Abstract—This work proposes the use of magnetic coupling for powering autonomous sensors in space-constrained applications, such as occupancy and belt detection in removable vehicle seats. The power demand of the autonomous sensor is considered between tens and hundreds of milliwatts. A theoretical analysis first highlights the critical parameters in order to achieve a large powering range and high efficiency. Series-resonant tanks are considered for both the primary and secondary networks. Because the intended application is space-constrained, small coils have to be used. In order to increase their quality factor, commercial ferrite-core coils are used. A class D power amplifier is proposed for the primary network. Experimental results show that a power of tens of milliwatts can be transferred to a 100 Ω load placed at the secondary network up to a distance of 2 cm, near seven times the radius of the coils (3 mm). The addition of a rectifier and a voltage regulator at the secondary network in order to properly power an autonomous sensor (3 V @ 30 mA) limits the powering range to 1 cm. Overall power efficiencies around 45 % and 20 % are achieved respectively at distances of 5 mm and 1 cm.


I. INTRODUCTION

Remote or wireless power transmission via inductive coupling has been around for a long time. High-power transfer includes battery recharging of electrical vehicles [1] and a broad range of industrial applications [2] whereas low-power transfer includes RFID systems [3], biomedical implants [4], or portable consumer electronic products [5].

Transmission power distance is, in general, shorter than the diameter of the powering coils. Recently, though, the possibility of effectively powering at higher distances has been demonstrated. In particular, a power transfer of 60 W with 40 % coil-to-coil (60 cm in diameter) efficiency over distances in excess of 2 m (ca. 3 times the diameter of the coils) was shown in [6]. The analysis is rather based on physical theory and more engineering focused approaches using circuit lumped circuits have appeared since then [7]. The same principle has also been explored for powering: multiple receivers from a single transmitter coil [8], biomedical implants [9], or even chips [10].

Vehicles can also benefit from inductive powering. In particular, some vans and minivans incorporate removable seats in order to flexibly arrange their internal space. Wiring that seats in order to incorporate, for instance, seat belt detectors, can become unpractical. So, in some vehicles a passive detection is performed via magnetic coupling. The addition of wireless power could allow the incorporation of new devices that require some amount of power, such as a seat occupancy sensor or a microcontroller that adds intelligence to the removable seat. In these applications, the available or acceptable space for the coils is rather limited.

Section II first presents a theoretical analysis of the power transferred to the load with a pair of magnetic coupled resonators. Powering range distance and power efficiency are also analyzed. Then, Section III presents the selected commercial coils. Quality factors are measured for different frequencies and exposure regulations are verified through simulations. Section IV presents the design of the primary and secondary networks. The primary network includes a class D power amplifier whereas the secondary network includes a rectifier stage and a voltage regulator in order to properly power the autonomous sensor. Section V shows the measured performance. Finally, Section VI concludes the work.

II. THEORETICAL ANALYSIS

Fig. 1 shows the equivalent circuit of a pair of magnetically coupled series resonators, being the left-hand and right-hand networks the primary and secondary, respectively. \( V_1 \) is a sinusoidal signal that models the output of the power amplifier that drives the primary network; \( I_1 \) and \( I_2 \) stand for the currents of the primary and secondary; \( L_1 \) and \( L_2 \) model the coils; \( C_1 \) and \( C_2 \) are the added capacitances to work at resonance; \( R_1 \) models the output resistance of the power amplifier, \( R_{1L} \) and \( R_{2L} \) model the losses of the coils, and \( R_{load} \) models the load; and finally \( M \) models the mutual inductance between the coils, where

\[
M = k \sqrt{L_1 L_2}
\]
being \( k \) the coupling factor between the coils. We neglect the losses of the capacitors as they usually are much lower than that of inductors.

\[ k_c = \frac{1}{\sqrt{Q_1 Q_2}}. \]

The parameter \( k_c \) is referred as critical coupling [12]. Whenever \( k = k_c \), \( \eta_1 = 0.5 \), i.e. \( R_{eq} = R_1 \), thus maximum power is transferred from the primary to the secondary circuit, being the load power

\[ P_{Load,\text{max}} = \frac{V_i^2}{4R_1} \eta_2. \]

The parameter \( d_c \) can be defined as the critical distance between the coils at which (10) holds. A lower \( k_c \) leads to a higher \( d_c \). By increasing \( Q_1 \) and \( Q_2 \) in (10), \( k_c \) is reduced and \( d_c \) is increased, thus increasing the powering distance range. For \( k > k_c \), i.e. \( d < d_c \), \( P_L \) decreases but \( \eta_1 \) monotonically increases towards the unity. The parameter \( \eta_2 \) does not depend on \( k \) and thus remains constant over the distance.

