NONSMOOTH PITCHFORK BIFURCATION IN A DC-DC CONVERTER: COEXISTING ATTRACTORS AND INTERMITTENCY

Abdelali El Aroudi\textsuperscript{1}, Vanessa Moreno-Font\textsuperscript{2}, and Luis Benadero\textsuperscript{2}

\textsuperscript{1}Departament d’Enginyeria Electrònica, Elèctrica i Automàtica (DEEEA), Universitat Rovira i Virgili Tarragona, Spain, e-mail: abdelali.elaroudi@urv.cat
\textsuperscript{2}Departament de Física Aplicada, Universitat Politècnica de Catalunya, Barcelona, Spain, (luis,vanessa)@fa.upc.edu

ABSTRACT
In this paper we deal with the analysis of nonlinear dynamical behavior of a single inductor two inputs two outputs (SITITO) power electronics DC-DC converter. The system can be used to regulate generally two outputs (one positive and one negative). Under certain operating conditions, the switching model can be approximated by a one dimensional piecewise constant vector field and, as a consequence, the corresponding map is piecewise linear (PWL). This model is derived and then it is used to study a nonsmooth pitchfork bifurcation in the system. Coexistence of attractors are detected by using the same model. Intermittent chaotic behavior is also addressed. Analytical results are confirmed by one dimensional and two-dimensional bifurcation diagrams.

Keywords: single inductor, DC-DC converters, nonlinear dynamics, bifurcations, chaos, pitchfork, nonsmooth.

1. INTRODUCTION
Switching power converters are widely used in the power management area, due to their potential of high efficiency, low cost and small size [1]. In most of the applications, these systems are used in situations where there is a need to stabilize an output voltage to a desired constant value.

The demand for portable equipments in modern vehicles such mobile phones, MP3 players, PDAs, GPS, has grown significantly during the last few years. The portable equipments usually include a variety of loads such as LCD displays, memories, microprocessors, Universal Series Bus (USB) and Hard Disk Drives (HDD). These loads require different operating voltages and load currents and are powered by the rechargeable batteries through DC-DC converters. To make longer the system run life and smaller its size, more and more system designers are focusing on improving the system power conversion efficiency with advanced power converters topologies. The traditional solution of using independent converters one for each of the outputs has the shortcomings of higher number of switches and magnetics components. Besides this option and other dual DC-DC converter configurations [2]-[9], single inductor multiple inputs multiple output (SIMIMO) DC-DC converters are, in general, convenient solutions for these low power applications.

Nowadays, there are many works dealing with nonlinear behavior in elementary stand alone [10] and other more complex power electronics circuits such as paralleled DC-DC converters [11], multi-cell and multi-level converters [12]-[14] and also for an example of single inductor two input two output (SITITO) converters [7].

In this paper, more insights into the modeling and analysis of a SITITO interleaved converter is presented. After giving the switched model, a systematic approach is described to obtain a simple PWL map that can be used to predict accurately the fast scale (switching) dynamics of the system. The model will be used to get some analytical conditions for stability of the system and for optimizing its performances.

The rest of the paper is organized as follows. Section II deals with the description of the SITITO DC-DC converter and the basis for the interleaving control circuit is presented. Then in Section III, the mathematical model is given. Under the assumption of perfect output regulation, the general PWL map is derived. Section IV will deal with stability analysis by using this PWL map. We will use the derived map to get some analytical expressions for stability conditions and to draw some bifurcation curves of the system. Section V includes some numerical simulations from the PWL map and one-dimensional and two-dimensional bifurcation diagrams of the system by using the PWL map. Basins of attraction are addressed in section VI. Finally, some conclusions are given in the last section.
2. SINGLE INDUCTOR TITO DC-DC REGULATOR

2.1. Power stage circuit

The schematic diagram in Figure 1 shows a DC-DC converter with a single inductor for two outputs $V_P$ and $V_N$ with opposite polarity. The voltages of both loads are regulated from the power source $V_{in}$ by generating a sequence of command signals driving switches $S_A$ and $S_B$.

