Sliding mode control applied to improve the dynamical response of a spherical 3D wind sensor for Mars atmosphere

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Abstract

The objective of this paper is to show how the structure of a thermal wind anemometer can be tuned to ensure a fast response when the sensor works in closed loop configuration (constant temperature operation). If the thermal filter associated to the sensor structure has only one significative time constant, the resulting system time response, working in closed loop, is enhanced beyond the natural limit imposed by its own thermal circuit. This effect is theoretically explained using the theory of sliding mode controllers. Experimental corroboration is presented by comparing the results obtained with two prototypes of a spherical wind sensor for Mars atmosphere. It will be shown that in case of having only one significant time constant, the time response in closed loop is much faster than the value associated with that time constant. It will be experimentally shown that this effect is lost when the thermal filter has more than one significant time constant. Diffusive representation is used to obtain arbitrary-order models of the thermal structures of the sensors.

Key words: Wind sensors, sliding mode controllers, diffusive representation, time-varying systems

1. Introduction

The objective of this paper is to present the improvements of the dynamical response of a spherical 3D wind sensor for Mars atmosphere. Wind sensing in Mars is a challenging task due to the rarefied atmosphere, mostly composed of CO$_2$, low pressure (in the range of 6-12mBar) and large temperature dynamical range: from 150 to 300K [1]. Several methods exist for measuring wind speed and direction on Mars surface. Thermal anemometry, though, is for now the most adequate method in this environment. Sensors based on ultrasound do not work well below 15mBar and require higher powers to operate [2]. Windsocks [3] and tell-tales [4] have been also used for detecting wind velocity and angle, but these methods require to process images to infer wind patterns. Thermal anemometry, then benefits from a general simplicity and robustness. This method detects the wind velocity by measuring the power losses of a heated element due to forced convection. These thermal sensors were first introduced for Mars missions in Viking Lander [5,6], in 1976. The wind sensor was formed by two small cylinders covered by a thin film of platinum that formed a resistance, which worked under constant temperature mode. The Pathfinder mission of 1997 [7,8] also included a wind detection sensor based on thermal anemometry. It consisted in a cylindrical mast with six platinum-iridium wires as heating elements. More recently, in August 2012, the Mars Science Laboratory (MLS) on board the Curiosity rover landed in the red planet. It included the REMS (Remote Environmental Monitoring Station) sensor suite that measures humidity, pressure, temperature, radiation and wind velocity [8,9,10]. For the wind velocity and direction detection, platinum resistors, fabricated on silicon technology, were employed. These resistors were grouped in sets of four in a coplanar plane, to achieve 2D wind sensitivity. The 3D direction was obtained placing three of these sets forming 120° with each other in a cylindrical boom. Two of these booms were incorporated in REMS. The same device concept is scheduled to fly again in the InSight Mission in 2018, the TWINS instrument (Temperature and Wind sensors for InSight mission). While Viking was able to send data for two years, Pathfinder sent wind...
direction data only. REMS is still operational, although some of the chips were damaged when landing [10].

Two sensor prototypes of an spherical anemometer, heritage of the REMS wind sensor, will be described in this paper. The main advantages of these sensors are their spherical geometry, which eases the 3D wind speed recovery reverse algorithm [11], and the simplicity of the system itself, with only a few heating elements. Furthermore, as the heating elements are inside a metallic shell, the sensor is more robust [11]. Figure 1 shows a picture of both sensors.

As it has been mentioned above, the dynamics of the two sensor prototypes is analyzed in this paper. In the more recent prototype, (Figure 1 B) the thermal coupling between the heating elements and the sectors conforming the shell of the sensor has been improved using bare platinum chip resistors instead of encapsulated commercial resistances (as in Figure 1 A).

The purpose of this paper is to experimentally assess how the dynamical response of the sensor, following a wind speed change, is significantly improved due to the better thermal coupling between the hot dice and the shell exposed to the wind. Moreover, in this paper we seek to find a general quantitative condition that thermal anemometers have to fulfill in order to have a closed loop response much faster than the open loop thermal time constants. This is the main reason why we compare sensors A and B having different thermal coupling to the ambient.

