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Research paper

Variability of wave impact measurements on vertical breakwaters

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ABSTRACT

The measurement of wave-induced impacts on structures involves significant scientific and engineering challenges. Regular nearly breaking pgb-gand broken waves generated following cnoidal and first-order Stokes wavemaker theory have been considered to study the variability of wave impacts on vertical breakwaters. Four small-scale hydraulic experiments were carried out and repeated 120 times using high-speed pressure and force measurement equipment. High variability in the measured pressure field and total force were observed, reflecting the random nature of the studied phenomena. The force variability was similar for the nearly breaking and broken waves. The maximum measured force was between 166 and 177% of the minimum measured force. In relation to the force distribution and maximum pressure points, the observed variability is higher for broken waves. To deal with such variability in the observations, suitable probability distributions for forces (*GEV*) and pressures (*Gamma*) are proposed.

Keywords: Impact loads; pressure/force measurements; test repeatability; vertical structures; wave impact; wave loads

1 Introduction

The continuous draft increase in big vessels implies the construction of harbour protection structures at deeper waters. At these depths, the amount of material needed for rubble mound breakwaters makes them less cost-effective than vertical structures. However, during severe storms, wave breaking often happens in front of vertical breakwaters. This is mainly due to the high reflection associated to vertical structures, which can induce very steep and plunging wave conditions in front of the structure. Oumeraci (1994) provided a comprehensive review of the most important historical failures of this type of structure, highlighting the impact induced by breaking clapotis as the main reason for failure. Large-scale hydraulic tests have been performed over the last 20 years, leading to

different formulations that allow calculation of loads induced by wave breaking on vertical breakwaters (Allsop, Kortenhaus, Oumeraci, & McConnell, 1999; Bullock, Obhrai, Peregrine, & Bredmose, 2007; Cuomo, Allsop, Bruce, & Pearson, 2010; Oumeraci et al., 2001). However, for impulsive loads on vertical breakwaters Takahashi's formula (Takahashi, 2002; Takahashi, Tanimoto, & Shimosako, 1994), based on small-scale hydraulic experiments, is one of the most frequently used. Takahashi's formula introduces an impulsive pressure coefficient into the Goda formula (Goda, 1974, 1985), initially developed for quasi-static loads. Most equations for impulsive loads on vertical breakwaters are based on semi-empirical models calibrated with wave flume (small and large-scale) data and, in some cases, with prototype-measured data (Minikin, 1963). For a comprehensive hydraulic design of vertical breakwaters, the following wave

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induced load-related variables should be considered: (i) maximum pressure; (ii) location of pressure maxima; and (iii) total force and its application point (overturning moment).

Vertical pressure distributions in hydraulic tests can be determined using the results from arrays of pressure transducers. The maximum pressure and its location can then be extracted from the pressure distribution diagram. During the design process, particular attention should be paid to where the structure joints are located (Cox & Cooker, 2001), as they should not be placed close to the maximum pressure point. For the hydraulic design of a vertical breakwater, sliding and overturning are the primary failure modes (Goda, 1974) and are directly associated with force/pressure distribution. During hydraulic tests, the total force can be calculated from data obtained by load cells or by computing the integral of the vertical pressure distribution. The resulting force application point can be calculated by finding the mass centre of the vertical pressure distribution, experimentally determined if a combination of two (or more) load cells have been used for measuring the total force. Historically, in most of the wave impact physical tests, a vertical array of pressure transducers was installed on the seaward side of the structure (e.g. Bullock et al., 2007; Cuomo et al., 2010). Only a few experiments were performed using load cells for the direct measurement of the total force (e.g. Hofland, Lech Kaminski, & Wolters, 2010). This is mainly due to the complexity, especially at large scale, of hanging large structures directly from the load cells.

The uncertainties regarding the measurements of pressures generated by impulsive wave conditions are very high (Bullock et al., 2007). Due to these uncertainties, generating correct measurements of wave impact loads on marine structures is probably one of the biggest challenges for coastal and maritime hydraulic modellers. Stagonas, Marzeddu, Gironella, and Sánchez-Arcilla (2016) performed a series of experiments to validate a pressure mapping system against classical measurements (pressure arrays and load cells). Their work showed that the spatial distribution of these impacts is characterized by pressure hot spots that appear slightly above the still water level (SWL). However, the horizontal position of the hot spots remains largely unpredictable (Bullock et al., 2007). The highly turbulent environment and the air/water mixture during wave breaking add further difficulty to the measurements. Furthermore, even if an almost perfect repeatability of the generated wave is guaranteed, this does not translate into perfect repeatability of the impact conditions (Stagonas et al., 2016).

The effects of sampling frequency on the measurement of wave induced impact loads poses another important problem (see e.g. Marzeddu, Gironella, Sachez-Arcilla, & Sutherland, 2014). At higher sampling frequencies, the magnitude of the measured load peak can vary up to 10 times. Bullock et al. (2007) studied the scale and laboratory effects on wave induced impact loads. Using a Froude-scale model, a good agreement was found between prototype and model results. Furthermore, they illustrated the importance of having the same spatial and temporal resolution (laboratory tests were performed with

lower spatial and temporal resolution) for a reliable comparison. According to Bullock et al. (2007), an unexpected behaviour of the measured loads at laboratory and prototype scales should not be excluded, due to the complex interactions occurring in the presence of air/water mixtures in a highly turbulent medium. A review of the literature on hydraulic tests dealing with wave-induced impact forces and pressures (Bullock et al., 2007; Marzeddu et al., 2014; Stagonas et al., 2016) revealed a limited understanding of the results variability, suggesting the need of a larger number of test repetitions for statistically robust and comparable datasets.

