

The relativistic spherical δ -shell interaction in \mathbb{R}^3 : spectrum and approximation

Albert Mas and Fabio Pizzichillo

*A. Mas, Departament de Matemàtiques i Informàtica,
Universitat de Barcelona. Gran Via Corts Catalanes 585, 08007 Barcelona (Spain).** and

*F. Pizzichillo, BCAM - Basque Center for Applied
Mathematics. Alameda de Mazarredo 14, 48009 Bilbao (Spain).†*

(Dated: August 24, 2017)

This note revolves on the free Dirac operator in \mathbb{R}^3 and its δ -shell interaction with electrostatic potentials supported on a sphere. On one hand, we characterize the eigenstates of those couplings by finding sharp constants and minimizers of some precise inequalities related to an uncertainty principle. On the other hand, we prove that the domains given by Dittrich, Exner and Šeba [*Dirac operators with a spherically symmetric δ -shell interaction, J. Math. Phys.* 30.12 (1989), 2875-2882] and by Arrizabalaga, Mas and Vega [*Shell interactions for Dirac operators, J. Math. Pures et Appl.* 102.4 (2014), 617-639] for the realization of an electrostatic spherical shell interaction coincide. Finally, we explore the spectral relation between the shell interaction and its approximation by short range potentials with shrinking support, improving previous results in the spherical case.

Keywords: Dirac operator, self-adjoint extension, spherical δ -shell interaction, singular integral, approximation by scaled regular potentials.

I. INTRODUCTION

Since 1929, when Paul Dirac formulated it, the “Dirac equation” has played a fundamental role in quantum mechanics attracting the attentions of both mathematicians and physicists. The newness of Dirac’s approach comes from the fact that he could describe the behaviour of the particles by using a 4-component spinor. With this approach one can observe the existence of negative-energy particles, see [17] and the references therein. By the way, it is not clear whether one should interpret the Dirac equation as a quantum mechanical evolution equation, like the Schrödinger equation for a single particle. An example that highlights this difference is, for instance, the comparison between the Schrödinger operator and the Dirac operator coupled with a singular potential supported on a hyper-surface: the so-called δ interaction.

The idea of coupling Hamiltonians with singular potentials supported on subsets of lower dimension with respect to the ambient space is quite classic in quantum mechanics. For the Schrödinger operator, this problem is described in the monograph [1] for finite and infinite δ -point interactions and in [9] for singular potentials supported on hyper-surfaces. Regarding the Dirac operator, in the 1-dimensional case the problem is well-understood. Thanks to [1, 10, 11] we get the description of

*Electronic address: albert.mas@ub.edu

†Electronic address: fpizzichillo@bcamath.org

the domain, some properties of the spectrum and a resolvent formula. In three dimensions the first result is [8]. By using the decomposition into spherical harmonics, they can reduce their analysis to a 1-dimensional question and they can construct the domain of the Dirac operator coupled with a singular potential supported on the sphere. In the case of a general surface Σ , the first work is [2]. In this work they characterize the domain of the δ -shell Dirac operator with coupling constant $\lambda \neq \pm 2$, by the interactions between certain functions $u \in H^1(\mathbb{R}^3)$ and $g \in L^2(\Sigma)$. Comparing this work with the general abstract theory given in [15], one could suppose that this kind of interaction is forcing g to be in $H^{1/2}(\Sigma)$. Indeed, recently, in [14] they proved that this conjecture is true. Moreover they also define the domain of δ -shell Dirac operator when the coupling constant $\lambda = \pm 2$. Finally, in [5, 7] they can define the domain of the δ -shell Dirac operator by using the abstract theory of boundary triples.

Nevertheless, one has to keep in mind that, even if this kind of model is easier to be mathematically understood, it is an ideal model that cannot be physically reproduced. This is the reason why it is interesting to approximate this kind of operators by more regular ones. And the results we obtain show the difference between the Schrödinger operator and the Dirac operator. For instance, in one dimension, if $V \in C_c^\infty(\mathbb{R})$ then

$$V_\epsilon(t) := \frac{1}{\epsilon} V\left(\frac{t}{\epsilon}\right) \rightarrow (\int V)\delta_0 \quad \text{when } \epsilon \rightarrow 0$$

in the sense of distributions, where δ_0 denotes the Dirac measure at the origin. In [1] it is proved that $\Delta + V_\epsilon \rightarrow \Delta + (\int V)\delta_0$ in the norm resolvent sense when $\epsilon \rightarrow 0$, and in [6] this result is generalized to higher dimensions for singular perturbations on general smooth hyper-surfaces.

These kind of results do not hold for the Dirac operator and it is physically known as ‘‘Klein’s paradox’’. In fact, in [16] it is proved that, in the 1-dimensional case, the convergence holds in the norm resolvent sense but the coupling constant does depend non-linearly on the potential V , unlike in the case of Schrödinger operators. Though, in the 3-dimensional case, the convergence holds in the strong resolvent sense for bounded smooth hyper-surfaces under certain hypothesis over the potential V , see [12].

This note revolves on the free Dirac operator in \mathbb{R}^3 and its δ -shell interactions with singular electrostatic potentials supported on a sphere. On one hand, we answer an open question posed in [3] which provides eigenstates of those couplings by finding sharp constants and minimizers of some precise inequalities related to an uncertainty principle (see Question III.4, Theorem III.6 and Corollary III.5). On the other hand, we prove that the domains given in [8] and [2] coincide in the spherical case and that the conjecture that comes from the comparison to [15] holds (see

Theorem II.2 and Remark II.3). Moreover we explore the spectral relation between the electrostatic δ -shell interaction and its approximation by the coupling of the free Dirac operator with shrinking short range potentials, improving the results in [12] in the case of a spherical shell interaction (see Theorem IV.2).

II. PRELIMINARIES

In this section we review some notation and basic rudiments on the construction of the electrostatic δ -shell interactions mentioned in the introduction. We first recall the approach presented in [2] for interactions on general smooth bounded domains and then we review the one from [8] in the case of a spherical interaction. At the end of the section we prove that both self-adjoint realizations coincide in the spherical case whenever the strength of the interaction differs from a concrete value.

Given $m \geq 0$, the free Dirac operator in \mathbb{R}^3 is defined by

$$H = -i\alpha \cdot \nabla + m\beta,$$

where $\alpha = (\alpha_1, \alpha_2, \alpha_3)$,

$$\alpha_i = \begin{pmatrix} 0 & \sigma_i \\ \sigma_i & 0 \end{pmatrix} \quad \text{for } i = 1, 2, 3, \quad \beta = \begin{pmatrix} \mathbb{I}_2 & 0 \\ 0 & -\mathbb{I}_2 \end{pmatrix}, \quad \mathbb{I}_2 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix},$$

$$\text{and } \sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

is the family of *Pauli's matrices*. Thus H acts on spinors of the form $\varphi : \mathbb{R}^3 \rightarrow \mathbb{C}^4$. Despite the massless case has its own interest, throughout this article we will assume that $m > 0$.

Let $\Omega \subset \mathbb{R}^3$ be a bounded smooth domain. We set $\Omega_+ = \Omega$ and $\Omega_- = \mathbb{R}^3 \setminus \overline{\Omega_+}$. Besides, let σ and ν denote the surface measure and unit outward (with respect to Ω) normal vector field on $\partial\Omega$, respectively. A fundamental solution of H is given by

$$\phi(x) = \frac{e^{-m|x|}}{4\pi|x|} \left(m\beta + (1 + m|x|) i\alpha \cdot \frac{x}{|x|^2} \right) \quad \text{for } x \in \mathbb{R}^3 \setminus \{0\},$$

see [2, Lemma 3.1]. Given $g \in L^2(\partial\Omega)^4$ we introduce the operator

$$\Phi(g)(x) = \int_{\partial\Omega} \phi(x-y)g(y) d\sigma(y) \quad \text{for } x \in \mathbb{R}^3 \setminus \partial\Omega. \quad (1)$$

By [2, Corollary 2.3], $\Phi : L^2(\partial\Omega)^4 \rightarrow L^2(\mathbb{R}^3)^4$ is linear and bounded. For $x \in \partial\Omega$ we also define

$$C_\sigma g(x) = \lim_{\epsilon \searrow 0} \int_{\partial\Omega \cap \{|x-y| > \epsilon\}} \phi(x-y)g(y) d\sigma(y) \quad \text{and} \quad C_\pm g(x) = \lim_{\Omega_\pm \ni y \rightarrow x} \Phi(g)(y), \quad (2)$$

where $\Omega_{\pm} \ni y \xrightarrow{nt} x$ means that y tends to x non-tangentially from Ω_{\pm} , respectively. From [2, Lemma 3.3] we know that both C_{σ} and C_{\pm} are linear and bounded in $L^2(\partial\Omega)^4$, and the following Plemelj-Sokhotski jump formulae holds:

$$C_{\pm} = \mp \frac{i}{2}(\alpha \cdot \nu) + C_{\sigma}. \quad (3)$$

Let $H^1(\mathbb{R}^3)^4$ denote the Sobolev space of \mathbb{C}^4 -valued functions such that all its components have all its zero and first order derivatives in $L^2(\mathbb{R}^3)$. It is well known that the trace operator on $\partial\Omega$, initially defined on smooth functions, extends to a bounded operator

$$\text{tr}_{\partial\Omega} : H^1(\mathbb{R}^3)^4 \rightarrow H^{1/2}(\partial\Omega)^4 \subset L^2(\partial\Omega)^4,$$

where $H^{1/2}(\partial\Omega)$ denotes the Sobolev-Slobodeckij space on $\partial\Omega$ of order 1/2.