To do so we have used the tools of Sliding Mode Control (SMC) to the sigma-delta modulator we have used to close the loop of the two sensors. It is well known in SMC theory that thermal anemometers have to fulfill in order to have a closed loop response much faster than the open loop thermal time constants. This is the main reason why we compare sensors A and B having different thermal coupling to the ambient.

The SMC parameters are found from open loop measure-ments at several wind velocities using modelling based on Diffusive Representation. These state-space models have proven to be very well suited for thermal applications [12]. The extraction method is different from the one we have used in [13] in which each thermal dynamical model was obtained using an experiment at constant wind speed.

Both tools, SMC and DR, are used together and are explained in this paper. Section 2 describes the sensors and how they are operated. Section 3 presents the thermal models extracted from both sensors prototypes and the sliding mode analysis of how the system will respond under closed loop operation. Section 4 analyzes the conditions for obtaining a matched or a mismatched response in the sensor. Finally, experimental results are shown in section 5.

2. Sensor Description

Figure 2 shows a schematic view of the two sensors. In the following, each sensor is going to be described briefly.

SENSOR A. The initially designed sensor is composed of two silver hemispheres connected to a PCB (Printed Circuit Board), which provides mechanical and electrical support. In order to reduce emissivity and avoid oxidation, the hemispheres are polished and sputtered with a thin layer of gold on their surface respectively. Each hemisphere integrates a SMD (Surface Mounted Device) Pt resistor. The resistors, $R_A$ and $R_B$, allow to heat the hemispheres and sense the temperature. In the PCB, two more Pt resistors are attached, $R_{core1}$ and $R_{core2}$, in order to heat the core of the sphere to reduce the conduction heat flux between the hemispheres and the supporting structure. The resistors of sensor A are commercial surface mounted 100 Ω platinum resistors with temperature coefficient $\alpha = 0.00385^\circ C^{-1}$ with different encapsulations; SMD0603 for the spherical sectors and SMD0805 for the PCB.

SENSOR B. The new design is composed of four silver sectors connected to a PCB. Now, the silver sectors (which have the same polishing and sputtering process as that of sensor A), form a tetrahedral sphere of same radius as the sphere formed in sensor A. Again, each sector and the PCB have a Pt resistor integrated on them ($R_A$, $R_B$, $R_C$, and $R_D$ in the sectors, and $R_{core1}$ and $R_{core2}$ in the PCB), which allow heating and sensing of the temperature and reducing conduction heat flux. A key design factor for improving the time response of the sensor is to ensure a good thermal connection between heaters and sectors. To achieve this goal non-encapsulated chip resistors have been used this time. Additionally, the electrical connexion between the PCB and the resistors is made by a wire bonding process. For this purpose, resistors have been fabricated in our clean room using a Silicon wafer. Silicon oxide ($SiO_2$) has been grown thermally for electrical isolation. On top of the $SiO_2$, the platinum resistors were patterned with a photolithography process. The Pt resistors of sensor B have a nominal value of 175Ω and a temperature coefficient $\alpha = 0.0031^\circ C^{-1}$.

2.1. Sensor Operation

Thermal anemometers can work in open or closed loop operation mode. In open loop mode, a constant current or voltage is injected into the heaters of the sensor while the temperature...
changes, due to convection losses to the fluid, are measured to infer the wind speed. In the closed loop, however, the temperature in the heated elements is forced to a constant value (CTA mode, Constant Temperature Anemometry), and the power required at every sector to keep constant the temperature is the output signal of the sensor. The wind velocity is calculated from the thermal conductance (ratio of the power and the temperature difference between the hot point and the ambient) which is related to the Nusselt number (embodiment of the Reynolds number and hence the wind speed). It is known that constant temperature operation mode is better in terms of time response [8]. In section 5 of this paper, both operation modes are tested and compared.

\[
A = \begin{bmatrix}
    \sum_{n.g(n)=w_1} \psi_1^{(n)}(t_0) & \ldots & \sum_{n.g(n)=w_1} \psi_K^{(n)}(t_0) \\
    \vdots & \ddots & \vdots \\
    \sum_{n.g(n)=w_1} \psi_1^{(n)}(t_f) & \ldots & \sum_{n.g(n)=w_1} \psi_K^{(n)}(t_f)
\end{bmatrix}, \quad [1]_K
\]

3. Dynamical Analysis

An analysis of the dynamics of the sensor structures has been carried out to understand and predict the time response of each sensor. In this section we present the tools of diffusive representation and sliding mode controllers, that will be used together. The same modelling and analysis method has been employed both in sensors A and B.