For most wave impact hydraulic tests, extreme long waves (d/L < 0.1) and wave periods between 15 and 25 s) in intermediate and shallow water-wave conditions need to be considered (Frostick, McLelland, & Mercer, 2011; Hughes, 1993). Therefore, the accuracy of the laboratory results depends on how precisely such nonlinear waves are generated. In many wave-impact hydraulic studies (e.g. Marzeddu et al., 2014; Stagonas et al., 2016), waves are generated by piston or flap wavemakers moving according to the linear wave theory (Biésel & Suguet, 1951). However, some studies (Goda, 1967, 1997; Morison, Crooke, & Washington, 1953; Oliveira, Sachez-Arcilla, Gironella, & Madsen, 2017) reported the presence of secondary harmonic free waves for extreme long wave conditions, when the waves were generated by sinusoidally moving a paddle. To prevent the generation of unwanted secondary waves, various authors proposed different wavemaker theories based on second-order wave theory (Hudspeth & Sulisz, 1991; Madsen, 1971; Schäffer, 1996), cnoidal wave theory (Goring, 1979), and stream function theory (Zhang & Schäffer, 2007).

In this study, we aim to understand the effects of hydraulic test repeatability on the pressure and force on vertical breakwaters under regular wave-induced impact conditions. For this, nearly breaking and broken wave conditions generated by cnoidal and Stokes first-order wavemaker theory have been analysed and compared. Nearly breaking waves generate higher impact pressures/forces, this clearly being a condition to be taken into account for the design of vertical breakwaters. On the other hand, broken wave attacks cause highly turbulent environments (complicating the measurement conditions) and are a common situation in designing rigid structures in shallow waters. Unlike linear waves, cnoidal waves represent relatively long waves in intermediate and shallow waters. Both types of wave theories are applied in this study to evaluate the laboratory effects on wave generation. Four different small-scale experiments have been carried out and repeated 120 times. High-speed recordings with load cells and pressure transducers are used in all the test runs.

This paper is divided into five sections. In Section 2 we present the set-up of the small-scale hydraulic tests. Results of the loads induced by nearly breaking and broken waves, as generated by cnoidal and Stokes first-order wavemaker theories, are described in Section 3. The discussion is in Section 4 and conclusions are in Section 5.

2 Physical model tests

2.1 Experimental layout

The experiments were carried out in the CIEMito wave flume at the *Laboratori d'Enginyeria Maritima* (LIM) of the *Universitat Politecnica de Catalunya* (LIM-UPC BarcelonaTech). The flume is 18 m long, 0.38 m wide, and 0.56 m deep; it is equipped with a piston type wave paddle that creates regular and irregular waves and a pumping system for currents. During the experiments, only 13.3 m of the flume was used and a vertical seawall model was placed at the end of a 1:15 smooth slope (Fig. 1). The water level was kept constant in all the experiments with a water depth of 0.285 m in front of the wave paddle and 0.115 m in front of the vertical wall. Figure 1a and 1c show the experimental layout for nearly breaking wave conditions. A $0.2 \times 0.38 \times 0.05$ m (length \times width \times height) block was placed at 0.5 m from the vertical wall to induce broken wave conditions (Fig. 1b and 1d).

2.2 Wave generation

Two types of regular waves were generated. The first type was based on cnoidal wavemaker theory (Goring, 1979), while for the second type, Stokes first-order wave generation (Biésel & Suquet, 1951) was used. According to Goring (1979), the piston paddle displacement to generate a free surface cnoidal wave, X_c , can be represented as:

$$X_c(t) = L\left(\frac{t}{T} - \frac{\theta_c}{2K(m)}\right) \tag{1}$$

where:

$$\theta_c = 2K(w) \left(\frac{t}{T} - \frac{X_c}{L}\right) \tag{2}$$

and L is the wave length, t is time, T is the wave period, and K(w) is the complete elliptic integral of the first kind and $w = HL^2/h^3$ is the elliptic parameter.

A Stokes first-order wave (linear wave) can be generated by a piston paddle displacement, X_l , as proposed by Biésel and Suquet (1951):

$$X_l(t) = 2a \left(\frac{\sinh(2kh) + 2kh}{8\sinh^2(kh)} \right) \sin(\omega t)$$
 (3)

where a is the wave amplitude, h is the water depth in front of the paddle, k is the wave number $(k = 2\pi/L)$, and ω is the angular frequency $(\omega = 2\pi/T)$.

For each type of regular wave, a train of three waves was generated. The first and third waves (ramp-up and ramp-down waves) were always smaller than the second wave (target wave height). The ramp-up wave was fully reflected by the structure, while the second wave resulted in an impact on the wall. Each wave train was repeated 120 times for each generation type (cnoidal and linear) and each impact condition (nearly breaking and broken waves), adding up to a total of 480 independent wave tests. Because the first wave of the train was always reflected by the structure, the impact of the second wave was considered the "cleanest" and most repeatable. The analysed wave could be considered as representative of an extreme wave in an irregular wave train. Thus, only the force/pressures induced by the impact of this wave were considered in the analysis.

2.3 Measurements

Eight resistance wave gauges were used for measuring the free surface evolution at 3.00, 5.00, 5.08, 5.20, 5.60, 10.50, 10.71, and 11.11 m from the wavemaker. A sampling rate of 100 Hz was used; the degree of accuracy of these sensors is

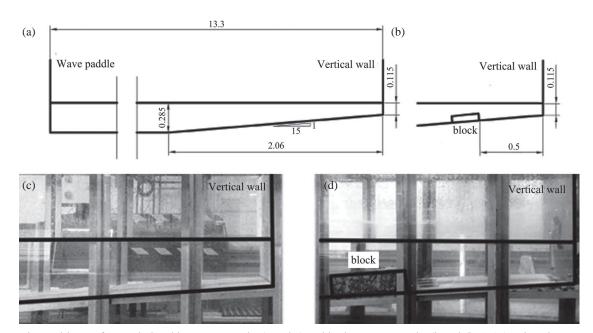


Figure 1 Experimental layout for nearly breaking wave attacks (a and c) and broken wave attacks (b and d). A 0.2 m length rectangular box was placed at 0.5 m from the vertical wall to induce broken waves onto the structure; all values in this figure are given in meter (m) unit