We are ready to introduce the electrostatic δ -shell interaction $H + \lambda\delta_{\partial\Omega}$ studied in [2]. Here, $\lambda \in \mathbb{R}$ is a parameter that represents the strength of the interaction. Following [2, Theorem 3.8], we define

$$\begin{aligned} D(H + \lambda\delta_{\partial\Omega}) &= \{u + \Phi(g) : u \in H^1(\mathbb{R}^3)^4, g \in L^2(\partial\Omega)^4, \lambda \text{tr}_{\partial\Omega} u = -(1 + \lambda C_{\sigma})g\}, \\ (H + \lambda\delta_{\partial\Omega})\varphi &= H\varphi + \lambda \frac{\varphi_+ + \varphi_-}{2} \sigma \quad \text{for } \varphi \in D(H + \lambda\delta_{\partial\Omega}), \end{aligned} \quad (4)$$

where $H\varphi$ in the right hand side of the second statement in (4) is understood in the sense of distributions and φ_{\pm} denotes the boundary traces of φ when one approaches to $\partial\Omega$ from Ω_{\pm} . In particular, one has $(H + \lambda\delta_{\partial\Omega})\varphi = Hu \in L^2(\mathbb{R}^3)^4$ for all $\varphi = u + \Phi(g) \in D(H + \lambda\delta_{\partial\Omega})$. Furthermore, $H + \lambda\delta_{\partial\Omega}$ is self-adjoint for all $\lambda \neq \pm 2$. For shortness sake, we put

$$H_{\lambda} = H + \lambda\delta_{\partial\Omega}.$$

Roughly speaking, to construct the interaction H_{λ} presented in [2], one finds the domain of definition of the operator looking in the space of distributions and imposing some restrictions to get a big enough domain for the adjoint. However, in [8], the electrostatic δ -shell interactions are obtained using self-adjoint extensions of the restricted operator $H|_{C_c^{\infty}(\mathbb{R}^3 \setminus \partial\Omega)^4}$, as we will see below.

Unless we say the contrary, from now on we restrict our study to the case

$$\Omega = \{x \in \mathbb{R}^3 : |x| < 1\}.$$

For clarity, let us denote $B_{\pm} = \Omega_{\pm}$ and $\mathbb{S}^2 = \partial\Omega$. We now review the approach from [8], where the authors construct self-adjoint and rotationally invariant extensions of $H|_{C_c^{\infty}(\mathbb{R}^3 \setminus \mathbb{S}^2)^4}$ by using the decomposition in the classical spherical harmonics Y_n^l . The indices refer to $n = 0, 1, 2, \dots$

and $l = -n, -n + 1, \dots, n$, and the functions satisfy $\Delta_{\mathbb{S}^2} Y_n^l = n(n+1)Y_n^l$, where $\Delta_{\mathbb{S}^2}$ denotes the usual spherical laplacian. Moreover, $\{Y_n^l\}_{l,n}$ is a complete orthonormal set in $L^2(\mathbb{S}^2)$. For $j = 1/2, 3/2, 5/2, \dots$ and $m_j = -j, -j + 1, \dots, j$ set

$$\begin{aligned}\psi_{j-1/2}^{m_j} &= \frac{1}{\sqrt{2j}} \begin{pmatrix} \sqrt{j+m_j} Y_{j-1/2}^{m_j-1/2} \\ \sqrt{j-m_j} Y_{j-1/2}^{m_j+1/2} \end{pmatrix}, \\ \psi_{j+1/2}^{m_j} &= \frac{1}{\sqrt{2j+2}} \begin{pmatrix} \sqrt{j+1-m_j} Y_{j+1/2}^{m_j-1/2} \\ -\sqrt{j+1+m_j} Y_{j+1/2}^{m_j+1/2} \end{pmatrix}.\end{aligned}$$

Then $\{\psi_{j\pm 1/2}^{m_j}\}_{j,m_j}$ is a complete orthonormal set in $L^2(\mathbb{S}^2)^2$. Moreover, if we set

$$r = |x|, \quad \hat{x} = x/|x| \quad \text{and} \quad L = -ix \times \nabla \quad \text{for } x \in \mathbb{R}^3 \setminus \{0\},$$

then

$$(\sigma \cdot \hat{x})\psi_{j\pm 1/2}^{m_j} = \psi_{j\mp 1/2}^{m_j}, \quad \text{and} \quad (1 + \sigma \cdot L)\psi_{j\pm 1/2}^{m_j} = \pm(j+1/2)\psi_{j\pm 1/2}^{m_j},$$

where $\sigma = (\sigma_1, \sigma_2, \sigma_3)$ is the vector of *Pauli's matrices*. For $k_j = \pm(j+1/2)$ we define

$$\Phi_{m_j, k_j}^+ = \begin{pmatrix} i\psi_{j\pm 1/2}^{m_j} \\ 0 \end{pmatrix} \quad \text{and} \quad \Phi_{m_j, k_j}^- = \begin{pmatrix} 0 \\ \psi_{j\mp 1/2}^{m_j} \end{pmatrix}.$$

Then, the set $\mathcal{B} = \{\Phi_{m_j, k_j}^+, \Phi_{m_j, k_j}^-\}_{j, k_j, m_j}$ is a complete orthonormal base of $L^2(\mathbb{S}^2)^4$. We refer to [17, Section 4.6] for the details.

We now introduce the subspaces

$$\mathcal{C}_{m_j, k_j} = \left\{ \frac{1}{r} \left(f_{m_j, k_j}^+(r) \Phi_{m_j, k_j}^+(\hat{x}) + f_{m_j, k_j}^-(r) \Phi_{m_j, k_j}^-(\hat{x}) \right) \in L^2(\mathbb{R}^3)^4 : f_{m_j, k_j}^\pm \in C_c^\infty(0, +\infty) \right\}.$$

From [17, Theorems 1.1 and 4.14] we know that the operator

$$\mathring{H} = H|_{C_c^\infty(\mathbb{R}^3 \setminus \{0\})^4}$$

is essentially self-adjoint and leaves the partial wave subspace \mathcal{C}_{m_j, k_j} invariant. Its action on each subspace is represented, in terms of the base \mathcal{B} , by the operator

$$D(\mathring{h}_{m_j, k_j}) = C_c^\infty(0, +\infty)^2, \quad \mathring{h}_{m_j, k_j} \begin{pmatrix} f \\ g \end{pmatrix} = \begin{pmatrix} m & -\frac{d}{dr} + \frac{k_j}{r} \\ \frac{d}{dr} + \frac{k_j}{r} & -m \end{pmatrix} \begin{pmatrix} f \\ g \end{pmatrix}. \quad (5)$$

By [17, Lemma 4.15], \mathring{h}_{m_j, k_j} is essentially self-adjoint and, if we set $h_{m_j, k_j} = \overline{\mathring{h}_{m_j, k_j}}$, we get that

$$H \cong \bigoplus_{j=\frac{1}{2}, \frac{3}{2}, \dots}^{\infty} \bigoplus_{m_j=-j}^j \bigoplus_{k_j=\pm(j+1/2)} h_{m_j, k_j},$$

where “ \cong ” means that the operators are unitarily equivalent and H is the free Dirac operator defined on $H^1(\mathbb{R}^3)$.

