suspended by increasing the content of zirconium. In addition,  
 14 in the case of the  $\text{Cu}_{50-x}\text{Ti}_{50+x}$  ( $x= 0, 7$  and  $9$ ) binary MGs, the slow  $\beta$  relaxation is

1 suppressed by increasing the content of titanium. Finally, rare earth-based binary MGs,  
2 i.e, Dy<sub>65</sub>Co<sub>35</sub>, Y<sub>65</sub>Co<sub>35</sub> ad Y<sub>67</sub>Cu<sub>33</sub>, show an evident slow  $\beta$  relaxation process. The  
3 activation energy of the  $\beta$  relaxation obeys the empirical relation:  $E_{\beta} \sim 26RT_g$ . The  $\beta$   
4 relaxation process is sensitive to the chemical composition of MGs. These binary MGs  
5 provide an ideal model alloy to shed light on the fundamental issues in MGs, i.e.  
6 relaxation processes, diffusiton behavior and mechanical properties.

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