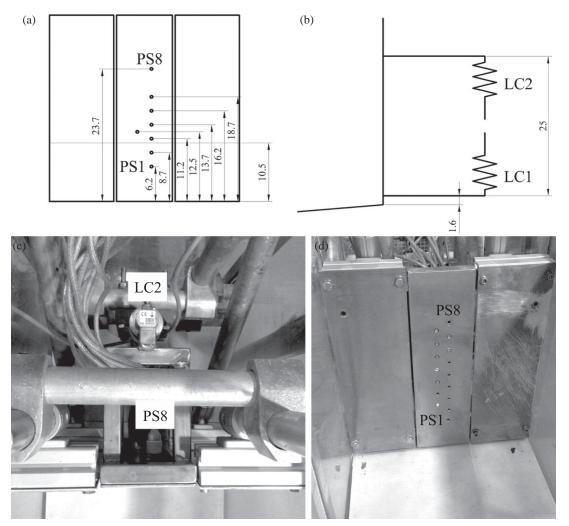


Figure 2 Experimental arrangement: (a) schematic frontal view of the wall showing the transducers (PS1–PS8); (b) schematic lateral view showing load cells (LC1–LC2); (c) top view photo of the instrumented vertical wall (d) front view photo of the vertical wall. All values in this figure are given in centimeter (cm) unit

approximately 0.001 m (Oliveira, Sanchez-Arcilla, & Gironella, 2012; Stagonas et al., 2016). The experimental arrangement consisted of an array of eight HBM P8AP pressure transducers (Hottinger Baldwin Messtechnik GmbH, Darmstadt, Germany) placed in the middle of the vertical wall. The vertical distances between the transducers (PS1 to PS3 and PS5 to PS8) are shown in Fig. 2a while an additional transducer PS4 was placed near the SWL and at a horizontal distance of 2.5 cm to the left of the array. The array was mounted on a segment of the wall (35 \times 10 cm) and supported by two load cells, LC1 and LC2 (Fig. 2b).

The force/pressure measurement systems were the same as those used by Marzeddu et al. (2014), Marzeddu, Gironella, and Sánchez-Arcilla (2013) and Stagonas et al. (2016). Marzeddu et al. (2013) carried out tests with non-impulsive conditions to evaluate the capability of P8AP and Z6FC3 sensors to record repeatable loads. Results showed a satisfactory repeatability of the measured loads with a root mean square error of 5 N m $^{-1}$, compared to normal range value of 150 N m $^{-1}$ (3%).

The HBM P8AP is an absolute pressure transducer for measuring static and dynamic liquid or gas induced pressures that

can be safely immersed in depths of up to 1 m. Each transducer is composed of a strain-gauge sensor and is provided already calibrated (by HBM), holding an EC declaration of conformity and a test certificate. The maximum measuring range of the P8AP sensors used for this study was 103 kN m⁻², with a reported accuracy of 0.3% of the maximum load, a 24 bit resolution and a natural frequency of the diaphragm of 12 kHz.

Force measurements were conducted using two Z6FC3 bending beam load cells (LC1-LC2 in Fig. 2b) with a nominal load of 50 kg, an accuracy of 0.009% of the maximum load and a 24-bit resolution. Similarly to pressure transducers, load cells were calibrated and could be immersed to a maximum depth of 1 m. An HBM QuantumX data acquisition system (Hottinger Baldwin Messtechnik GmbH) was used to simultaneously obtain samples from each load cell and pressure transducer at a sampling frequency of 4.8 kHz. However, the system could be amplified and sample up to 16 channels with a maximum sampling rate of 19.2 kHz. For experiments with pressure transducers, (Marzeddu et al., 2013,

2014) concluded that a satisfactory description of the pressure pulse can be obtained with a sampling frequency higher than 2.4 kHz.

3 Results

3.1 Wave height

Table 1 summarizes wave characteristics at 3 m from the paddle for cnoidal and linear waves. The values correspond to the second wave in the wave train (wave highlighted in Fig. 3a), representing nearly breaking and broken waves. The repeatability of the generated waves was tested; the root mean square

error (E_{RMSE}) and percentage error (computed on wave height) are shown in Table 1. The target wave height and period were H=0.16 m and T=2.3 s respectively. At a water depth of 0.285 m, the target wave presented a large Ursell number (Ur = 94.95) and low relative water depth (h/L=0.08), indicating highly nonlinear wave characteristics with corresponding implication for their evolution (Oliveira, Gironella, Sanchez-Arcilla, Sierra, & Celigueta, 2009; Oliveira et al., 2017). The wave periods of the generated cnoidal and linear waves matched the target period perfectly. The generated wave height was 0.120 and 0.155 m for cnoidal and linear waves, respectively. These differences between target and generated wave height can be due in part to the leakage around the paddle, which can decrease

Table 1 Incoming wave parameters, E_{RMSE} , and percentage of error on surface elevation measurements for 120 waves from each category. Wave characteristics and wave height error calculation were computed from wave probe number 1 at 3 m from the wave paddle

							E_{RMS}	<i>SE</i> (m)	Erro	or (%)
	H(m)	T(s)	d(m)	N° rep.	Ur	d/L	Max	Mean	Max	Mean
Target	0.160	2.3	0.285	120	95.0	0.08				
Cnoidal	0.120	2.3	0.285	120	71.2	0.08	0.0035	0.0011	2.9	1.0
Linear	0.155	2.3	0.285	120	92.0	0.08	0.0021	6.4e-4	1.3	0.4