![Figure 1](image)

**Figure 1.** Schematic diagram of a single inductor DC-DC converter with positive $v_P$ and negative $v_N$ output voltages.

2.2. Pulse Width Modulation Strategy with Interleaving

One control strategy for the SITITO topology of Figure 1, which is based on an interleaved current mode control, was proposed in [9] (See Figure 2). This control strategy works as follows: the two ramp reference signals, which are associated to each of the channels, are phase-shifted (interleaved) so that in the steady state regime, a sequence of time intervals, once per channel, will be produced. In this way, the ramp reference signal $v(t)$ ($j \in \{A, B\}$) corresponding to channel $j$ is shifted by a phase shift $\phi_j$. Based on current mode control, the general idea is to obtain a current reference $v(t)$ for each channel (output). These signals are to be compared to the inductor current $i_L$, Figure 2-b, shows the relevant control signals and parameters of the system.

![Figure 2](image)

**Figure 2.** Control strategy of single inductor two outputs converter of Figure 1 based on interleaved pulse width modulation. The feedback current is $i_L$. Control signals for the SITITO interleaved DC-DC converter

3. CLOSED LOOP MATHEMATICAL MODELING

3.1. Switched model for the SITITO converter

The switched model gives the set of five equations for the different configurations of the system. It contains the binary signals $u_A$ and $u_B$ and the instantaneous state variables. This model can be easily obtained by applying standard Kirchoff’s voltage law to the circuit [7]:

\[
\begin{align*}
\frac{dv_P}{dt} & = \frac{1}{C_P} \left( i_L(1-u_B) - \frac{v_P}{R_P} \right) \\
\frac{dv_N}{dt} & = \frac{1}{C_N} \left( -i_L(1-u_A) - \frac{v_N}{R_N} \right) \\
\frac{ds_P}{dt} & = \frac{1}{\tau_P} (V_P - v_P) \\
\frac{ds_N}{dt} & = \frac{1}{\tau_N} (v_N - V_N) \\
\frac{di_L}{dt} & = \frac{1}{L}(u_B-1)v_P + \frac{1}{L}(1-u_A)v_N + \frac{1}{L}(-r_Li_L + V_{in}u_A)
\end{align*}
\]

where $v_P$ is the voltage across the capacitor $C_P$, $v_N$ is the voltage across the capacitor $C_N$, $i_L$ is the current through the inductor $L$ whose equivalent series resistance is $r_L$. $s_P$ and $s_N$ are the integral variables and $\tau_P$ and $\tau_N$ are the time constants in the corresponding PI integrators. $u_A$ and $u_B$ are the driving binary signal which are 1 if the corresponding switch is ON and 0 if the switch is OFF. $v^A_P$ and $v^B_P$ are the reference currents that can be expressed as:

\[
\begin{align*}
v^A_P(t) & = g_P(V_P - v_P + s_P) + v^A_{ramp}(t) \\
v^B_P(t) & = g_N(v_N - V_N + s_N) + v^B_{ramp}(t)
\end{align*}
\]

where

\[
\begin{align*}
v^A_{ramp}(t) & = V_u - (V_u - V_i) \mod \left( \frac{t}{T} - 1 \right) \\
v^B_{ramp}(t) & = V_u - (V_u - V_i) \mod \left( \frac{t}{T} - \phi_B \right) + 1
\end{align*}
\]

mod($\cdot$) stands for the modulus function and $\phi_A = 1 - \phi_B$. The first subset of two equations (Eq. 1) refers to the dynamics of the voltage output $v_P$ and $v_N$. $R_P$ and $R_N$ are the resistive load for both outputs. Additional subset (Eq. 3 and Eq. 4) deal with two equations for each of the integral terms $s_P$ and $s_N$, being $V_P$ and $V_N$ the reference voltages. The last equation (Eq. 5) deals with inductor current dynamics. During each switching period, the
switch $S_A$ (resp. $S_B$) is closed for a time duration $d_A T$ (resp. $d_B T$). The duty ratios $d_A$ and $d_B$ for each ON sub-interval, which are determined by the action of comparator, can be expressed implicitly in terms of the state variables as follows:

$$i_p^A(t) - i_L = 0$$
$$v_p^B(t) - i_L = 0$$

where $i_p^A(t) = v_p^A(t)/r_{S,j}$, $j \in \{A, B\}$. Obtaining $d_A$ and $d_B$ requires solving Eqs. (10) and (11). If these equations are not feasible, the duty cycles are saturated to their limit values 0 or 1 depending on the relation between the control signal $r_{SjL}$ and the ramp signals $v_p^j$. The above set of five state equations is the closed switched continuous time model that can be used for computer simulations of the whole system.

3.2. Piecewise linear map

In this section we will give a one dimensional map that can capture the fast scale dynamic of the SITITTO interleaved converter. Under the assumption that outputs $v_p$ and $v_N$ are well regulated to their desired values $V_p$ and $V_N$ respectively, this map is a one-dimensional map and the state variable is the discrete-time current $i_n := i_L(nT)$. This assumption was done in [8] for elementary DC-DC converters. Due to the different operating modes that the system can present ($d_A$ and/or $d_B$ non-saturated and saturated), the one-dimensional map has different forms. Independently on the operating mode, it can be shown that the mapping can be written in the following form:

$$P(i_n) = P_A \circ P_B(i_n)$$

where $P_j$ is the local sub-mapping defined by:

$$P_j(x) = \begin{cases} x + m_{on} \phi_j T, & \text{if } d_j > \phi_j \\ x + m_{off}^{j}(\phi_j - d_j)T, & \text{if } d_j < \phi_j \end{cases}$$

where $x$ is the state variable at the beginning of phase $j$ and $m_{on}$ and $m_{off}^{j}$ are the approximates of the inductor current during the ON and the OFF phases and that are given by:

$$m_{on} = \frac{V_{in} - r_{L}I_{L}}{L}$$
$$m_{off}^{j} = \frac{V_{in} - V_{P} - r_{L}I_{L}}{L}$$
$$m_{off}^{N} = \frac{V_{N} - r_{L}I_{L}}{L}$$

In the expressions of $m_{on}$, $m_{off}^{P}$ and $m_{off}^{N}$ the term $r_{L}I_{L}$ which appears in the inductor current state equation was approximated by $r_{L}I_{L}$. The parameter $I_{L}$ is obtained by means of the averaged model and it is given by [7]:

$$I_{L} = \frac{V_{in}}{2r_{L}} + \sqrt{\frac{V_{in}^2}{4r_{L}^2} - \frac{1}{r_{L}} \left( \frac{V_{P}^2}{R_{P}} + \frac{V_{N}(V_{N} - V_{in})}{R_{N}} \right)}$$

Figure 3 shows the cobweb plot of the map $P$ for different values of $r_{S}$. Note the richness of behavior that the map can exhibit. Namely, in Fig. 3-a and Fig. 3-e, the system has only one real fixed point. The other two fixed points are virtual. For Fig. 3-b and Fig. 3-f, the virtual fixed points become real through a non smooth pitchfork bifurcation as we will see later. In Fig. 3-c and Fig. 3-g, two attractors are coexisting, one is chaotic and the other one is a fixed point. For Fig. 3-d and Fig. 3-h, a single chaotic attractor exists after a boundary crisis.