For m_j and k_j as above, let us define

$$D(\hat{h}_{m_j, k_j}) = C_c^\infty((0, 1) \cup (1, +\infty))^2 \subset D(\mathring{h}_{m_j, k_j}), \quad \hat{h}_{m_j, k_j} \varphi := \mathring{h}_{m_j, k_j} \varphi, \text{ for all } \varphi \in D(\hat{h}_{m_j, k_j}).$$

Assume that $\lambda \in \mathbb{R} \setminus \{\pm 2\}$ and set

$$M_\lambda^\pm = \begin{pmatrix} \lambda/2 & \pm 1 \\ \mp 1 & \lambda/2 \end{pmatrix}.$$

In [8] it is proved that the operator $h(\lambda)_{m_j, k_j}$ defined by

$$D(h(\lambda)_{m_j, k_j}) = \left\{ \varphi \in L^2(0, +\infty)^2 : h_{m_j, k_j} \varphi \in L^2(0, +\infty)^2, \varphi \in AC((0, 1) \cup (1, +\infty))^2, \right. \\ \left. M_\lambda^- \varphi(1^+) + M_\lambda^+ \varphi(1^-) = 0 \right\}, \quad (6)$$

$$h(\lambda)_{m_j, k_j} \varphi = h_{m_j, k_j} \varphi \quad \text{for all } \varphi \in D(h(\lambda)_{m_j, k_j})$$

is a self-adjoint extension of \hat{h}_{m_j, k_j} . Here, $AC((0, 1) \cup (1, +\infty))$ denotes the space of absolutely continuous functions on the open set $(0, 1) \cup (1, +\infty)$. Furthermore, if one sets

$$\delta_1(\varphi) = \frac{\varphi(1^+) + \varphi(1^-)}{2},$$

then $h(\lambda)_{m_j, k_j} = \mathring{h}_{m_j, k_j} + \lambda \delta_1$ on $D(h(\lambda)_{m_j, k_j})$, with the understanding that here \mathring{h}_{m_j, k_j} just means the differential operator given by the matrix on the right hand side of (5) acting in the sense of distributions. Let us finally introduce the subspaces

$$\mathcal{H}(\lambda)_{m_j, k_j} = \left\{ \frac{1}{r} \left(f_{m_j, k_j}^+(r) \Phi_{m_j, k_j}^+(\hat{x}) + f_{m_j, k_j}^-(r) \Phi_{m_j, k_j}^-(\hat{x}) \right) \in L^2(\mathbb{R}^3)^4 : f_{m_j, k_j}^\pm \in D(h(\lambda)_{m_j, k_j}) \right\}.$$

The electrostatic δ -shell interaction with strength λ studied in [8] is given by

$$D(\hat{H}(\lambda)) = \bigoplus_{j=\frac{1}{2}, \frac{3}{2}, \dots}^{\infty} \bigoplus_{m_j=-j}^j \bigoplus_{k_j=\pm(j+1/2)} \mathcal{H}(\lambda)_{m_j, k_j}, \\ \hat{H}(\lambda) \cong \bigoplus_{j=\frac{1}{2}, \frac{3}{2}, \dots}^{\infty} \bigoplus_{m_j=-j}^j \bigoplus_{k_j=\pm(j+1/2)} h(\lambda)_{m_j, k_j}, \quad (7)$$

which is a self-adjoint operator.

In order to compare the notions of a δ -shell interaction given by (4) in the spherical case and (7), let us first prove an auxiliary result.

Lemma II.1. *Let $\Omega \subset \mathbb{R}^3$ be a bounded domain of class C^2 . Then*

$$H^1(\mathbb{R}^3 \setminus \partial\Omega)^4 = \{u + \Phi(g) : u \in H^1(\mathbb{R}^3)^4, g \in H^{1/2}(\partial\Omega)^4\}.$$

Proof. If $f = u + \Phi(g)$ for some $u \in H^1(\mathbb{R}^3)^4$ and $g \in H^{1/2}(\partial\Omega)^4$, by [11, Lemma 3.1] we have that $f \in H^1(\mathbb{R}^3 \setminus \partial\Omega)^4$.

Let us consider now $f \in H^1(\mathbb{R}^3 \setminus \partial\Omega)^4$. Since $f \in H^1(\Omega_\pm)^4$, by the trace theorem we also have $f_\pm \in H^{1/2}(\partial\Omega)^4$. Set

$$g := i(\alpha \cdot \nu)(f_+ - f_-) \in H^{1/2}(\partial\Omega)^4$$

and $u = f - \Phi(g)$. Once again, [11, Lemma 3.1] shows that $u \in H^1(\mathbb{R}^3 \setminus \partial\Omega)^4$. Moreover, by (3),

$$u_+ - u_- = f_+ - f_- - C_+g + C_-g = f_+ - f_- + i(\alpha \cdot \nu)g = 0,$$

thus $u_+ = u_-$ and u has a well defined boundary trace in $H^{1/2}(\partial\Omega)^4$. This implies that actually $u \in H^1(\mathbb{R}^3)^4$, and we are done since $f = u + \Phi(g)$. \square

Theorem II.2. *Assume that $\Omega = \{x \in \mathbb{R}^3 : |x| < 1\}$. For any $\lambda \in \mathbb{R} \setminus \{\pm 2\}$, the self-adjoint realizations H_λ and $\widehat{H}(\lambda)$ defined by (4) and (7), respectively, coincide.*

Proof. Consider the operator

$$\begin{aligned} D(\widetilde{H}_\lambda) &= \{u + \Phi(g) : u \in H^1(\mathbb{R}^3)^4, g \in H^{1/2}(\mathbb{S}^2)^4, \lambda \operatorname{tr}_{\mathbb{S}^2} u = -(1 + \lambda C_\sigma)g\}, \\ \widetilde{H}_\lambda &= H_\lambda|_{D(\widetilde{H}_\lambda)}. \end{aligned}$$

Since $H^{1/2}(\mathbb{S}^2)^4 \subset L^2(\mathbb{S}^2)^4$, by construction we get $\widetilde{H}_\lambda \subset H_\lambda$. We are going to prove that

$$\widehat{H}(\lambda) \subset \widetilde{H}_\lambda. \quad (8)$$

With this at hand, we deduce that $\widehat{H}(\lambda) \subset \widetilde{H}_\lambda \subset H_\lambda$ and, since both $\widehat{H}(\lambda)$ and H_λ are self-adjoint operators for $\lambda \neq \pm 2$, we finally conclude that $\widehat{H}(\lambda) = H_\lambda$ and the theorem follows. Let us focus on (8). Fixed m_j and k_j as in (7), for simplicity of notation we put

$$\begin{aligned} f^\pm(r) &= f_{j,m_j}^\pm(r), \\ \Phi^\pm(\hat{x}) &= \Phi_{j,m_j}^\pm(\hat{x}), \\ \mathcal{H}(\lambda) &= \mathcal{H}(\lambda)_{m_j,k_j}, \\ h(\lambda) &= h(\lambda)_{m_j,k_j}. \end{aligned}$$

Thus, any $\varphi \in \mathcal{H}(\lambda)$ can be written as

$$\varphi(x) = \frac{1}{r} (f^+(r)\Phi^+(\hat{x}) + f^-(r)\Phi^-(\hat{x})) = \frac{1}{r} \begin{pmatrix} f^+(r) \\ f^-(r) \end{pmatrix} \cdot \begin{pmatrix} \Phi^+(\hat{x}) \\ \Phi^-(\hat{x}) \end{pmatrix}.$$

In the last expression above, “ \cdot ” just means “scalar product”. As before, we denote by φ_\pm the boundary values of φ when we approach \mathbb{S}^2 from Ω_\pm . Let \mathcal{M}_λ^\pm be the operator defined on $\mathcal{H}(\lambda)$ by the action of the matrix M_λ^\pm on the basis $\{\Phi^+, \Phi^-\}$, that is, for any $\hat{x} \in \mathbb{S}^2$,

$$\begin{aligned} \mathcal{M}_\lambda^+ \varphi_+(\hat{x}) &= \frac{1}{r} \left(M_\lambda^+ \begin{pmatrix} f^+(1^+) \\ f^-(1^+) \end{pmatrix} \right) \cdot \begin{pmatrix} \Phi^+(\hat{x}) \\ \Phi^-(\hat{x}) \end{pmatrix}, \\ \mathcal{M}_\lambda^- \varphi_-(\hat{x}) &= \frac{1}{r} \left(M_\lambda^- \begin{pmatrix} f^+(1^-) \\ f^-(1^-) \end{pmatrix} \right) \cdot \begin{pmatrix} \Phi^+(\hat{x}) \\ \Phi^-(\hat{x}) \end{pmatrix}. \end{aligned}$$

So, in particular, we have that

$$\mathcal{M}_\lambda^+ \varphi_-(\hat{x}) + \mathcal{M}_\lambda^- \varphi_+(\hat{x}) = 0 \quad \text{for all } \hat{x} \in \mathbb{S}^2. \quad (9)$$