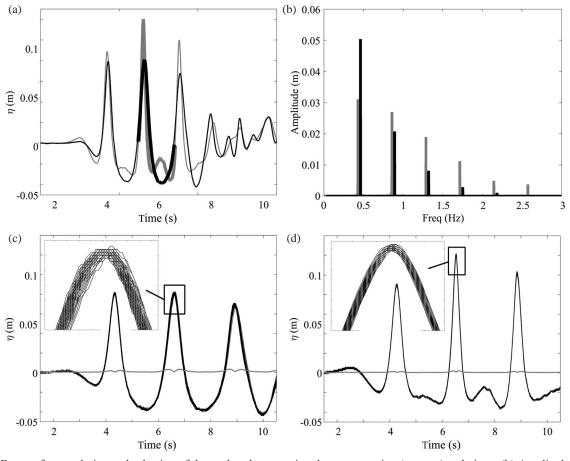


Figure 3 (a) Free surface evolution and selection of the analysed wave using the zero-crossing (zero-up) technique. (b) Amplitude spectrum (the frequencies for the cnoidal generation have been moved slightly forward in order to avoid overlap and make the figure clearer). Black and grey lines represent the results obtained with cnoidal and linear wave generation, respectively. (c) and (d) show for cnoidal and linear wave generation, respectively, the free surface evolution, η_i , of all 120 test repetitions (black lines) and the corresponding E_{RMSE} (grey line)

Table 2 Amplitudes (m) and energy (m² s²) calculated as the integral of the spectrums of the two types of wave attacks. Different intervals for the integral calculation were used in order to calculate the energy associated at each sub-harmonic

		0.43 Hz	0.86 Hz	1.29 Hz	Total spectrum
Cnoidal	Energy $(m^2 s^2)$				0.00103
	Amplitude (m)				0.06
Linear	Energy $(m^2 s^2)$	5.65e-04	3.54e-04	1.04e-04	0.00106
	Amplitude (m)	0.031	0.027	0.019	0.076

the amplitude of a generated wave (Hughes, 1993; Madsen, 1971). Wave heights in Table 1 were calculated as the difference between the maximum (wave crest) and the minimum (wave trough) free surface level. Regarding E_{RMSE} and error, hardly any variations were determined due to test repeatability (Table 1).

For each time instant, t, the E_{RMSE} was calculated using:

$$E_{RMSE}(t) = \sqrt{\frac{\sum_{i=1}^{n} (\eta_i(t) - \bar{\eta}(t))^2}{n}}$$
 (4)

where $\eta_i(t)$ is the value of the free surface for repetition i at time instant t, $\bar{\eta}(t)$ is the average free surface at time instant t, and n = 120 is the total number of repetitions.

In Fig. 3c and 3d all the 120 η_i time series and E_{RMSE} are plotted from t = 1.5 to t = 10.5 for cnoidal and linear waves, respectively.

The error is defined for each time instant by:

$$error(t) = \frac{E_{RMSE}(t)}{H_{\eta}}$$
 (5)

where $H_{\eta} = (\eta_{max} - \eta_{min})$ is calculated for the second wave of the train (target wave).

Figure 3a shows the differences between the free surface evolution at 3 m from the wave paddle obtained with cnoidal and Stokes first-order wavemaker theory. An unwanted free harmonic can be observed in the case of linear waves (two distinct waves can be seen in the trough area). The amplitude spectra of cnoidal and linear waves are shown in Fig. 3b. Table 2 summarizes the energy and amplitude for the first three harmonics. Although the total energy is similar in both types of generated waves, there are significant differences in energy distribution through the harmonics. The cnoidal wave height calculated as H_{η} , is similar to the cnoidal height approximation calculated as $H_a = 2a_1$. The values are 0.120 and 0.100 m, respectively, with η the surface elevation and a_1 the amplitude of the first harmonic. However, a huge difference was found for H_n and H_a for the linear waves: 0.155 and 0.062 m, respectively. This difference is due to the presence of undesired free harmonics, with amplitudes of the same order of magnitude as the first harmonic, (Madsen, 1971) in linear waves.

A graphic summary of the time history of nearly breaking and broken waves colliding directly with a vertical wall in a cnoidal wave test is shown in Fig. 4.

3.2 Nearly breaking wave attack loads

Figure 5 shows the vertical pressure distributions induced by the two types of generated waves for nearly breaking conditions. Dimensional (Fig. 5a) and dimensionless results (Fig. 5b and 5c) were considered. The maximum pressure in cnoidal waves was always recorded at PS5 or PS6, placed at + 0.025 and + 0.051 m above the SWL, respectively. In PS6, where maximum pressure was recorded in 86% of the repetitions, the ratio between the maximum and minimum pressure was 482%,

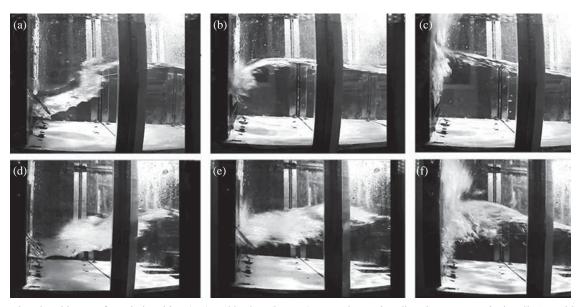


Figure 4 Snapshot time history of nearly breaking (top) and broken (bottom) waves impacting directly onto a vertical wall. A cnoidal wave case is seen in the photographs