The duty cycles $d_A$ and $d_B$ are obtained by assuming that the inductor current is linear during each charging phase, then, the following expressions are obtained:

$$d_A = \frac{I^A_p - i_n^A}{(m_{on} - m_{r})T}$$
$$d_B = \frac{I^B_p - i_n^B}{(m_{on} - m_{r})T}$$

where $i_n^A = i_n$, $i_n^B = P_B(i_n)$, $I^A_p$ and $I^B_p$ are respectively the value of $i_p^A(t)$ and $i_p^B(t)$ at the beginning of phases $A$ and $B$, and

$$m_{r} = \frac{V_{L} - V_{L}}{r_{S}T}$$

The fixed points are obtained by forcing that $i_{n+1} = i_n$. Depending on the parameters of the circuit, there can exist one (nonsaturated), or three (one unsaturated and two saturated) fixed points (see Fig. 3). The unsaturated fixed point $i_d$ is given by the following expression:

$$i_d = I_{L} + \left( \frac{D_{A}^2 + D_{B}^2}{4} - \phi_{A}D_{A} - D_{B} \right) m_{on}T$$
$$+ \left( D_{B} \left( 1 - \frac{D_{B}}{2} \right) - \phi_{B} \left( \phi_{A} + \phi_{B} \right) \right) m_{off}T$$
$$+ \left( \phi_{A}D_{A} - \frac{D_{A}^2}{2} - \phi_{B}^2 \right) m_{off}T$$

where the average duty cycles corresponding to each phase can be obtained from the averaged model and they are given by [9]:

$$D_{A} = \phi_{A} - \frac{V_{N}}{R_{N}I_{L}}$$
$$D_{B} = \phi_{B} - \frac{V_{P}}{R_{P}I_{L}}$$

The two other fixed points when they exist are given by:

$$i_{dA} = I_{L}^A - m_{r}\phi_{B}T + m_{off}A m_{on} - m_{r}A T$$
$$i_{dB} = I_{L}^B - (m_{on} - m_{r}) \left( \phi_{B} - \frac{m_{off}}{m_{on} - m_{r}} \right) T$$

These fixed points can only exist if the derivative of the map at the nonsaturated fixed point is bigger than 1.
Figure 3. The form of the map $P$ and the cobweb plot near the fixed points for different values of $r_S$. (a, e) $r_S = 1 \, \Omega$, $V_P = 18 \, V$, (b, f) $r_S = 2 \, \Omega$, $V_P = 18 \, V$, (c, g) $r_S = 3 \, \Omega$, $V_P = 18 \, V$ and (d, h) $r_S = 5 \, \Omega$, $V_P = 21 \, V$. (a-d) overview of the PWL map. (e-h) zoom near the fixed points.

4. NONSMOOTH PITCHFORK BIFURCATION

4.1. The stability indexes

A sufficient condition for stability of the nonsaturated fixed point is that the absolute value of the derivative $\lambda$ of the map $P$ evaluated at this fixed point is smaller than 1. It can be shown that this derivative is:

$$\lambda(i_d) = \frac{(m_{P,off} - m_r)(m_{N,off} - m_r)}{(m_{on} - m_r)^2}$$  \hspace{1cm} (25)

If for a certain choice of parameters, this derivative becomes larger than one then, the two saturated fixed points are created (they pass from virtual to real). The derivative of the map at $i_{dA}$

$$\lambda(i_{dA}) = \frac{(m_{P,off} - m_r)}{(m_{on} - m_r)}$$  \hspace{1cm} (26)

Likewise, the derivative of the map at $i_{dB}$ is

$$\lambda(i_{dB}) = \frac{(m_{N,off} - m_r)}{(m_{on} - m_r)}$$  \hspace{1cm} (27)

Let us define the following set of parameters:

$$\kappa = \frac{L(V_u - V_l)}{r_S T}$$  \hspace{1cm} (28)

$$\alpha = r_L I_L - V_{in} + V_P$$  \hspace{1cm} (29)

$$\beta = r_L I_L + V_N$$  \hspace{1cm} (30)

$$\gamma = r_L I_L - V_{in}$$  \hspace{1cm} (31)