III. COILS, COUPLING FACTOR AND FIELD EXPOSURE

For the intended application, small-size coils are sought. At the same time, in order to comply with the reference levels for general public exposure to time-varying electric and magnetic fields [13], and to reduce the power losses of the power amplifier, frequency resonance was limited to less than 150 kHz. So, in order to increase the quality factor of the coils with these constraints, the use of magnetic-core material was considered as an appropriate solution.

Figure 2. Dimensions of the selected 1 mH coils (in millimeters). Source: http://www.fastrongroup.com/.
In order to obtain the experimental quality factors of the coils at different frequencies, we used a series-resonant network and measured the resistance at the resonant frequency by using an HP4194A impedance analyzer. Appropriate values of capacitors were used in order to tune the resonant frequency. Fig. 3 shows both the measured resistances and the resulting quality factors. As can be seen, resistance values increased with frequency, which is due to the joint combination of skin and proximity effects and the losses of the ferrite [14]. Quality factor increased steeply at low frequencies achieving a maximum (ca. 40) around 40 kHz.

![Figure 3. Values of resistance and quality factor for the selected commercial coil of 1 mH.](image)

In order to estimate the coupling factor ($k$) over the separation distance ($d$) of the coils, we used the simulation program COMSOL. Fig. 4 shows an axisymmetric model for the primary and secondary coils, and illustrates the parameter $d$. Sizes of the inductors were in accordance with those presented in Figure 2. The contour areas R1 to R6 were defined as ferrite whereas C1 and C2 (wire coil) were defined as copper. A relative permittivity ($\mu_r$) of 2000 was used for the ferrite. Spherical domain boundaries were used and set to zero magnetic insulation.

![Figure 4. Modeling of the primary and secondary coils using COMSOL.](image)

IV. CIRCUIT DESIGN

A. Primary Network

Fig. 6 shows the circuit schematic of the primary network. We used a class D power amplifier based on a low-cost commercial self-oscillating half-bridge driver (IR2153) and two external N-channel MOSFETs (BS108), M1 and M2. The driver, powered ($V_{cc}$) at 12 V dc (battery voltage in vehicles), alternatively activates the two MOSFETs, thus injecting a square wave signal into the series-resonant network. The oscillation frequency is selectable via an RC network ($R_b, C_b$) up to 1 MHz. A potentiometer was used to fine tune the desired resonant frequency. Following the manufacturer guidelines, a bootstrap capacitor ($C_c$) was used to properly activate M1.

![Figure 6. Circuit schematic of the primary network.](image)

For moderate to high quality factors, only the first harmonic will generate a current through the network, being its amplitude of $7.64 \, V \left(2V_{cc}/\pi \right)$ and its rms value of 5.4 V.
The MOSFET manufacturer publishes an on resistance of 8 Ω (@ Vgs = 2.8 V). In our case, we measured a lower value, around 4 Ω, due to the higher value of Vgs. This value corresponds to R in Figure 1.

B. Secondary Network

Fig. 7 shows the schematic circuit of the secondary network. As can be seen, in order to obtain a DC signal, a full bridge rectifier was jointly used with a stabilization capacitor (C0). An ensuing voltage regulator (LP2980) was added to provide a voltage (Vr) of 3 V across the load (RL). This voltage value is appropriate for low-power commercial transceivers, where current consumption is in the order of units to tens of milliamps. Here we consider a range of 3 mA to 30 mA. This leads to an equivalent RL of 100 Ω to 1 kΩ. A 10 V zener diode Ds was used for protecting the voltage regulator from overvoltages.

Figure 7. Circuit schematic of the secondary network.

V. EXPERIMENTAL SETUP AND RESULTS

Fig. 8 shows the mechanical setup fabricated to fix the distance between the primary and secondary networks, which were implemented in two separate PCB boards. The main support and the fixing screws were made of nylon. For the measurements, distance was adjusted manually from 0 cm to 3 cm in 0.5 cm steps. For the series-resonant networks, we used the commercial coils of 1 mH presented in section III and capacitors (C1, C2) of 1.8 nF. The frequency of the primary driver was adjusted to the experimental resonant frequency.