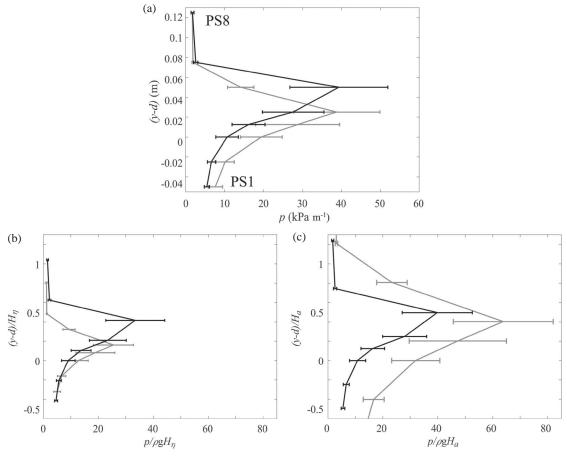


Figure 5 Comparison of vertical pressure distribution (means and standard deviations). Cnoidal (black) and linear (grey) wave generation theories: (a) dimensional plot; (b) dimensionless plot in relation to H_a . Nearly breaking wave attack conditions

and the maximum/mean pressure ratio was 209%. For linear waves, the maximum pressure was always recorded at PS4 or PS5, placed at + 0.125 and + 0.025 m above the SWL, respectively. In PS5, where the maximum was recorded in 78% of the repetitions, the maximum/minimum pressure ratio was 419%, and the maximum/mean pressure ratio was 178%.

The dimensionless pressures were calculated by dividing the recorded pressure by ρgH , where ρ is the water density and g is the acceleration due to gravity. Both H_n and H_a are used in the dimensionless analysis. The relative vertical position (yd/H was selected as a dimensionless parameter, where y is the distance to the bottom and d the water depth at the toe of the structure. For both maximum pressures and forces a number of parametric distributions have been tested (for both nearly breaking and broken waves) to find the best fit for the studied datasets. The selection of the "best distribution" has been done using a Bayesian information criterion (BIC). A chi-squared test has been performed to evaluate the goodness of fit of the "best distribution". All the observed populations accept the null hypothesis (the observed sample proportions are equal to those expected with a significance level of 0.05). The scores (p-values) of the chi-square test are reported in Tables 3–6.

The *Gamma* distribution probability density function presents the best fit for the maximum pressure recorded by PS6

in cnoidal waves and by PS5 in linear waves:

$$q_p = \frac{1}{b^s \Gamma(z)} p^{s-1} e^{-p/b} \tag{6}$$

where q_p is the probability density function, p is the maximum pressure, s is the shape parameter and b is the scale parameter. The results and Gamma distribution parameters for the pressures induced by nearly breaking waves are summarized in Table 3. The statistical distribution of the pressures, recorded by the sensors with the higher probability to record the maximum pressures (PS6 for the cnoidal and PS5 for the linear), with the Gamma distribution and the histograms of the two datasets are presented in Fig. 6.

Figure 7 shows the statistical distribution of the total force measured for the dimensional (Fig. 7a) and dimensionless (Fig. 7b and 7c) analyses. A *GEV* distribution was fitted and the histograms of the two sets of data are presented in the figure. In the case of cnoidal waves, the variability found for pressure was also observed for the force and the force application point. The maximum force was 168% of the minimum force and 133% of the most probable force value. For linear waves, similar ratios were determined, with a maximum recorded force 174% of the minimum and 138% of the most probable force value.

Table 3 Summary of the maximum pressure recorded for nearly breaking wave attack conditions

			Maxim	um pressu	ire							
		Position		Statistics	Recorded pressures				PDF Gamma distribution			
		Symbol	PS	Value		Symbol	Max	Min	Mean \pm SD	S	b	p-value
Cnoidal	Dimensional	y-d (m)	PS6	0.051	PS5 14% PS6 86%	$P(\text{kPa m}^{-1})$	82	17	39.4 ± 12.6	10.1	3.9	0.9
	Dimensionless H_{η}	$\frac{y-d}{H_{\eta}}$		0.42		$\frac{p}{ ho g H_{\eta}}$	70	14	33.4 ± 10.7		3.3	
	Dimensionless H_a	$\frac{y-d}{H_a}$		0.49		$\frac{p}{\rho g H_a}$	83	17	39.8 ± 12.7		3.9	
Linear	Dimensional	y-d (m)	PS5	0.025	PS4 22% PS5 78%	$P(kPa m^{-1})$	69	16	38.8 ± 11.1	12	3.23	0.92
	Dimensionless H_{η}	$\frac{y-d}{H_{\eta}}$		0.16		$\frac{p}{ ho g H_{\eta}}$	45	11	25.5 ± 7.3		2.13	
	Dimensionless H_a	$\frac{y-d}{H_a}$		0.4		$\frac{p}{\rho g H_a}$	114	27	63.8 ± 18.2		5.32	

Table 4 Summary of the maximum force recorded in the case of nearly breaking wave attack conditions

		Application point position			Recorded force				PDF GEV distribution				
		Symbol	Max	Min	Mean	Symbol	Max	Min	Mode	r	σ	μ	p-value
Cnoidal	Dimensional	y-d (m)	0.034	0.02	0.03	$F(\text{Nm}^{-1})$	2210	1311	1662	- 0.214	160	1626	0.09
	Dimensionless H_{η}	$\frac{y-d}{H_{\eta}}$	0.28	0.17	0.25	$\frac{F}{\rho g H_{\eta} d}$	16	10	12		1.182	12.009	
	Dimensionless H_a	$\frac{y-d}{H_a}$	0.34	0.2	0.3	$\frac{F}{\rho g H_a d}$	19	12	15		1.407	14.296	
Linear	Dimensional	y-d (m)	0.023	0.01	0.02	$F(\mathrm{Nm}^{-1})$	1885	1082	1367	-0.061	169	1350	0.25
	Dimensionless H_{η}	$\frac{y-d}{H_{\eta}}$	0.15	0.06	0.11	$\frac{F}{\rho g H_{\eta} d}$	11	6	8		0.934	7.723	
	Dimensionless H_a	$\frac{y-d}{H_a}$	0.23	0.1	0.17	$\frac{F}{\rho g H_a d}$	27	15	20		2.334	19.31	