With these definitions the stability indexes $\lambda$ become

$$\lambda(i_d) = \frac{(\kappa - \alpha)(\kappa - \beta)}{(\kappa - \gamma)^2}$$

$$\lambda(i_{dA}) = \frac{\kappa - \alpha}{\kappa - \gamma}$$

$$\lambda(i_{dB}) = \frac{\kappa - \beta}{\kappa - \gamma}$$  \hspace{1cm} (32)

It can be observed that as $\lambda(i_d) = \lambda(i_{dA})\lambda(i_{dB})$ and that after a pitchfork bifurcation $\lambda(i_d) > 1$, it is always true that $|\lambda(i_{dA})| > 1$ or $|\lambda(i_{dB})| > 1$. This means that the saturated fixed points cannot be both stable. This characteristic is different from the case of a smooth pitchfork bifurcation where the stability of the created fixed points is symmetric.

4.2. Nonsmooth pitchfork bifurcation

If for a certain choice of parameter values one has $\lambda(i_d) = 1$ and $\lambda(i_{dA}) < 0$ and $\lambda(i_{dB}) < 0$, then a nonsmooth pitchfork bifurcation will occur. The critical value of $r_S$ for which this bifurcation occurs is given by the expression:

$$r_{S,cri} = \frac{L(V_u - V_l)(V_{in} + V_P - V_N)}{T(V_N(V_{in} - V_P) + r_L I_L(V_{in} + V_P - V_N) - V_{in}^2)}$$  \hspace{1cm} (33)

This expression will be plotted later together with the 2-D numerically obtained bifurcation diagram.
5. BIFURCATION BEHAVIOR FROM NUMERICAL SIMULATION

Let us consider the circuit of Figure 1 with the control scheme of Figure 2. Let us consider also the values of parameters shown in Table 1. The bifurcation parameters used here are the current gain sensor \( r_S \) and the reference voltage for the positive output channel \( V_P \). The bifurcation parameter \( r_S \) is varied in the range \((1, 5) \) \( \Omega \) and \( V_P \) is varied in the range \((1, 21) \) \( \text{V} \). The 2-D bifurcation diagram is shown in Fig. 4. In this figure, the period of the orbit was color coded. Only five colors were used to determine the different periodic modes that the system can present. Using more colors does not change the results because the system does not present stable periodic orbit with period bigger than four. Yellow color represent the useful stable zone (nonsaturated fixed point \( i_d \)). Brown color corresponds to the the saturated fixed point \( i_{dA} \). Zones of higher periods and chaotic behavior are also shown as parameters \( r_S \) and \( V_P \) vary. Their period are labeled in the same diagram. One can observe that there exists a period doubling region for low values of \( V_P \) and pitchfork bifurcation region for high values of \( V_P \). These regions are also labeled in the same figure. The period doubling region is studied by the authors in a separate paper. Here the study is limited to the pitchfork bifurcation region. In the coexistence region, a chaotic attractor coexists with the saturated fixed point \( i_{dA} \).

Figure 5 shows the one dimensional bifurcation diagram. The fixed points are also plotted in the same figure. We can see clearly that the branches of the saturated fixed points, after the nonsaturated fixed point becomes unstable, form a shape like a pitchfork bifurcation. However, this is a nonsmooth bifurcation and as a consequence the created fixed points have not the same stability. As a result, instead of coexisting two stable fixed points after the bifurcation, here a chaotic attractor coexists with a stable fixed point. The size of chaotic attractor increases between \( r_S = 2 \) \( \Omega \) and \( r_S = 3.75 \) \( \Omega \) when this attractor undergoes a boundary crisis (the attractor collides with \( i_d \) which is a boundary between the basins of attraction) and disappears. Between \( r_S = 3.75 \) \( \Omega \) and \( r_S = 5 \) \( \Omega \) only the saturated fixed point attractor exists and the basin of attraction of the chaotic attractor becomes null. In this case, the basin of attraction of the fixed point is the entire domain of definition of the PWL map.