Moreover, since $\varphi \in H^1(\mathbb{R}^3 \setminus \mathbb{S}^2)^4$, using Lemma II.1 we can write $\varphi = u + \Phi(g)$ for some $u \in H^1(\mathbb{R}^3)^4$ and $g \in L^2(\mathbb{S}^2)^4$. Then, since $\nu(\hat{x}) = \hat{x}$ for all $\hat{x} \in \mathbb{S}^2$, using (3) we see that (9) is equivalent to

$$\begin{aligned} 0 &= (\mathcal{M}_\lambda^+ + \mathcal{M}_\lambda^-) \operatorname{tr}_{\mathbb{S}^2} u(\hat{x}) + (\mathcal{M}_\lambda^+ C_+ + \mathcal{M}_\lambda^- C_-) g(\hat{x}) \\ &= (\mathcal{M}_\lambda^+ + \mathcal{M}_\lambda^-) \operatorname{tr}_{\mathbb{S}^2} u(\hat{x}) + \frac{1}{2} (\mathcal{M}_\lambda^- - \mathcal{M}_\lambda^+) i(\alpha \cdot \hat{x}) g(\hat{x}) + (\mathcal{M}_\lambda^+ + \mathcal{M}_\lambda^-) C_\sigma g(\hat{x}). \end{aligned} \quad (10)$$

Since $M_\lambda^+ + M_\lambda^- = \lambda \mathbb{I}_2$, where \mathbb{I}_2 denotes the 2×2 identity matrix, we get that, for $\hat{x} \in \mathbb{S}^2$,

$$(\mathcal{M}_\lambda^+ + \mathcal{M}_\lambda^-) u(\hat{x}) = \lambda u(\hat{x}), \quad (11)$$

$$(\mathcal{M}_\lambda^+ + \mathcal{M}_\lambda^-) C_\sigma g(\hat{x}) = \lambda C_\sigma g(\hat{x}). \quad (12)$$

Note also that

$$\frac{1}{2} (M_\lambda^- - M_\lambda^+) = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix},$$

that is the matrix that represent the operator $-i(\alpha \cdot \hat{x})$ on the basis $\{\Phi^+, \Phi^-\}$ (see [17, Equation 4.123]). So

$$\frac{1}{2} (\mathcal{M}_\lambda^- - \mathcal{M}_\lambda^+) (i\alpha \cdot \hat{x}) g(\hat{x}) = g(\hat{x}) \quad (13)$$

for $\hat{x} \in \mathbb{S}^2$. Combining (11), (12) and (13), (10) becomes

$$0 = \lambda \operatorname{tr}_{\mathbb{S}^2} u + (1 + \lambda C_\sigma) g.$$

In conclusion, we have seen that if $\varphi \in \mathcal{H}(\lambda)$ then $\varphi \in H_\lambda$. Since these arguments are valid for any m_j and k_j , (8) follows. \square

Remark II.3. From the proof of Theorem II.2 we also see that if $\lambda \neq \pm 2$ then $\tilde{H}_\lambda = H_\lambda$, which means that the condition $\lambda \operatorname{tr}_{\mathbb{S}^2} u = -(1 + \lambda C_\sigma) g$ in (4) forces g to belong to $H^{1/2}(\mathbb{S}^2)^4$.

III. ON THE SPECTRUM FOR THE SPHERICAL δ -SHELL INTERACTION

In this section we answer affirmatively a question posed in [3, Section 4.2.3]. As commented there, this yields a relation between the eigenvalues in the gap $(-m, m)$ for the electrostatic spherical δ -shell interaction and the minimizers of some precise quadratic form inequality. Before going further, we must recall some rudiments from [3, Section 4]. Throughout this section, Ω denotes the unit ball and $\partial\Omega = \mathbb{S}^2$. Given $a \in [-m, m]$, set

$$k^a(x) = \frac{e^{-\sqrt{m^2 - a^2}|x|}}{4\pi|x|} \mathbb{I}_2 \quad \text{and} \quad w^a(x) = \frac{e^{-\sqrt{m^2 - a^2}|x|}}{4\pi|x|^3} \left(1 + \sqrt{m^2 - a^2}|x|\right) i\sigma \cdot x$$

for $x \in \mathbb{R}^3 \setminus \{0\}$. Given $f \in L^2(\sigma)^2$ and $x \in \mathbb{S}^2$, set

$$K^a f(x) = \int_{\mathbb{S}^2} k^a(x-z) f(z) d\sigma(z) \quad \text{and} \quad W^a f(x) = \lim_{\epsilon \searrow 0} \int_{\{|x-z| > \epsilon\} \cap \mathbb{S}^2} w^a(x-z) f(z) d\sigma(z).$$

Then

$$C_\sigma^a = \begin{pmatrix} (a+m)K^a & W^a \\ W^a & (a-m)K^a \end{pmatrix},$$

where C_σ^a is defined as C_σ in (2) replacing the kernel ϕ by

$$\phi^a(x) = \frac{e^{-\sqrt{m^2-a^2}|x|}}{4\pi|x|} \left(a + m\beta + \left(1 + \sqrt{m^2-a^2}|x|\right) i\alpha \cdot \frac{x}{|x|^2} \right) \quad \text{for } x \in \mathbb{R}^3 \setminus \{0\}.$$

The following corresponds to [3, Lemma 4.3].

Lemma III.1. *Given $a \in (-m, m)$, there exist positive numbers $d_{j\pm 1/2}$ and purely imaginary numbers $p_{j\pm 1/2}$ for all $j = 1/2, 3/2, 5/2, \dots$, and $m_j = -j, -j+1, \dots, j$, such that*

(i) $K^a \psi_{j\pm 1/2}^{m_j} = d_{j\pm 1/2} \psi_{j\pm 1/2}^{m_j}$ and $\lim_{j \rightarrow \infty} d_{j\pm 1/2} = 0$. Moreover,

$$0 \leq d_{j\pm 1/2} \leq d_0 = \frac{1 - e^{-2\sqrt{m^2-a^2}}}{2\sqrt{m^2-a^2}}.$$

(ii) $W^a \psi_{j\pm 1/2}^{m_j} = p_{j\pm 1/2} \psi_{j\mp 1/2}^{m_j}$ and $p_{j+1/2} = -p_{j-1/2}$. Moreover,

$$|p_{j\pm 1/2}|^2 = \frac{1}{4} - (m^2 - a^2)d_{j+1/2}d_{j-1/2} \geq \frac{1}{4} e^{-2\sqrt{m^2-a^2}} \left(2 - e^{-2\sqrt{m^2-a^2}} \right).$$

The following result allows us to construct eigenstates for H_λ from the eigenfunctions of K^a ; it corresponds to [3, Lemma 4.6].

Lemma III.2. *Let H_λ be as in (4). If $\lambda > 0$ and $a \in (-m, m)$ satisfy*

$$\frac{\lambda^2}{4} - ((m+a)d_{j\mp 1/2} - (m-a)d_{j\pm 1/2})\lambda = 1 \quad \text{for some } j, \quad (14)$$

then, for any m_j , $\psi_{j\pm 1/2}^{m_j}$ gives rise to an eigenfunction for H_λ with eigenvalue a .

Remark III.3. In Lemma III.2, the expression ‘‘gives rise to an eigenfunction’’ means that, if one defines

$$g = \begin{pmatrix} f \\ h \end{pmatrix} \in L^2(\mathbb{S}^2)^4, \quad \text{where } h = \psi_{j\pm 1/2}^{m_j} \quad \text{and} \quad f = -(1/\lambda + (a+m)K^a)^{-1}W^a h,$$

setting $\varphi = \phi * (a\Phi^a(g)) + \Phi(g)$ one gets that $H_\lambda\varphi = a\varphi$. Here, Φ^a is defined as Φ in (1) replacing ϕ by ϕ^a .

In [3, Question 4.7], the following question was raised:

Question III.4. *Let $d_{j\pm 1/2}$ be the coefficients given by Lemma III.1. Is it true that $d_{j+1/2}d_{j-1/2} < d_1d_0$ for all $j = 3/2, 5/2, 7/2, \dots$?*

Theorem III.6 answers it in the affirmative and, as commented at the end of [3, Section 4.2.3], it yields the following result related to Lemma III.2. We first recall the values of d_0 and d_1 from Lemma III.1 (computed in [3]) and a precise constant d_* that will appear below, see [3, equations (4.31), (4.32) and (4.39), respectively]:

$$\begin{aligned} d_0 &= \frac{1 - e^{-2\sqrt{m^2 - a^2}}}{2\sqrt{m^2 - a^2}}, \\ d_1 &= \frac{1}{2\sqrt{m^2 - a^2}} \left(1 - \frac{1}{m^2 - a^2} + \left(1 + \frac{1}{\sqrt{m^2 - a^2}} \right)^2 e^{-2\sqrt{m^2 - a^2}} \right), \\ d_* &= \frac{1}{2\sqrt{m^2 - a^2}} - \frac{1}{2} \left(1 + \frac{1}{\sqrt{m^2 - a^2}} \right) e^{-2\sqrt{m^2 - a^2}}. \end{aligned}$$

Corollary III.5. *Let $a \in (-m, m)$ and $\lambda > 0$. Then, for any $f \in L^2(\sigma)^2$,*

$$\begin{aligned} \int_{\mathbb{S}^2} |f|^2 d\sigma &\leq \frac{1/\lambda + (m+a)d_0}{2d_*^2} \int_{\mathbb{S}^2} (1/\lambda + (m+a)K^a)^{-1} W^a f \cdot \overline{W^a f} d\sigma \\ &\quad + \frac{1}{2(1/\lambda + (m+a)d_0)} \int_{\mathbb{S}^2} (1/\lambda + (m+a)K^a)(\sigma \cdot \nu) f \cdot \overline{(\sigma \cdot \nu) f} d\sigma. \end{aligned} \quad (15)$$

The equality in (15) is only attained at linear combinations of ψ_1^l for $l \in \{-1/2, 1/2\}$. If

$$\frac{\lambda^2}{4} - ((m+a)d_0 - (m-a)d_1)\lambda = 1 \quad (16)$$

then the minimizers of (15) give rise to eigenfunctions of H_λ . Besides, these conclusions also hold if we exchange the roles of d_0 and d_1 in (15) and (16) and we replace ψ_1^l by ψ_0^l (that is, we exchange the roles of $j+1/2$ and $j-1/2$ for $j=1/2$).