3.1. Thermal modelling: Diffusive Representation

The theory of Diffusive Representation allows obtaining exact and approximate state realizations of a wide class of integral operators. This mathematical tool is suitable for system identification of any physical phenomena based on diffusion [14].

Given a non-rational transfer function, \( H(p) \), associated with a convolution causal operator denoted by \( H(\hat{\xi}(\cdot)) \), the diffusive realization of this operator is expressed by the following input \( u \) – output \( y \) state space realization of \( u \mapsto y = H(\hat{\xi})u = h \ast u \) of the form \([15]\):

\[
\frac{\partial \hat{\xi}(\xi, t)}{\partial t} = -\xi \hat{\xi}(\xi, t) + u(t), \quad \psi(\xi, 0) = 0
\]

\[
y(t) = \int_{-\infty}^{0} \eta(\xi, t)\psi(\xi, t)d\xi
\]

where \( \xi \in \mathbb{R} \) is frequency, \( \eta(\xi, t) \) is the diffusive symbol of \( H(\hat{\xi}(\cdot)) \) that represents how the system behaves, and the state-variable \( \psi(\xi, t) \) is a time-frequency representation of the input, called the diffusive representation of \( u(t) \) \([15]\).

To be able to handle experimental data, a discrete approximation of \( H(\hat{\xi}(\cdot)) \) can be built discretizing the continuous variable \( \xi \) into \( [\xi_1, \ldots, \xi_K] \), where \( K \) is the order of the discretized model. This leads to an input – output approximation \( u \mapsto \hat{y} \approx y = H(\hat{\xi}/\Delta t)u \).

The dynamical system can be written as:

\[
\begin{align*}
\psi_1^{(n)}(t) & = 0 & t \in [t_0, t_n] \\
\dot{\psi}_k^{(n)}(t) & = -\xi_k \psi_k^{(n)}(t) + u(t) & t \in [t_n, t_{n+1}] \\
\dot{\psi}_k^{(n)}(t) & = -\xi_k \psi_k^{(n)}(t) & t \geq t_{n+1}
\end{align*}
\]

\[
\hat{y}(t) = \sum_{n,k}^{K} \eta_k^{(n)}(\psi_k^{(n)}(t)) + \sum_k K c_k e^{-\xi_k(t-t_0)}
\]

where \( \eta_k^{(n)} \in \mathbb{R}^K \) is the diffusive symbol associated to the conditions of the system in the \( n \)-th interval in \( t \in [t_n, t_{n+1}] \). \( c_k \in \mathbb{R}^K \) represents the initial conditions of the system at the beginning of the measurements, at \( t = t_0 \). These intervals are a discretization of time for which the dissipative symbols can be considered constant in a wind experiment. This model extraction method, is different from the method described in \([13]\), where the dissipative symbols for each wind velocity were inferred from long open-loop measurements at a single wind velocity. In that case, every time the wind was switched, the memory of the system was lost, and therefore, it was necessary to find new initial conditions. It was not possible to predict the thermal dynamics without having before hand the closed-loop measurements under switching wind.

On the contrary, in this paper, the thermal characterization of the wind sensors is made from open-loop measurements in which wind speed is being continuously switched between several wind velocities. The diffusive symbols corresponding to each wind velocity are hence inferred from a single experiment, with \( t \in [t_0, t_f] \). The number of wind velocities for which the sensor will be characterized is \( J \). The wind speeds, \( [w_1, \ldots, w_J] \), are a uniformly distributed random sequence. Each wind speed is applied in an interval \( t \in [t_n, t_{n+1}] \), where \( N \) is the number of wind events, such that, \( N \gg J \). All the wind velocities are applied for a short time \( \Delta t = (t_{n+1} - t_n) \), where \( \Delta t \ll t_f \), for all \( n \). In the meantime, a PRBS heating current is applied to the resistors in open-loop, with a period \( T_{PRBS} \ll \Delta t \ll t_f \).