### Table 1. Parameter values used in numerical simulations.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value</th>
<th>Parameter</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>( V_{in} )</td>
<td>12 ( V )</td>
<td>( V_f )</td>
<td>0</td>
</tr>
<tr>
<td>( L )</td>
<td>640 ( \mu H )</td>
<td>( V_u )</td>
<td>1</td>
</tr>
<tr>
<td>( r_L )</td>
<td>0.7 ( \Omega )</td>
<td>( V_P )</td>
<td>+18 ( V )</td>
</tr>
<tr>
<td>( R_N )</td>
<td>33 ( \Omega )</td>
<td>( V_N )</td>
<td>−22 ( V )</td>
</tr>
<tr>
<td>( R_P )</td>
<td>16.5 ( \Omega )</td>
<td>( f_s )</td>
<td>10 ( kHz )</td>
</tr>
<tr>
<td>( r_S )</td>
<td>varying</td>
<td>( \phi_A = \frac{1}{2} )</td>
<td>( \phi_B = \frac{1}{2} )</td>
</tr>
</tbody>
</table>

Figure 4. 2-D bifurcation diagram of the map \( P \) taking \( r_S \) and \( V_P \) as bifurcation parameters. The diagram is plotted starting from an initial condition near the average value of inductor current \( I_L \).

Figure 5. 1-D bifurcation diagram taking \( r_S \) as bifurcation parameter. The diagram is plotted starting from different initial conditions in order to uncover the possible coexisting attractors. The fixed points and their stability are also shown.

#### 5.1. Coexisting period-1 orbit and chaotic attractor:

\[ \lambda(i_{dA}) < -1 \text{ and } -1 < \lambda(i_{dB}) < 0 \]

As in the case of a smooth pitchfork, two fixed points are created and, depending on their eigenvalues, we can have the possibility of two attractors for the same parameter values. The basin of attraction, which is defined as the set of points in the state space such that initial conditions chosen in this set dynamically evolve to an attractor, can be obtained by sweeping the initial condition in a certain interval. These sets can be obtained by numerical simulations. The evolution of the basins of attraction in terms of the bifurcation parameter \( r_S \) is shown in Figure 6. The boundaries between different basins of attraction are dif-
different unstable periodic orbits with different periodicity. The expressions of these unstable orbits can be obtained in closed form but they hold too much space and they are not given here. Initials conditions selected in the white zones will make the system to evolve to the stable non-saturated 1-periodic orbit. Initials conditions selected in the gray zones will drive the system to the stable saturated 1-periodic orbit $i_{dA}$ while those selected in the black zones will drive the system to the chaotic attractor.

Figure 6. Evolution of the basins of attraction in term of the bifurcation parameter $r_S$.

5.2. Intermittent chaotic behavior: $\lambda(i_{dA}) < -1$ and $\lambda(i_{dB}) < -1$

If both saturated fixed points exist and are unstable ($\lambda(i_{dA}) < -1$ and $\lambda(i_{dB}) < -1$), then a chaotic behavior is possible. The size of the attractor is larger for smaller values of $\lambda(i_{dB})$. When the attractor collides with the nonsaturated fixed point, its size is suddenly enlarged and the attractor goes near the fixed point $i_{dB}$ which is also a repeller point and again the trajectory moves to the fixed point $i_{dA}$ region. This moving from one side of the nonsaturated fixed point to another gives rise to intermittent behavior as shown in Fig. 7.

6. CONCLUSIONS

In this paper a single inductor two inputs two outputs (SITITO) DC-DC converter with a power stage configuration under current mode control and interleaving is studied. Its dynamics is described by a nonlinear modeling approach. An expression for a simple PWL map is derived to obtain accurate information about the dynamics of the system. This map can predict many nonlinear behaviors such as bifurcations, chaos, coexisting attractors and the associated hysteresis phenomenon and intermittent behavior. A nonsmooth pitchfork bifurcation is detected and analyzed for the system. Closed form expressions for the critical parameter are given in suitable design parameter space. Numerical simulations confirm the theoretical predictions.

7. REFERENCES