Theorem III.6. *Let $d_{j\pm 1/2}$ be the coefficients given by Lemma III.1. Then,*

$$d_{j\pm 1/2} = I_{(j+1/2)\pm 1/2}(\sqrt{m^2 - a^2}) K_{(j+1/2)\pm 1/2}(\sqrt{m^2 - a^2}), \quad (17)$$

where I and K denote the standard second order Bessel's functions. Moreover,

$$d_{j+1/2}d_{j-1/2} < d_0d_1 \quad \text{for all } j = 3/2, 5/2, 7/2, \dots \quad (18)$$

Proof. Let us first compute $d_{j\pm 1/2}$ in terms of Bessel's functions. Fixed m_j and k_j , due to [17, Lemma 4.15] and Theorem II.2 it is enough to find some $a \in (-m, m)$ which is an eigenvalue for the operator $h(\lambda)_{m_j, k_j}$. We want to find some

$$\begin{pmatrix} f \\ g \end{pmatrix} \in D(h(\lambda)_{m_j, k_j})$$

verifying the following system of differential equations:

$$\begin{cases} (m-a)f + \left(-\frac{d}{dr} + \frac{k_j}{r}\right)g = 0, \\ \left(\frac{d}{dr} + \frac{k_j}{r}\right)f - (m+a)g = 0. \end{cases} \quad (19)$$

Set $M = \sqrt{m^2 - a^2}$. Since $k_j = \pm(j+1/2)$, we set

$$f(r) = \begin{cases} A\sqrt{r} I_{(j+1/2)\pm 1/2}(Mr) & \text{if } r < 1 \\ B\sqrt{r} K_{(j+1/2)\pm 1/2}(Mr) & \text{if } r > 1 \end{cases}, \quad g(r) = \begin{cases} \frac{AM}{m+a}\sqrt{r} I_{(j+1/2)\pm 1/2}(Mr) & \text{if } r < 1 \\ -\frac{BM}{m+a}\sqrt{r} K_{(j+1/2)\pm 1/2}(Mr) & \text{if } r > 1 \end{cases}, \quad (20)$$

for some $(A, B) \neq (0, 0)$. If we put

$$\varphi = \begin{pmatrix} f \\ g \end{pmatrix},$$

then $\varphi \in L^2(0, +\infty)^2$, $h_{m_j, k_j} \varphi \in L^2(0, +\infty)^2$, $\varphi \in AC((0, 1) \cup (1, +\infty))^2$ and φ satisfies (19). Thus, to get that φ is an eigenvector for the operator $h(\lambda)_{m_j, k_j}$ it remains to prove that $\varphi \in D(h(\lambda)_{m_j, k_j})$, that is we have to show that $M_\lambda^- \varphi(1^+) + M_\lambda^+ \varphi(1^-) = 0$. In other words, the following linear system must hold:

$$\begin{cases} A \left(\frac{M}{a+m} I_{(j+1/2)\mp 1/2}(M) + \frac{\lambda}{2} I_{(j+1/2)\pm 1/2}(M) \right) \\ \quad + B \left(\frac{M}{a+m} K_{(j+1/2)\mp 1/2}(M) + \frac{\lambda}{2} K_{(j+1/2)\pm 1/2}(M) \right) = 0, \\ A \left(\frac{\lambda M}{2(a+m)} I_{(j+1/2)\mp 1/2}(M) - I_{(j+1/2)\pm 1/2}(M) \right) \\ \quad + B \left(K_{(j+1/2)\pm 1/2}(M) - \frac{\lambda M}{2(a+m)} K_{(j+1/2)\mp 1/2}(M) \right) = 0. \end{cases}$$

Since this is a 2×2 homogeneous linear system on A and B and we are supposing that $(A, B) \neq (0, 0)$, we deduce that the associated matrix has null determinant. This means that

$$\begin{aligned} 0 &= -\frac{\lambda^2 M}{4(a+m)} \left(I_{(j+1/2)\pm 1/2}(M) K_{(j+1/2)\mp 1/2}(M) + I_{(j+1/2)\mp 1/2}(M) K_{(j+1/2)\pm 1/2}(M) \right) \\ &\quad + \frac{\lambda}{m+a} \left((m+a) I_{(j+1/2)\pm 1/2}(M) K_{(j+1/2)\pm 1/2}(M) - (m-a) I_{(j+1/2)\mp 1/2}(M) K_{(j+1/2)\mp 1/2}(M) \right) \\ &\quad + \frac{M}{m+a} \left(I_{(j+1/2)\pm 1/2}(M) K_{(j+1/2)\mp 1/2}(M) + I_{(j+1/2)\mp 1/2}(M) K_{(j+1/2)\pm 1/2}(M) \right). \end{aligned} \tag{21}$$

By [13, Equation 10.20.2] we get that

$$I_{(j+1/2)\pm 1/2}(M) K_{(j+1/2)\mp 1/2}(M) + I_{(j+1/2)\mp 1/2}(M) K_{(j+1/2)\pm 1/2}(M) = \frac{1}{M}. \tag{22}$$

Finally, combining (21) and (22) we see that the following must hold:

$$\begin{aligned} \frac{\lambda^2}{4} - \left((m+a) I_{(j+1/2)\pm 1/2}(M) K_{(j+1/2)\pm 1/2}(M) \right. \\ \left. - (m-a) I_{(j+1/2)\mp 1/2}(M) K_{(j+1/2)\mp 1/2}(M) \right) \lambda - 1 = 0. \end{aligned} \tag{23}$$

In conclusion, if we define

$$\begin{aligned} D_{j\pm 1/2}(a, \lambda) &= \frac{\lambda^2}{4} - \left((m+a) I_{(j+1/2)\pm 1/2}(M) K_{(j+1/2)\pm 1/2}(M) \right. \\ &\quad \left. - (m-a) I_{(j+1/2)\mp 1/2}(M) K_{(j+1/2)\mp 1/2}(M) \right) \lambda - 1 \end{aligned} \tag{24}$$

and we take $\varphi = \begin{pmatrix} f \\ g \end{pmatrix}$ with f and g given by (20), then φ is an eigenfunction for $h(\lambda)_{m_j, k_j}$ with eigenvalue a if and only if $D_{j\pm 1/2}(a, \lambda) = 0$. In this case the function

$$\psi(x) = \frac{1}{r} \left(f(r) \Phi_{m_j, k_j}^+(\hat{x}) + g(r) \Phi_{m_j, k_j}^-(\hat{x}) \right)$$

is an eigenfunction for H_λ with eigenvalue a . For this reason, a comparison of (23) and (14) yields (17), as desired.