Table 5 Summary of the maximum pressure recorded for broken wave conditions

			Max	imum pr	essure							
		Position		Statistics	Rec	PDF <i>Gamma</i> distribution						
		Symbol	PS	Value		Symbol	Max	Min	Mean \pm SD	S	b	p-value
Cnoidal	Dimensional	y– d (m)	PS2	- 0.025	PS2 42% PS3 38%	<i>p</i> (kPa m ^{−1})	26	4	8.2 ± 3.2	8.47	0.968	0.47
	Dimensionless H_{η}	$\frac{y-d}{H_{\eta}}$		- 0.17	PS4 15% PS5 5%	$rac{p}{ ho g H_{\eta}}$	22	3.4	7 ± 2.7		0.822	
	Dimensionless H_a	$\frac{y-d}{H_a}$		- 0.25	PS1.1% PS2 34%	$\frac{p}{\rho g H_a}$	26	4	8.3 ± 3.2		0.979	
Linear	Dimensional	y- d (m)	PS2	0	PS3 20% PS4 28%	<i>p:</i> (kPa m ⁻¹)	26	4	8.2 ± 3.2	11.03	0.502	0.07
	Dimensionless H_{η}	$\frac{y-d}{H_{\eta}}$		0	PS5 13%	$rac{p}{ ho g H_{\eta}}$	22	3.4	7 ± 2.7		0.33	
	Dimensionless H_a	$\frac{y-d}{H_a}$		0	PS6 4%	$\frac{p}{\rho g H_a}$	26	4	8.3 ± 3.2		0.826	

Table 6 Summary of the maximum force recorded for broken wave conditions

		Application point position			R		PDF GEV distribution					
		Symbol	Max Min	Mean	Symbol	Max	Min	Mode	r	σ	μ	p-value
Cnoidal	Dimensional	y - d(m)	0.02 - 0.014	0.007	$F(\text{Nm}^{-1})$	442	267	315	- 0.063	28	312	0.17
	Dimensionless H_{η}	$\frac{y-d}{H_{\eta}}$	0.18 - 0.12	0.06	$\frac{F}{\rho g H_{\eta} d}$	3.3	2.0	2.3	N/A	0.21	2.31	N/A
	Dimensionless H_a	$\frac{y-d}{H_a}$	0.20 - 0.14	0.07	$\frac{F}{\rho g H_a d}$	3.9	2.3	2.8	N/A	0.25	2.75	N/A
Linear	Dimensional	y - d(m)	0.19 - 0.04	0.08	$F(\mathrm{Nm}^{-1})$	380	215	278	-0.135	26	274	0.15
	Dimensionless H_{η}	$\frac{y-d}{H_{\eta}}$	1.23 - 0.25	0.52	$\frac{F}{\rho g H_{\eta} d}$	2.2	1.3	1.6	N/A	0.15	1.57	N/A
	Dimensionless H_a	$\frac{y-d}{H_a}$	3.17 - 0.66	1.30	$\frac{F}{\rho g H_a d}$	5.4	3.1	4.0	N/A	0.37	3.92	N/A

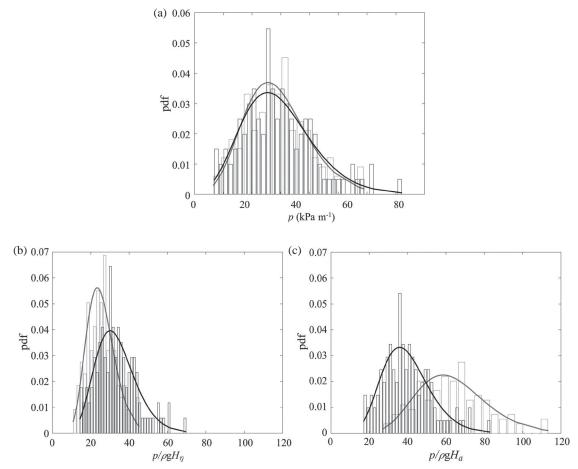


Figure 6 Theoretical *Gamma* pdf with histograms of the real datasets. Comparison of cnoidal (black) and linear (grey) wave generation theories. (a) Dimensional plot; (b) dimensionless plot in relation to H_{η} ; (c) dimensionless plot in relation to H_{α} . Nearly breaking wave attack conditions. The sensors with the highest percentage of recorded maximum pressures are the PS6 for cnoidal and PS5 for linear

The dimensionless forces presented in Fig. 7b and 7c were calculated by dividing the recorded force by ρgHd , using H_{η} and H_{a} .

Force results were fitted with a set of parametric probability distributions and the *GEV* function gave the best fit:

$$q_{F} = \left(\frac{1}{\sigma}\right) exp \left[-\left(1 + r \frac{F - \mu}{\sigma}\right)^{-1/r}\right] \times \left(1 + r \frac{F - \mu}{\sigma}\right)^{-1 - (1/r)}$$
(7)

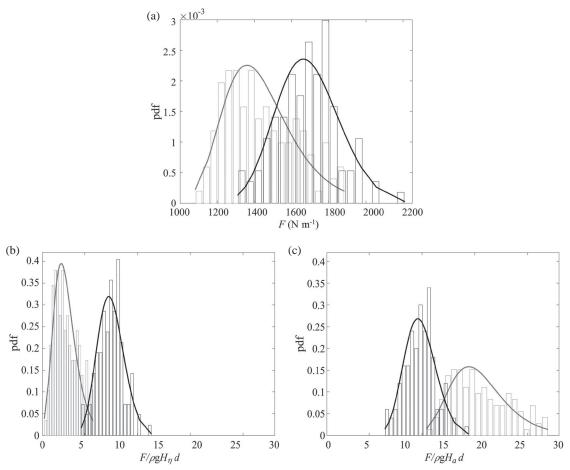


Figure 7 Theoretical GEV pdf with histograms of the real datasets. Comparison of cnoidal (black) and linear (grey) wave generation theories. (a) Dimensional plot; (b) dimensionless plot in relation to H_n ; (c) dimensionless plot in relation to H_a . Nearly breaking wave attack conditions

where q_F is the probability density function, F is the force, r is the shape parameter, σ is the scale parameter, and μ is the location parameter. The results and GEV parameters for the analysed force and application point datasets are summarized in Table 4.