The solution to the identification problem is found solving the finite least squares problem:

\[
\min_{\hat{\xi} \in \mathbb{R}^K} ||A\hat{\xi} - Y||^2
\]

where \( A \) is of the form of Eq. [1] and the vector \( \hat{\xi} \) is that of Eq. [5].

In matrix of Eq. [1], the function \( g(n) \) returns the wind \( w_j \) applied at every wind interval \( n \). The last column of this matrix, \([1]_K \in \mathbb{R}^K \) is an all ones vector and

\[
e^{-\xi(t-t_0)}[1] = [e^{-\xi(t-t_0)}, \ldots, e^{-\xi(t-t_0)}] \]

The diffusive symbol vector that is going to be inferred is of the form of Eq. [5]:

\[
\hat{\xi}^{T} = ([\hat{\xi}_1^{(n)}, \ldots, \hat{\xi}_K^{(n)}]^T, \ldots, [\hat{\xi}_1^{(n)}, \ldots, \hat{\xi}_K^{(n)}]^T, [c_1, \ldots, c_K]^T] \]

where \([\hat{\xi}_k^{(n)}]^T \) is the diffusive symbol corresponding to the wind \( w = w_j \). The measurements vector is \( Y^T = [T(t_0), \ldots, T(t_f)] \).
being \(T(t_i)\) the temperature in the instant \(t_i \in [t_0, t_F]\) in our case. The solution to the inference problem is classically given by:

\[
\dot{\hat{y}} = [A^* A]^{-1} A^* Y
\]  

(6)

3.2. Constant Temperature Anemometer dynamics

The analysis of the heat flow dynamics in an anemometer is contemplated under constant temperature mode operation. This study will be undertaken considering that the controller is a sigma-delta modulator, which is a well-known type of analog-to-digital converter [16]. A thermal sigma-delta modulator, applies a modulated power to maintain the temperature constant in the sensor shell. In this control, at each sampling period, \(T_s\), the temperature of the hot element, \(T_n = T(nT_s)\), is compared to the desired target temperature, \(\Delta T\). If \(T_n - \Delta T\), \(P_{off}\) is delivered during the following sampling period, and if \(T_n < \Delta T\), \(P_{on}\) is injected instead. Sigma-delta modulators can be seen as an example of a sliding mode controller (SMC). [17]. These type of controllers change the dynamics of the system by applying a discontinuous control signal so that, under some conditions, the system “slides” on a certain control surface [13]. In the case of sigma-delta modulation, the typical hysteresis used in SMC is substituted by a discretization of time.

To analyze the dynamics of the system for the general case of finite order thermal systems a control surface has to be defined. This is:

\[
\sigma(\psi(t)) = \Delta T - \dot{\bar{y}}(t) = \Delta T - \sum K k(t) \eta_k(t) \psi_k(t) \]  

(7)

The objective is to place the system within the control surface \(\sigma(\psi(t)) = 0\) in finite time, in our case to have a constant temperature. This is possible if, [13]:

\[
P_{off} \sum K \frac{\eta_k(t)}{\xi_k} < \Delta T < P_{on} \sum K \frac{\eta_k(t)}{\xi_k} \]  

(8)

The condition \(\sigma(\psi(t)) = 0\) gives the necessary equivalent control under a sliding motion to maintain the temperature of the shell constant and for a system described by diffusive representation is given by Eq. (9):

\[
u_{eq}(t) = \Gamma^{-1} \sum K k(t) \eta_k(t) \xi_k \psi_k(t) \]  

(9)

where \(\Gamma = \sum K \eta_k(t)\).