Let us finally prove (18). We put $n = j + 1/2 \in \mathbb{N}$. Since $j > 1/2$, we have $n > 1$. Then (18) is equivalent to

$$d_n d_{n-1} < d_0 d_1, \quad \text{for all } n \geq 2. \quad (25)$$

We are going to show (25) by induction. For $n = 2$, we have to check that $d_1 d_2 < d_0 d_1$, which is equivalent to $d_1(d_2 - d_0) < 0$. Since $d_n \geq 0$ for all $n > 1$, it is enough to show that $d_2 - d_0 < 0$. But, from (17) we easily get that

$$d_2 - d_0 = \frac{3(M^3 + 2M^2 + 3M + 3) \sinh(M) - 3M(M^2 + 3M + 3) \cosh(M)}{e^M M^5} < 0.$$

Let us now suppose that (25) holds for $n - 1$. Then, we can split

$$d_{n-1} d_n - d_0 d_1 = d_{n-1}(d_n - d_{n-2}) + d_{n-1} d_{n-2} - d_0 d_1.$$

On one hand, $d_{n-1} d_{n-2} - d_0 d_1 < 0$ by (25). On the other hand, $d_n - d_{n-2} \leq 0$ by [4, Theorem 2] and $d_{n-1} \geq 0$. Thus (25) holds for all $n \geq 2$. \square

IV. APPROXIMATION BY SHORT - RANGE POTENTIALS

In this section we investigate the spectral relation between the electrostatic δ -shell interaction on the boundary of a smooth domain and its approximation by the coupling of the Dirac operator with a short-range potential which depends on a parameter $\epsilon > 0$ in such a way that it shrinks to the boundary of the domain as $\epsilon \rightarrow 0$; see the definition of $H_{\lambda,\epsilon}$ below. From [12, Theorem 1.2] we know that if $a \in \sigma(H_\lambda)$, where here $\sigma(\cdot)$ denotes the spectrum, then there exists a sequence $\{a_\epsilon\}$ such that $a_\epsilon \in \sigma(H_{\mu,\epsilon})$ and $a_\epsilon \rightarrow a$ for $\epsilon \rightarrow 0$, where $\lambda = 2 \tan(\frac{\mu}{2})$. However, the vice-versa spectral implication may not hold in general. In this section we are going to show that the reverse does hold in the spherical case, that is, if $a_\epsilon \rightarrow a$ with $a_\epsilon \in \sigma(H_{\mu,\epsilon})$, then $a \in \sigma(H_\lambda)$ (see Theorem IV.2 below). In particular this means that, when passing to the limit, we don't lose any element of the spectrum for electrostatic interactions with potentials shrinking on \mathbb{S}^2 .

Given $\epsilon > 0$ and $x \in \mathbb{R}^3$, we define

$$V_\epsilon(x) = \frac{1}{2\epsilon} \chi_{(1-\epsilon, 1+\epsilon)}(|x|) \quad \text{and} \quad \mathbf{V}_\epsilon = V_\epsilon \mathbb{I}_4,$$

where \mathbb{I}_4 denotes the 4×4 identity matrix. For $\lambda \in \mathbb{R}$, we also introduce the operators

$$\begin{aligned} D(\mathring{H}_{\lambda,\epsilon}) &= C_c^\infty(\mathbb{R}^3 \setminus \{0\})^4 \quad \text{and} \quad \mathring{H}_{\lambda,\epsilon} = H + \lambda \mathbf{V}_\epsilon, \\ D(H_{\lambda,\epsilon}) &= H^1(\mathbb{R}^3)^4 \quad \text{and} \quad H_{\lambda,\epsilon} = H + \lambda \mathbf{V}_\epsilon. \end{aligned}$$

Since $|V_\epsilon| \leq \frac{1}{2\epsilon}$, $\dot{H}_{\lambda,\epsilon}$ is essentially self-adjoint and $H_{\lambda,\epsilon}$ is self-adjoint by [17, Theorem 4.2]. Moreover $\sigma_{ess}(H_{\lambda,\epsilon}) = \sigma_{ess}(H) = \sigma(H) = (-\infty, -m] \cup [m, +\infty)$. For this reason we are looking for some $a \in (-m, m)$ eigenvalue of $H_{\lambda,\epsilon}$.

Our aim is to find a precise relation between a , λ and ϵ , say $R_\epsilon(a, \lambda)$, which must hold in order to get an eigenfunction for $H + \lambda V_\epsilon$ with eigenvector a . Then, we will take the limit of $R_\epsilon(a, \lambda)$ for $\epsilon \rightarrow 0$ and we will compare the result to (14). To do so, we use the same approach developed in Section III.

Clearly, if $\lambda = 0$ we get that $H_{\lambda,\epsilon} = H$, i.e. we are not perturbing the free Hamiltonian H , thus we can exclude this case in our study. Assuming that $\lambda \neq 0$, we note that if a is an eigenvalue of $H_{\lambda,\epsilon}$ with eigenfunction $\psi = \begin{pmatrix} \phi \\ \chi \end{pmatrix}$ then $-a$ is an eigenvalue of $H_{-\lambda,\epsilon}$ with eigenfunction $\tilde{\psi} = \begin{pmatrix} -\chi \\ \phi \end{pmatrix}$. For this reason, from now on, we will further assume that $\lambda > 0$.

Observe that $\dot{H}_{\lambda,\epsilon}$ leaves the partial wave subspace \mathcal{C}_{m_j, k_j} invariant. Its action on each subspace is represented with respect to the basis $\{\Phi_{m_j, k_j}^+, \Phi_{m_j, k_j}^-\}$ by the operator

$$D(\dot{h}(\lambda, \epsilon)_{m_j, k_j}) = C_c^\infty(0, +\infty)^2, \quad (26)$$

$$\dot{h}(\lambda, \epsilon)_{m_j, k_j} \begin{pmatrix} f \\ g \end{pmatrix} = \begin{pmatrix} m + \frac{\lambda}{2\epsilon} \chi_{(1-\epsilon, 1+\epsilon)} & -\frac{d}{dr} + \frac{k_j}{r} \\ \frac{d}{dr} + \frac{k_j}{r} & -m + \frac{\lambda}{2\epsilon} \chi_{(1-\epsilon, 1+\epsilon)} \end{pmatrix} \begin{pmatrix} f \\ g \end{pmatrix}.$$

Since $\dot{H}_{\lambda,\epsilon}$ is essentially self-adjoint $\dot{h}(\lambda, \epsilon)_{m_j, k_j}$ too, thus setting $h(\lambda, \epsilon)_{m_j, k_j} = \overline{\dot{h}(\lambda, \epsilon)_{m_j, k_j}}$ we get

$$D(h(\lambda, \epsilon)_{m_j, k_j}) = \{\varphi \in L^2(0, +\infty)^2 : h(\lambda, \epsilon)_{m_j, k_j} \varphi \in L^2(0, +\infty)^2 \text{ and } \varphi \in AC(0, +\infty)^2\} \quad (27)$$

and the action of $h(\lambda, \epsilon)_{m_j, k_j}$ on its domain of definition is formally given by the right hand side of the second equation in (26). Moreover, $a \in (-m, m)$ is an eigenvalue for $H_{\lambda,\epsilon}$ if and only if a is an eigenvalue for $h(\lambda, \epsilon)_{m_j, k_j}$ for some $\{m_j, k_j\}$. For this reason, we want to solve

$$\begin{cases} (m-a)f + \left(-\frac{d}{dr} + \frac{k}{r}\right)g = 0 \\ \left(\frac{d}{dr} + \frac{k}{r}\right)f - (m+a)g = 0 \end{cases} \quad \text{if } 0 < r < 1-\epsilon \text{ or } r > 1+\epsilon,$$

$$\begin{cases} \left(m-a + \frac{\lambda}{2\epsilon}\right)f + \left(-\frac{d}{dr} + \frac{k}{r}\right)g = 0 \\ \left(\frac{d}{dr} + \frac{k}{r}\right)f - \left(m+a - \frac{\lambda}{2\epsilon}\right)g = 0 \end{cases} \quad \text{if } 1-\epsilon < r < 1+\epsilon$$

for $\begin{pmatrix} f \\ g \end{pmatrix} \in D(h(\lambda, \epsilon)_{m_j, k_j})$.

Since $k_j = \pm(j + 1/2)$, a non-trivial solution is given by

$$\begin{aligned}
 f(r) &= \begin{cases} A\sqrt{r} I_{(j+\frac{1}{2})\pm\frac{1}{2}}(Mr) & r < 1 - \epsilon \\ B_1\sqrt{r} J_{(j+\frac{1}{2})\pm\frac{1}{2}}(Lr) + B_2\sqrt{r} Y_{(j+\frac{1}{2})\pm\frac{1}{2}}(Lr) & 1 - \epsilon < r < 1 + \epsilon \\ C\sqrt{r} K_{(j+\frac{1}{2})\pm\frac{1}{2}}(Mr) & r > 1 + \epsilon \end{cases} \\
 g(r) &= \begin{cases} \frac{AM}{a+m}\sqrt{r} I_{(j+1/2)\mp 1/2}(Mr) & 0 < r < 1 - \epsilon, \\ \frac{L\sqrt{r}}{a-\frac{\lambda}{2\epsilon}+m} \left(B_1 J_{(j+\frac{1}{2})\mp\frac{1}{2}}(Lr) + B_2 Y_{(j+\frac{1}{2})\mp\frac{1}{2}}(Lr) \right) & 1 - \epsilon < r < 1 + \epsilon, \\ -\frac{CM}{a+m}\sqrt{r} K_{(j+1/2)\mp 1/2}(Mr) & r > 1 + \epsilon, \end{cases}
 \end{aligned} \tag{29}$$

where J and Y denote the first order Bessel's functions and I and K the second order Bessel's functions,