Figure 8. Mechanical setup used to fix the distance between the primary and the secondary networks.

First, only the resonant tank illustrated in Fig. 1 was used at the secondary network. Resistors of 100 Ω and 1 kΩ were used for RL. For each distance both PLoad and η were estimated. PLoad was estimated by measuring the voltage drop across RL with a floating oscilloscope. η was estimated by dividing PLoad by the generated power from the 12 V dc source of the primary network.

Fig. 9 shows PLoad. As can be seen, the critical distance, dc, increased from 0.5 cm for RL = 1 kΩ to 1 cm for RL = 100 Ω. This agrees with the theoretical predictions and simulations of section II and III. From (6), an increase of RL leads to a decrease of Qc. Thus, from (10), kc decreases leading to an increase of dc (see Fig. 5). The value of PLoad,max for RL = 100 Ω was 138 mW, which nearly agrees with the predicted value of 134 mW obtained from (11). This predicted value was obtained considering Vr = 5.4 V, Rs = 4 Ω (see section IV), and RL = 36 Ω. The value of RL was estimated by measuring the voltage drop across a resistor momentarily added in series with the resonant tank. At a distance of 2 cm, near seven times the internal radius of the coils (3 mm), ca. 55 mW were transmitted to the load.

Figure 9. PLoad for L1=L2=1 mH at a resonance frequency of 117 kHz.

Fig. 10 shows power efficiency. As can be seen, efficiency was higher for RL = 100 Ω than for RL = 1 kΩ down to 0.5 cm. Efficiencies for RL = 100 Ω at 0.5 cm and 1 cm were around 50% and 40%, respectively. Conversely, for a distance of 0 cm efficiency decreased for RL = 100 Ω because of the losses of the half-bridge driver of the primary network, ca. 15 mW. This power loss was relatively less important for RL = 1 kΩ because the higher value of the load power at 0 cm (see Fig. 9), achieving an efficiency near 60%.

Figure 10. η for L1=L2=1 mH at a resonance frequency of 117 kHz.
Then, the circuit of Fig. 7 was used at the secondary network with resistors of 100 \( \Omega \), 1 k\( \Omega \), and 1 M\( \Omega \) for \( R_L \). Now, for each distance, \( V_L \) (Figure 11.) was measured and power efficiency was again estimated (Fig. 12). For \( R_{load} = 100 \Omega \), the desired voltage of 3 V, and thus a load power of 90 mW, was achieved for distances from 0.5 cm to 1 cm. Corresponding efficiencies ranged from 45% to 20%. For \( R_{load} = 1 \text{k}\Omega \), the desired voltage of 3 V, and thus a load power of 9 mW, was achieved for distances up to 1.5 cm. Efficiencies were rather low in this case. Additionally, \( R_{load} = 1 \text{M}\Omega \) was considered, which emulates the case when the autonomous sensor demands a low current (3 \( \mu \)A), i.e. is in a sleep mode. Here, a suitable voltage was also achieved for distances up to 1.5 cm. Efficiency was now extremely low, which is logical considering the losses of the primary driver and the low power demanded by the load (9 \( \mu \)W).

![Figure 11. Load voltage (\( V_L \)) across \( R_L \) for the secondary network of Fig. 7.](image)

![Figure 12. Overall efficiency for the secondary network of Fig. 7.](image)

**VI. CONCLUSIONS**

The principle of magnetic coupling resonance has been recently proposed to power portable devices. Here, we have used such a principle for exploring the possibility of powering autonomous sensors, e.g. seat belt and occupancy detectors, in removable vehicle seats. First, a theoretical analysis highlights the need of using high quality coils in order to achieve a large powering range and high efficiency. Additionally, the application is space-constrained. Thus, commercial ferrite-core coils of small-size have been used. The quality factor of the coils has been measured and simulations have demonstrated that the magnetic field is below the reference levels for general public exposure at distances higher than 5 mm.

A class D power amplifier has been used for the primary network. Experimental results have shown that a power of tens of milliwatts can be transferred to a load of 100 \( \Omega \) placed at the secondary network up to a distance of 2 cm, near seven times the radius of the coils (3 mm). The addition of a rectifier and a voltage regulator in order to properly power an autonomous sensor (3 V @ 30 mA) limits the powering range to 1 cm. Overall efficiencies around 45% and 20% have been achieved at distances of 5 mm and 1 cm, respectively.

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