Then, the dynamics of the system in the control surface is determined when the equivalent control is applied:

\[
\dot{\psi}_k(t) + \xi_k \psi_k(t) = \Gamma^{-1} \sum K \eta_k(t) \xi_k \psi_k(t) \]  

(10)

If the equilibrium point (\(\dot{\psi}_k(t) = 0\)) is reached, the system takes the steady-state values \(\psi_k(\xi_k) = \xi_k / \xi_k\) and \(u_0 = \Delta T / \sum K \eta_k / \xi_k\).

The time evolution predicted by the SMC analysis is obtained taking into account the changes in wind speed. The equivalent control to maintain the desired target temperature can be predicted using the diffusive symbols inferred in the open-loop characterization. The wind velocities for which the analysis will be done must be the same as the ones for which the wind sensor has been characterized previously. Now, in this experiment, the wind is kept constant within the time intervals \(\Delta t_k = (t'_{k+1} - t'_k)\), for \(n = 0, \ldots, N'\), where \(N'\) is the number of wind intervals such as \(N' \geq J\).

The first step to calculate the equivalent control is to obtain the point \(t'_k \in [t'_n, t'_{n+1}]\) at which the system reaches the control surface, \(\sigma = 0\). From a zero initial condition, a constant \(u(t) = P_{on}\) is applied till the control surface is reached. The intersection point of this initial trajectory with the surface is the initial condition for the sliding motion on the control surface following expressions (9) and (10) for the first wind step of the experiment. Now, Eq. (10) can be expressed as:

\[
\dot{\psi}_k(t') = -D(t') \psi(t'), \quad t' \geq t'_k \]  

(11)

where \(D(t')\) collects the terms in Eq. (10). Since wind velocities are considered constant during each interval \(\Delta t_k\) of the experiment, \(D(t')\) is piecewise constant and therefore for \(t' \in [t'_n, t'_{n+1}]\):

\[
\psi'(t') = e^{D(t'-t'_0)} \psi(t'_0) \]  

(12)

with \(\eta_k(t)\) the diffusive symbol of the sensor working under the wind of the \(n\)-th time interval.

Replacing Eq. (12) in the equivalent control expression of Eq. (9) we reach:

\[
u_{eq}(t') = H(t') \psi(t') \]  

(13)

where:

\[
H(t') = \Gamma^{-1} \sum K k(t') \xi_k e^{-D(t'-t'_0)} e^{-D(t'-t'_1)} \cdots e^{-D(t'-t'_N)} \xi_k \]  

(14)

where it has been assumed that \(t'_k > t'_k\).

The extracted diffusive symbols from [3.1] allow to predict the system behaviour under closed-loop operation control in the presence of wind variations. The equivalent control to maintain constant the temperature is directly obtained from the application of Eq. (13).

4. Reduction of the number of state variables for ensuring matched disturbances for time response acceleration

Wind velocity changes can be seen as an external disturbance of the system, represented by an additional term: \(\theta_k(\psi, t) = (\theta_1(\psi, t), \ldots, \theta_N(\psi, t))^T\). Since, \(\psi(0) = 0\), the thermal system can also be expressed as [13]:

\[
\dot{\psi}_k(t) = -\xi_k \psi_k(t) + \frac{\eta_k}{2} (P_{off} + P_{on})
+ \theta_k(\psi, t) + \theta_k(\psi, t)
\]  

(15)

\[
\sigma(\psi(t)) = \Delta T - \sum K \eta_k \psi_k(t) \]
where $u = \text{sgn}(\sigma^r(\psi(t)))$ and the control surface is $\sigma^r(\psi(t)) = \Delta T = \sum \chi_k \psi_k(t)$.

Assuming the sliding motion is taking place, this is $d\sigma^r/dt = 0$, the equivalent control is now:

$$u_{eq} = \frac{\sum_k \xi_k \psi_k(t) - \frac{1}{2}(P_{on} + P_{off}) - \chi(\psi, t)}{\chi(\psi, t)} \in [-1, 1] \quad (16)$$

with $\chi(\psi, t) = \sum \theta_k(\psi, t)$. The new dynamics in the control surface is therefore determined by:

$$\dot{\psi}_k(t) = -\xi_k \psi_k(t) + \frac{\theta_k(\psi, t)}{2}(P_{on} + P_{off}) + \frac{\sigma_k(\psi, t)}{2} \left[ \sum \xi_k \psi_k(t) - \frac{1}{2}(P_{on} + P_{off}) - \chi(\Psi, t) \right] \quad (17)$$