3.3 Broken wave attack loads

Figure 8 shows vertical pressure distributions for broken wave attack. Dimensional (Fig. 8a) and dimensionless results (Fig. 8b and 8c) are presented. In the case of cnoidal waves, the maximum pressure was recorded by PS2 in 42% of the cases (PS2 was the most probable point). Although PS2 recorded more maximum pressures events (42%), PS3 recorded higher pressure values. This leads to higher mean maximum pressures at PS3 than PS2. At PS2, the maximum/minimum pressure ratio was 650%, while the maximum/mean pressure ratio was 318%. In the case of linear waves, the maximum pressure was recorded in 34% of the cases by PS2. For linear waves the maximum/minimum pressure ratio was 500%, while the maximum/mean pressure ratio was 291%. For broken waves, the Gamma distribution function provides a best fit to the maximum pressure recorded by PS2 for both the cnoidal and the linear cases (Eq. (6)). Table 5 summarizes the results for pressure induced by broken waves. The statistical distribution of the pressures (recorded by the sensors with the higher probability to register the maximum pressures) with the fitted *Gamma* distribution and the histograms of the two datasets are presented in Fig. 9.

Figure 10 shows the statistical distribution of the total force measured for broken waves. A *GEV* distribution was selected because, as in the case of nearly breaking waves, this is the distribution that provides the best fit. The histograms of the two sets of data are presented, showing that in the case of cnoidal waves, pressure variability (Figs 9 and 10) is clearly higher than force variability. The maximum force was 166% of the minimum force and 140% of the most probable force value. Similar ratios were determined for linear waves, with a maximum recorded force of 177% and 137% of the minimum and the most probable values, respectively. Results for the forces induced by broken wave conditions are summarized in Table 6.

4 Discussion

For nearly breaking cnoidal waves, a large range of pressures and forces was found. In this case, the maximum/minimum ratios for total force and maximum pressure were 168% and 482%, respectively. These high ratios highlight the significant

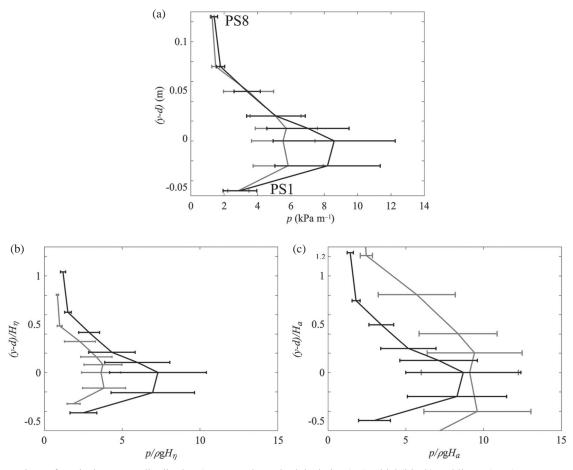


Figure 8 Comparison of vertical pressure distribution (means and standard deviations). Cnoidal (black) and linear (grey) wave generation theories. (a) Dimensional plot; (b) dimensionless plot in relation to H_n ; (c) dimensionless plot in relation to H_n . Broken wave attack conditions

variability of the phenomena and the difficulties of predicting reasonable values during the design process of vertical breakwaters when considering wave impact conditions. We found that the measurements of forces and maximum pressures do not follow the same statistical distributions. *GEV* and *Gamma* are the three parameters distributions that best fit the set of data for forces and maximum pressures, respectively. Variability was also observed for the location of maximum pressure and force application point. However, these positions are always located between half and one wave amplitude above the SWL.

For broken cnoidal waves, the maximum/minimum ratio of total force was 166%, similar to what was determined by nearly breaking cnoidal waves. However, this similarity was not observed for the maximum/minimum ratio of pressures. *GEV* and *Gamma* were also the distributions that provide a best fit for the total force and the maximum pressure for broken waves. For the broken cnoidal waves, the application force point and maximum pressure point are located between a half wave amplitude below and above the SWL.

When carrying out experiments with linear wave generation, the variability observed for total force, force application point, and point of maximum pressure are similar to those found for nearly breaking and broken choidal waves. However, the force application point and the point of maximum pressure have a slightly lower position in the case of linear waves. In all four tested conditions, the highest variations in pressure measurements were found close to the pressure peak point.

For linear waves (with the presence of unwanted harmonics), the method for defining wave height can significantly affect the dimensionless analysis, a typical engineering artefact to extrapolate laboratory results to a prototype. When considering wave height as H_{η} , for cnoidal and linear waves, dimensionless and dimensional quantitative analysis of pressure and force are similar. However, when considering wave height as H_a , dimensionless and dimensional results are significantly different for linear waves. Differences of over 100% in magnitude and positions of forces and pressures were observed between the two dimensionless analyses used for the linear waves. Far from being a comprehensive study on the repeatability of wave generated impact forces/pressures, these results highlight the significant variability of these phenomena and the challenges in evaluating the correct design loads for vertical structures.