$$M = \sqrt{m^2 - a^2}, \quad L = \sqrt{\left(\frac{\lambda}{2\epsilon} - a\right)^2 - m^2}$$

and $(A, B_1, B_2, C) \neq 0$ are some constants. Note that $M \in \mathbb{R}$ by the assumptions on a , but L could be complex. Note also that $f, g \in H^1((0, +\infty) \setminus \{1 - \epsilon, 1 + \epsilon\}, r dr)$. To ensure that they belong to $D(h(\lambda, \epsilon)_{m_j, k})$ we have to verify that both f and g are continuous in $1 - \epsilon$ and $1 + \epsilon$, which means that the following linear system must hold:

$$\left\{ \begin{aligned}
 0 &= A\sqrt{1 - \epsilon} I_{(j+1/2)\pm 1/2}(M(1 - \epsilon)) - B_1\sqrt{1 - \epsilon} J_{(j+1/2)\pm 1/2}(L(1 - \epsilon)) \\
 &\quad - B_2\sqrt{1 - \epsilon} Y_{(j+1/2)\pm 1/2}(L(1 - \epsilon)), \\
 0 &= A \frac{\sqrt{1 - \epsilon} M I_{(j+1/2)\mp 1/2}(M(1 - \epsilon))}{a + m} - B_1 \frac{2\epsilon L\sqrt{1 - \epsilon} J_{(j+1/2)\mp 1/2}(L(1 - \epsilon))}{2a\epsilon - \lambda + 2m\epsilon} \\
 &\quad - B_2 \frac{2\epsilon L\sqrt{1 - \epsilon} Y_{(j+1/2)\mp 1/2}(L(1 - \epsilon))}{2a\epsilon - \lambda + 2m\epsilon}, \\
 0 &= B_1\sqrt{1 + \epsilon} J_{(j+1/2)\pm 1/2}(L(1 + \epsilon)) + B_2\sqrt{1 + \epsilon} Y_{(j+1/2)\pm 1/2}(L(1 + \epsilon)) \\
 &\quad - C\sqrt{1 + \epsilon} K_{(j+1/2)\pm 1/2}(M(1 + \epsilon)), \\
 0 &= B_1 \frac{2\epsilon L\sqrt{1 + \epsilon} J_{(j+1/2)\mp 1/2}(L(1 + \epsilon))}{2a\epsilon - \lambda + 2m\epsilon} + B_2 \frac{2\epsilon L\sqrt{1 + \epsilon} Y_{(j+1/2)\mp 1/2}(L(1 + \epsilon))}{2a\epsilon - \lambda + 2m\epsilon} \\
 &\quad + C \frac{\sqrt{1 + \epsilon} M K_{(j+1/2)\mp 1/2}(M(1 + \epsilon))}{a + m}.
 \end{aligned} \right. \tag{30}$$

Since this is a 4×4 homogeneous linear system on A, B_1, B_2 and C and we are assuming that

$(A, B_1, B_2, C) \neq 0$, we deduce that the associated matrix has null determinant. So, if we set

$$\begin{aligned}
D_{j\pm 1/2}^\epsilon(a, \lambda) := & \frac{2(a+m) K_{(j+1/2)\pm 1/2}(\sqrt{m^2 - a^2}(1+\epsilon))}{\epsilon(-2a\epsilon + \lambda - 2m\epsilon)^2} \\
& \times \left\{ -2L\epsilon(a+m) I_{(j+1/2)\pm 1/2}(\sqrt{m^2 - a^2}(1-\epsilon)) \right. \\
& \quad \times \left[J_{(j+1/2)\mp 1/2}(L(1+\epsilon)) Y_{(j+1/2)\mp 1/2}(L(1-\epsilon)) \right. \\
& \quad \quad \left. - J_{(j+1/2)\mp 1/2}(L(1-\epsilon)) Y_{(j+1/2)\mp 1/2}(L(1+\epsilon)) \right] \\
& \quad - \sqrt{m^2 - a^2}(2a\epsilon - \lambda + 2m\epsilon) I_{(j+1/2)\mp 1/2}(\sqrt{m^2 - a^2}(1-\epsilon)) \\
& \quad \times \left[J_{(j+1/2)\pm 1/2}(L(1-\epsilon)) Y_{(j+1/2)\mp 1/2}(L(1+\epsilon)) \right. \\
& \quad \quad \left. - J_{(j+1/2)\mp 1/2}(L(1+\epsilon)) Y_{(j+1/2)\pm 1/2}(L(1-\epsilon)) \right] \left. \right\} \quad (31) \\
& + \frac{\sqrt{m^2 - a^2} K_{(j+1/2)\mp 1/2}(\sqrt{m^2 - a^2}(1+\epsilon))}{\epsilon^2 L(2a\epsilon - \lambda + 2m\epsilon)} \\
& \times \left\{ -2L\epsilon(a+m) I_{(j+1/2)\pm 1/2}(\sqrt{m^2 - a^2}(1-\epsilon)) \right. \\
& \quad \times \left[J_{(j+1/2)\pm 1/2}(L(1+\epsilon)) Y_{(j+1/2)\mp 1/2}(L(1-\epsilon)) \right. \\
& \quad \quad \left. - J_{(j+1/2)\mp 1/2}(L(1-\epsilon)) Y_{(j+1/2)\pm 1/2}(L(1+\epsilon)) \right] \\
& \quad + \sqrt{m^2 - a^2}(2a\epsilon - \lambda + 2m\epsilon) I_{(j+1/2)\mp 1/2}(\sqrt{m^2 - a^2}(1-\epsilon)) \\
& \quad \times \left[J_{(j+1/2)\pm 1/2}(L(1+\epsilon)) Y_{(j+1/2)\pm 1/2}(L(1-\epsilon)) \right. \\
& \quad \quad \left. - J_{(j+1/2)\pm 1/2}(L(1-\epsilon)) Y_{(j+1/2)\pm 1/2}(L(1+\epsilon)) \right] \left. \right\},
\end{aligned}$$

then the determinant of the matrix associated to the linear system (30) is $\frac{\epsilon(\epsilon^2-1)}{(a+m)^2} D_{j\pm 1/2}^\epsilon(a, \lambda)$, which vanishes if and only if

$$D_{j\pm 1/2}^\epsilon(a, \lambda) = 0. \quad (32)$$

We can conclude that, given

$$\varphi = \begin{pmatrix} f \\ g \end{pmatrix} \quad \text{with } f \text{ and } g \text{ as in (29),}$$

φ is an eigenfunction for $h(\lambda, \epsilon)_{m_j, k_j}$ with eigenvalue a if and only if $D_{j\pm 1/2}^\epsilon(a, \lambda) = 0$. This means that the function

$$\psi(x) = \frac{1}{r} \left(f(r) \Phi_{m_j, k_j}^+(\hat{x}) + g(r) \Phi_{m_j, k_j}^-(\hat{x}) \right)$$

is an eigenfunction for $H_{\lambda, \epsilon}$ with eigenvalue a .

In order to compare (31) and (23), let us draw some pictures of these relations for some concrete values of the underlying parameters, say $m = 1$, $k = 1$ and $\epsilon = 2^{-10}$. Figures 1 and 2 describe the set of $(a, \lambda) \in (-1, 1) \times (0, 10)$ that verify (32) and (23), respectively.

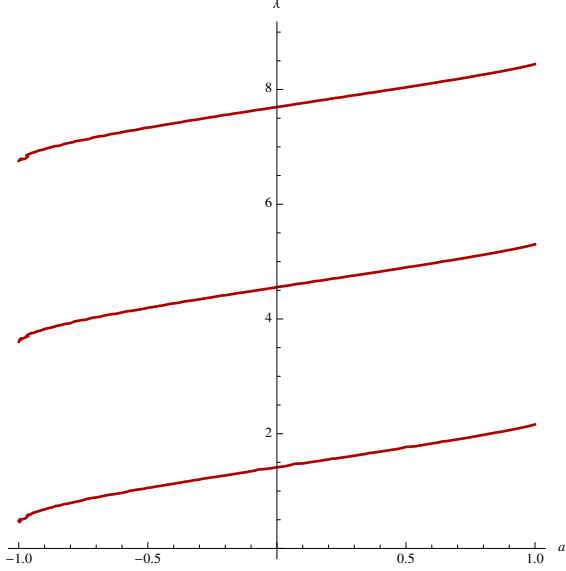


FIG. 1: The set of points (a, λ) satisfying (32).