As it is well known, disturbances can be decomposed in a matched and a mismatched component $[17]$. In the case of a matched disturbance, $\theta_k(t) = -\frac{\sigma_k(\psi, t)}{2}$, the disturbance is proportional to the diffusive symbol, i.e., $\theta_k(\psi, t) = \sigma_k \psi$ for some $\sigma \in \mathbb{R}$. Under this condition, the control method will be able to handle successfully the disturbances with a fast (almost instantaneous) response, since the dynamics of the system on the control surface would remain undisturbed. However, in the general case where a mismatched component is present, any disturbance will generate a change in the dynamics of the system on the control surface, therefore generating a slow dynamics in the equivalent control. Usually, the dynamics triggered by a mismatched component is expected to be of the same order as the one in the open loop response, i.e., slow.

One way of ensuring that any disturbance is a matched disturbance is the very simple case of having one state variable. This means that in a general case, it is very convenient to have only
one significant state variable of the thermal system, or what is the same, only one node of temperature in the interface between the heaters and the sensor’s sectors. This is achieved in our prototypes if the thermal connection between the heated devices and the sectors is good enough. Experimental corroboration of this fact is shown in next section, with the thermal models of the sensors obtained using Diffusive Representation.

5. Experimental Results

Open and closed-loop experiments have been carried out to compare the time response of both sensors. For the experiments, the sensors were enclosed separately in an hypobaric stainless steel chamber. The pressure of the chamber was set to 45 mBar and the measurements were made at room temperature. An automatic fan, located in front of the sensor, provides the necessary wind changes.

5.1. Thermal model extraction

First, the DR models for both wind sensors have been obtained. For this task, in open-loop mode operation, the Pt resistor of one hemisphere of sensor A is excited with a 10Hz PRBS current sequence, between 10 and 20mA, while the rest of the resistor components are controlled to maintain them at a constant temperature. For sensor B this operation is the same with the exception that the sensor is excited with a current between 8.9 and 19.9mA. Simultaneously, the wind velocity of the fan is changed following a uniformly distributed random sequence. The wind velocities, \(w_1, \ldots, w_J\), the fan can take are: 0.5, 0.7, 0.9, 1.1, 1.3 and 1.5 m/s.

The total experiment duration for each sensor has been 6 hours. During the whole experiment, the wind has changed automatically 180 times with a period \(T_w = 2\) min. The frequency mesh of the diffusive symbol is set according to the spectral contents of the measurements, therefore, it is chosen in accordance with the duration and sampling period of the experiment and is geometrically spaced. The minimum frequency is set to \(f_{\text{min}} = \frac{1}{T_w} \approx 2.7 \times 10^{-2}\)Hz, where \(T = 21600\)s is the total duration of the experiment and \(J\) is the number of wind flows applied during the experiment, in this case \(J = 6\). The maximum frequency is chosen accordingly with the sampling period. In this case, the temperature of the Pt resistors is sampled with a rate \(T_s = 5\)ms, therefore, \(f_{\text{max}} = \frac{1}{T_s} = 10\)Hz, which is the frequency at which the spectrum of the PRBS signal used in the experiment drops to zero. [13]

Applying the least-squares problem of Eq. (4), the DR models for each sensor have been inferred. In Figures 5 and 6 the matching between the fitted and the experimental data of sensor A and B, together with the applied wind sequence, can be observed. The model order chosen for the fitting is \(K = 9\) in sensor A and \(K = 8\) in sensor B. The diffusive symbols inferred for each wind velocity at each sensor are shown in Figure 4. As it can be observed, the diffusive symbol of sensor A has two significant maxima at approximately \(f \approx 0.015\)Hz and \(f \approx 0.2\) Hz. In sensor B, however, the maximum appears only at \(f \approx 0.025\) Hz. The fact that the diffusive symbol of sensor B has only one significant pole, imply that the disturbances generated by the wind are matched, since the condition \(\theta_k(\theta_k, t) = \alpha_\theta k\) is accomplished because there is only one significant variable. This means that the thermal system has a single node of temperature in the interface between the heating devices and the tetrahedral sectors. As it will be seen below in another experiment, the reduction to only one significant time constant will lead to a faster time response of this sensor under closed loop, where the control will almost instantaneously absorb the changes due to wind variations.