The high variability found in the study could be partly attributed to the spatial density of pressure transducers and to the randomness of wave breaking phenomena. A low spatial density of pressure transducers may hinder the detection of the

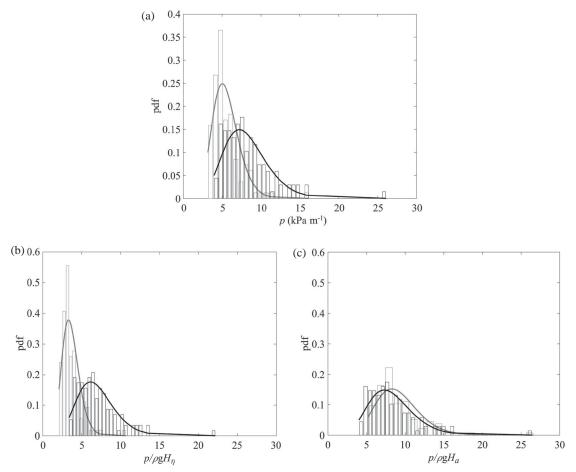


Figure 9 Theoretical *Gamma* pdf with histograms of the real datasets. Comparison of cnoidal (black) and linear (grey) wave generation theories. (a) Dimensional plot; (b) dimensionless plot in relation to H_{η} ; (c) dimensionless plot in relation to H_a . Broken wave attack conditions. The sensors with the highest percentage of recorded maximum pressures are the one analysed (PS2 for both cnoidal and linear)

maximum pressure point. As shown by the results a higher variability is found around the maximum pressure point. In this regard, a denser sensor distribution should decrease the uncertainties on the characterization of pressures induced by wave impacts on vertical breakwaters. Moreover, the lower variability found for force measurements, obtained with load cells, confirms the effect of sensor distribution on pressure characterization. The mixing of air and water during wave breaking could be considered as an almost random phenomenon and therefore hardly repeatable. Because of that, the randomness of wave breaking phenomena can induce variability both in pressure and force measurements.

It is important to remember that both scale and laboratory effects can play a major role in the evaluation of such loads. The hydraulic tests presented here correspond to typical work scales (between 1/50 and 1/80), and therefore scale effects are to be expected. Moreover, it is also important to separate the variability induced by sensor distribution and wave breaking from the uncertainties generated by laboratory effects, such as the wave generation method and the accuracy of sensors. For example, it would be possible to extrapolate our results to irregular wave tests, considering the analysed regular wave as the maximum wave height in an irregular wave train. The

probability of appearance of this type of loads is, in an irregular wave train, very low and a very small number of waves could be analysed (e.g. one to two waves) in comparison to the total number of generated waves (e.g. 1000 waves).

5 Conclusions

Test control and repeatability are essential to gain understanding in wave-induced pressure and forces on vertical structures. The analysis performed of impact conditions for nearly breaking and broken waves has allowed determination of total force variability, maximum pressure, and their corresponding locations. Based on four small-scale hydraulic experiments, probability distributions for regular wave-induced impact measurements (force and maximum pressure) are proposed. *GEV* and *Gamma* are the distributions that best fit the set of data for forces and maximum pressures, respectively. This proposal facilitates the setting up of hydraulic tests and selection of the appropriate wave conditions when designing vertical breakwaters. For all tested conditions, high variability in the measured pressure field and total force according to the random features of the studied phenomena were observed. The variability in the force

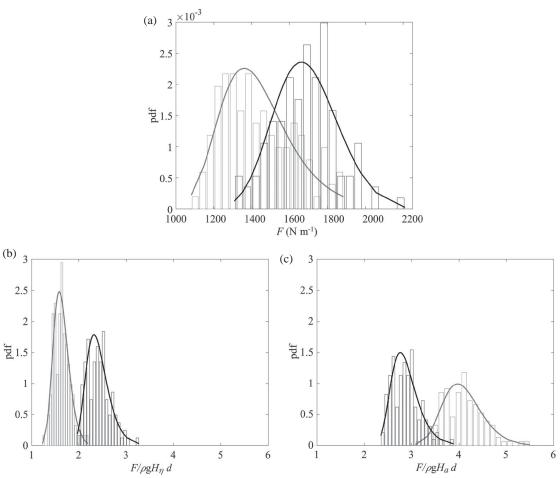


Figure 10 Theoretical *GEV* pdf with histograms of the real datasets. Comparison of cnoidal (black) and linear (grey) wave generation theories. (a) Dimensional plot; (b) dimensionless plot in relation to H_{η} ; (c) dimensionless plot in relation to H_{α} . Broken wave attack conditions

was similar for the nearly breaking and broken waves. The maximum measured force was between 166 and 177% of the minimum measured force and higher variability was found for more local data such as the maximum pressure. In this case, the maximum measured value was between 430 and 650% of the minimum measured value. For the force application point and maximum pressure point, variability was higher for broken waves due to the inherent randomness of breaking induced turbulence.

Performing the proposed type of analysis (120 repetitions for each wave attack condition) for long series of irregular waves would be extremely time-consuming. However, with a small number of irregular wave tests repetitions and the suggested non-dimensional distributions, it would be possible to have an idea of the variability. Hydraulic test repeatability studies should lead in the future to a significant reduction of experimental times and improved quality of the resulting structural design.

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Notations

a = wave amplitude (m)

b = scale parameter of the Gamma pdf(-)

d = water depth at the toe of the structure (m)

 E_{RMSE} = root mean square error

F = force (N m⁻¹)

G = gravity acceleration (m s⁻²)

h = water depth in front of the paddle (m)

H = wave height (m) k = wave number (m⁻¹)

K(w) = complete elliptic integral of the first kind (-)

L = wave length (m)

= maximum pressure (kPa m⁻¹) = force probability density function (-) q_F = pressure probability density function (-) q_p = shape parameter of the GEV pdf (-)= shape parameter of the *Gamma* pdf (-)S = time (s) T= wave period (s) Ur = Ursell number (-)= elliptical parameter (-) X_c = paddle displacement for a cnoidal wave (m) = paddle displacement for a linear wave (m) X_{l} = vertical dimension (m) $\Gamma(z)$ = gamma function (-)= water surface elevation (m) water surface η elevation (m) = location parameter of the GEV pdf (-) μ = water density (kg m^{-3}) ρ = scale parameter of the GEV pdf (-)σ

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= angular frequency (s^{-1})

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