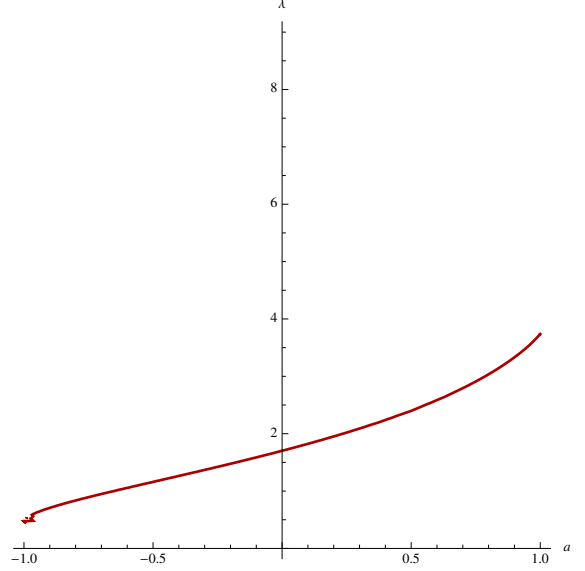


FIG. 2: The set of points (a, λ) satisfying (23).

Looking at Figures 1 and 2 we note that there is no apparent relation between (32) and (23).

However, the next result proves that there is indeed a precise connection between both equations when one takes the limit $\epsilon \rightarrow 0$ in $D_{j\pm 1/2}^\epsilon(a, \lambda)$.

Lemma IV.1. *Let $j = 1/2, 3/2, \dots$ and $D_{j\pm 1/2}^\epsilon$ and $D_{j\pm 1/2}$ be defined by (31) and (24), respectively. Then, for any $\lambda > 0$,*

$$\lim_{\epsilon \rightarrow 0} D_{j\pm 1/2}^\epsilon(a, \lambda) = \frac{4(a+m)}{\lambda\pi \left(1 + \tan\left(\frac{\lambda}{2}\right)^2\right)} D_{j\pm 1/2}\left(a, 2 \tan\left(\frac{\lambda}{2}\right)\right) \quad \text{uniformly on } a \in (-m, m).$$

Proof. Note that $L \rightarrow +\infty$ uniformly in $a \in (-m, m)$ when $\epsilon \rightarrow 0$, thus we can use the asymptotics

$$\begin{aligned} J_n(z) &= \sqrt{\frac{2}{\pi z}} \left(\cos\left(z - \frac{1}{2}n\pi - \frac{1}{4}\pi\right) + e^{|\Im(z)|} o(1) \right) \quad \text{for } |z| \rightarrow +\infty, \\ Y_n(z) &= \sqrt{\frac{2}{\pi z}} \left(\sin\left(z - \frac{1}{2}n\pi - \frac{1}{4}\pi\right) + e^{|\Im(z)|} o(1) \right) \quad \text{for } |z| \rightarrow +\infty, \end{aligned}$$

see [13, Equation 10.7.8]. Inserting these two relations in (31) and taking $\epsilon \rightarrow 0$, we get that, uniformly on $a \in (-m, m)$,

$$\begin{aligned} &\lim_{\epsilon \rightarrow 0} D_{j\pm 1/2}^\epsilon(a, \lambda) \\ &= \frac{4}{\lambda\pi} \left(M I_{(j+1/2)\mp 1/2}(M) \left((a+m) \cos(\lambda) K_{(j+1/2)\pm 1/2}(M) - M \sin(\lambda) K_{(j+1/2)\mp 1/2}(M) \right) \right. \\ &\quad \left. + (a+m) I_{(j+1/2)\pm 1/2}(M) \left((a+m) \sin(\lambda) K_{(j+1/2)\pm 1/2}(M) + M \cos(\lambda) K_{(j+1/2)\mp 1/2}(M) \right) \right). \end{aligned}$$

Setting $t = 2 \tan\left(\frac{\lambda}{2}\right)$, we know that $\sin(\lambda) = \frac{t}{1+\frac{t^2}{4}}$ and $\cos(\lambda) = \frac{1-\frac{t^2}{4}}{1+\frac{t^2}{4}}$. Using (22), hence

$$\lim_{\epsilon \rightarrow 0} D_{j \pm 1/2}^\epsilon(a, \lambda) = \frac{16(a+m)}{\lambda\pi(4+t^2)} \left(\frac{t^2}{4} - \left((m+a) I_{(j+1/2) \pm 1/2}(M) K_{(j+1/2) \pm 1/2}(M) \right. \right. \\ \left. \left. - (m-a) I_{(j+1/2) \mp 1/2}(M) K_{(j+1/2) \mp 1/2}(M) \right) t - 1 \right),$$

which coincides with (24) if one replaces λ by $t = 2 \tan\left(\frac{\lambda}{2}\right)$ in there. \square

The following result resumes what we have proven so far with the aid of Lemma IV.1.

Theorem IV.2. *Let $\mu \in \mathbb{R} \setminus \{0\}$ and*

$$\lambda = 2 \tan\left(\frac{\mu}{2}\right).$$

Let $h(\lambda)_{m_j, k_j}$ be as in (6) and, for $\epsilon > 0$, let $h(\mu, \epsilon)_{m_j, k_j}$ be as in (27). If $a_\epsilon \in \sigma_p(h(\mu, \epsilon)_{m_j, k_j})$ and $\lim_{\epsilon \rightarrow 0} a_\epsilon = a$ for some $a \in (-m, m)$, then $a \in \sigma_p(h(\lambda)_{m_j, k_j})$.

V. ACKNOWLEDGEMENTS

We would like to thank Aingeru Fernández-Bertolín for addressing us the reference [4] and for the enlightening discussions. Mas was supported by the *Juan de la Cierva* program JCI2012-14073 and the funding projects MTM2011-27739 and MTM2014-52402 from the Government of Spain. Pizzichillo was supported by the MINECO project MTM2014-53145-P, by the Basque Government through the BERC 2014-2017 program and by the Spanish Ministry of Economy and Competitiveness MINECO: BCAM Severo Ochoa accreditation SEV-2013-0323. Both authors were also supported by the ERCEA, Advanced Grant project 669689 HADE.

-
- [1] S. Albeverio, F. Gesztesy, R. Hoegh-Krohn, and H. Holden. *Solvable models in quantum mechanics*. Springer Science & Business Media, 2012.
 - [2] N. Arrizabalaga, A. Mas, and L. Vega. Shell interactions for Dirac operators. *Journal de Mathématiques Pures et Appliquées*, 102(4):617–639, 2014.
 - [3] N. Arrizabalaga, A. Mas, and L. Vega. Shell interactions for Dirac operators: on the point spectrum and the confinement. *SIAM Journal on Mathematical Analysis*, 47(2):1044–1069, 2015.
 - [4] Á. Baricz and S. Ponnusamy. On Turán type inequalities for modified Bessel functions. *Proceedings of the American Mathematical Society*, 141(2):523–532, 2013.
 - [5] J. Behrndt, P. Exner, M. Holzmam, and V. Lotoreichik. On the spectral properties of Dirac operators with electrostatic δ -shell interactions. *arXiv preprint arXiv:1609.00608*, 2016.
 - [6] J. Behrndt, P. Exner, M. Holzmam, and V. Lotoreichik. Approximation of Schrödinger operators with δ -interactions supported on hypersurfaces. *Mathematische Nachrichten*, 290(8-9):1215–1248, 2017.
 - [7] J. Behrndt, and M. Holzmam. On Dirac operators with electrostatic δ -shell interactions of critical strength. *arXiv preprint arXiv:1612.02290*, 2016.
 - [8] J. Dittrich, P. Exner, and P. Šeba. Dirac operators with a spherically symmetric δ -shell interaction. *Journal of Mathematical Physics*, 30(12):2875–2882, 1989.

- [9] P. Exner. Leaky quantum graphs: a review. In *Analysis on graphs and its applications*, volume 77 of *Proc. Sympos. Pure Math.*, pages 523–564. Amer. Math. Soc., Providence, RI, 2008.
- [10] F. Gesztesy and P. Šeba. New analytically solvable models of relativistic point interactions. *letters in mathematical physics*, 13(4):345–358, 1987.
- [11] A. Mantile, A. Posilicano, and M. Sini. Self-adjoint elliptic operators with boundary conditions on not closed hypersurfaces. *Journal of Differential Equations*, 261(1):1–55, 2016.
- [12] A. Mas and F. Pizzichillo. Klein’s paradox and the relativistic δ -shell interaction in \mathbb{R}^3 . *submitted*, 2017.
- [13] F. W. Olver. *NIST handbook of mathematical functions*. Cambridge University Press, 2010.
- [14] T. Ourmières-Bonafos and L. Vega. A strategy for self-adjointness of dirac operators: Applications to the MIT bag model and delta-shell interactions. *arXiv:1612.07058*, 2016.
- [15] A. Posilicano. Self-adjoint extensions of restrictions. *Operators and Matrices*, 2(4):483–506, 2008.
- [16] P. Šeba. Klein’s paradox and the relativistic point interaction. *Letters in Mathematical Physics*, 18(1):77–86, 1989.
- [17] B. Thaller. *The Dirac equation*, volume 31. Springer-Verlag Berlin, 1992.