Additionally, both sensors have been tested in open loop mode. In Figures 5 and 6 the temperature changes corresponding to wind velocity changes when a constant current is being injected in the hemispherical and the tetrahedral sensor respectively can be observed. A 15mA current has been injected into a sector of both sensors while the rest of the resistances were controlled to a constant temperature. The changes in the wind velocities went from 0.5 to 1.5m/s in ascendant steps of 0.2m/s.
every 300s and equally from 1.5 to 0.5m/s in descending order. As it can be seen in the inner figures, the open loop time response is slow, approximately 150s for sensor A and 100s for sensor B. In open-loop configuration the time response difference between both sensors is not substantial, and the one that exists is probably because the size of the sectors of sensor B; as they are smaller than the sectors in A they will have less mass and will heat up faster.

As it can be seen in the inner figures, the open loop time response A in blue and the the result of applying the sliding mode analysis using the approximately 90s in sensor A and of 5s in case of sensor B (note the thermal models based on Di 435

5.2. Analysis of closed loop operation

Under closed-loop operation mode, the system is maintained at a constant target temperature $\Delta T = 7.2\,K$ for sensor A and $\Delta T = 11.5\,K$ for sensor B. The wind velocities, the names for which the system has been already characterized, changed every 300s with the same wind steps as in the constant current open-loop experiment. In Figure 7, the experimental control waveform for sensor A (top) and B (bottom) is shown together with the sliding mode analysis prediction. As it can be observed there is a good match between the predictions of the SMC theory using the thermal models based on Diffusive Representation.

On the other hand, the time response to a wind change is of approximately 90s in sensor A and of 5s in case of sensor B (note that the fan has some few seconds of delay until the wind speed is achieved). Note that the time response in closed loop (5s) is much smaller than the most significant time constant of sensor B (100s). This experimental result indeed confirms the predictions made in the previous sections. Sensor A, presents more than one time constant, and therefore the main gain of using a closed loop operation is not obtained since this implies the presence of mismatched components in the disturbances generated by the wind.

6. Conclusions

The dynamics of a new spherical wind sensor for Mars atmosphere under constant temperature control has been enhanced improving the thermal connection of the heating elements with the device sectors. For this sensor structure, this change generates a thermal filter with only one significant time constant. Any wind change can be understood as a disturbance of the system dynamics. Since closed loop operation can be seen as an example of a Sliding Mode Controller, by ensuring the presence of only one time constant, any wind disturbance becomes a matched disturbance. This results in a closed loop time response much faster than the time constant of the thermal sensor itself. Experimental corroboration of this analysis has been presented using two sensor prototypes with different thermal filters, one of which has only one significant time constant.

Biographies

Maria-Teresa Atienza received the B.Sc. degree in Electronic Engineering and the M. Sc. degree in Physics Engineering from the Basque Country University, Spain, in 2012 and 2013 respectively. She joined Micro and Nano Technologies Research Group of the Universitat Politècnica de Catalunya (UPC), Spain, as a Ph.D. degree student where she is working on the identification of thermal systems applied to wind sensors for space applications.

Łukasz Kowalski obtained MSc degree in Electronics and Telecommunications from the Technical University of Lodz, Poland in 2005. He obtained PhD from Universitat Politècnica de Catalunya, Spain in 2016 for thesis ‘Contribution to advanced hot wire wind sensing’. He has joined MNT group in September 2005. He has participated in space exploration projects: REMS(MSL), MEIGA (MetNet), MEDA (Mars2020). His work is aimed to the development of thermal anemometers for the atmosphere of Mars. His research areas include dimensionless analysis, FEM/CFD simulations, thermal modeling, inverse algorithm optimization for wind speed and incidence angle retrieval. He has co-authored two patents and several scientific papers.
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